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SPACE DEPENDENT MODEL  
FOR THE SLOWING DOWN OF FAST NEUTRONS

Thomas Martin Mills



United States  
Naval Postgraduate School



THEESIS

SPACE DEPENDENT MODEL FOR  
THE SLOWING DOWN OF FAST NEUTRONS

by

Thomas Martin Mills

Thesis Advisor:

T. J. Williamson

June 1971

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Space Dependent Model for  
the Slowing Down of Fast Neutrons

by

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Submitted in partial fulfillment of the  
requirements for the degree of

MASTER OF SCIENCE IN PHYSICS

from the

NAVAL POSTGRADUATE SCHOOL  
June 1971

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## ABSTRACT

The slowing down of fast neutrons was analyzed by a multi-group method of discrete time and energy states coupled with a spatial harmonic expansion method to determine the neutron density in a homogeneous, isotropically scattering slab. Five neutron source geometries were studied for both a fissioning and a non-fissioning system.

Numerical results were obtained for the neutron flux, mean neutron energy and the neutron spectra for the one dimensional system using a harmonic mode expansion of up to six terms to determine the time-energy-space dependence.



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## I. INTRODUCTION

Existing and future experimental work in the study of pulsed fast-neutron assemblies require the ability to determine the time dependent solution of the neutron transport equation accurately in the time period extending from the nanosecond to the millisecond range. This work allows the experimenter to compare both measured response and the analytical predictions for the overall improvement of each. Safe operation and control of fast nuclear reactors demands a fast, accurate mathematical simulation.

Current research on the response of fast assemblies to pulsed neutron sources considers that the time dependence can be represented by an exponential decay of the form:

$$T(t) = \sum_i A_i \exp(-\lambda_i t)$$

where the constants  $A_i$  and  $\lambda_i$  can be related to the composition and geometry of the system. A common assumption is that all terms above the first or fundamental will decay very quickly and that all measurable quantities can be described by a single fundamental or pseudo-fundamental value for  $\lambda_i$ . Experimental effort is being directed to answering the question of whether or not such a pseudomode does in fact exist, and if not, to understand why not.

This work will investigate several of the aspects of this problem.



In their movement through a material, neutrons undergo interactions, scattering or absorption, with the atoms and suffer changes in direction and velocity as the result of both elastic and inelastic collisions. While this is a random process for an individual neutron, the overall process for considering the time history and motion of the bulk of the neutrons is amenable to a statistical analysis. The Boltzman time-dependent transport equation:

$$\frac{\partial N}{\partial t}(\underline{r}, v, \underline{\Omega}, t) = -v \underline{\Omega} \cdot \nabla N(\underline{r}, v, \underline{\Omega}, t) - v \Sigma_t(v) N(\underline{r}, v, \underline{\Omega}, t) + \int_{\underline{\Omega}} \int v' \Sigma(v') N(\underline{r}, v', \underline{\Omega}', t) \cdot f(v', \underline{\Omega}' \rightarrow v, \underline{\Omega}) d\underline{\Omega}' dv' + S(\underline{r}, v, \underline{\Omega}, t)$$

can describe the net processes occurring in any elemental volume of a physical system to account for all changes in the density of neutrons, subject to the usual conventions:

$\underline{r}$	Position of the volume element $dV$ .
$v$	Speed.
$\underline{\Omega}$	Direction. (The vector velocity would be defined as $\underline{v} = v \underline{\Omega}$ ).
$N(\underline{r}, v', \underline{\Omega}', t)$	Density of neutrons at time $t$ having velocity $v' \underline{\Omega}'$ .
$\Sigma(v')$	Speed dependent cross section for neutrons with speed $(v')$
$f(v', \underline{\Omega}' \rightarrow v, \underline{\Omega})$	Probability that a neutron having velocity $(v' \underline{\Omega}')$ will interact to become one velocity $(v \underline{\Omega})$ .



The neutron density variation with time considers:

(a) leakage of neutrons:

$$-v\Omega \cdot \nabla N(\underline{r}, v, \underline{\Omega}, t),$$

(b) losses due to interactions (scattering or absorption) in the volume element:

$$-v\Sigma_t(v) N(\underline{r}, v, \underline{\Omega}, t),$$

(c) scattering of neutrons into the element phase space (position-velocity-direction) due to interaction of neutrons having other initial parameters:

$$\iint v' \Sigma_t(v') N(\underline{r}, v', \underline{\Omega}', t) C_t f(v', \underline{\Omega}' - v, \underline{\Omega}) dV' d\underline{\Omega}'$$

neutrons introduced into the element of phase space from a source in that element:

$$S(\underline{r}, v, \underline{\Omega}, t).$$

Complete derivations of the neutron transport equation can be found in Tait [1] and Davison [2] and is discussed in detail in many other texts on nuclear reactor physics.

During the past twenty years, numerous methods and approximations have been formulated to solve this integro-differential equation, either in a closed analytical form, or by a numerical approximation, for both the time dependent and time-independent cases. While a major portion of this effort has been in obtaining the solutions applicable to thermal reactors, uranium-235 systems, the need to conserve



fuels into the distant future requires fast breeder reactors. Adequate understanding of the complete time, space and energy dependence of fast neutron populations requires the development of new analytic and simulation models for the fast neutron transport and fission problem.

Analytical solutions are often based on extension of the existing models for thermal systems. These solutions usually assume  $1/v$  or constant neutron cross sections, with the age-diffusion theory [3] considering the steady state, time independent cases and with perturbation analysis for time dependent solutions [4]. These provide representations of the steady state systems, and simple kinetic models for fast reactors.

(1) Numerical methods generally deal with the time independent solutions of the transport equation by two principal methods:

Multi-group theory accounts for all energy dependent parameters by dividing the range of neutron energies into some finite number of intervals and defining average values of the energy dependent terms for each energy group. This method is frequently used to study the energy and time relations and the energy and position dependence [5]. This method has been used to consider up to several hundred groups between fission energies and the thermal condition [6].

(2). Monte Carlo is a statistical method to consider the interactions of individual neutrons to determine principal integral or system values for several spatial regions using multiple energy groups. The study of the individual



neutron "life history" from its birth to its eventual loss, allows determination of such parameters as the effective multiplication factor ( $K_{\text{eff}}$ ). Computational requirements for this method limit its usefulness mainly to time independent calculations, and to a limited extent [7], the time dependent problem using very few velocity groups.

The description of the slowing down and thermalization of neutrons as a probabilistic Markov chain process was proposed by Perkel [8] in 1960, and expanded to consider the thermal energy region (0.00002 to 1.034 eV) by Ohanian and Daitch [9] for the study of the time dependent neutron spectra. Jenkins and Daitch [10] have extended this technique to formulations of the study of a pseudo-fundamental mode decay for time-energy response of pulsed fast systems, assuming that the system response can be represented by the lowest Fourier spatial mode of the time dependent diffusion equation for two energy regions: .001 to 1.0 eV and 1.0 to 1000 Kev.

Williamson and Albrecht [11] have extended the Markov chain method of treatment for the slowing down problem to cover the continuous range from 10.5 MeV to thermal energies.

Menzel, *et al*, [12] have employed the method of Ohanian and Daitch to provide a complete harmonic expansion method for the study of the space-time-energy response of pulsed systems in the interval 0.01 to 1.0 eV.

The purpose of this thesis is to apply the numerical methods of the Markov chain process to describe the energy



and time response of pulsed fast systems, with the harmonic expansion techniques for the spatial dependence similar to Menzel, in the interval 10.5 Mev to thermal energies.

Current work on the pulsed neutron response of fast assemblies, considers that the time decay can be represented by some exponential function of the form:

$$T(t) = \sum_i A_i \exp(-\lambda_i t)$$

where the constants  $\lambda_i$  and  $A_i$  can be related to the composition and geometry of the assembly.

In this study, the fundamental and higher harmonic space modes are combined with energy-time expressions to determine the effectiveness of this technique in simulating the experiments conducted in the laboratory. Experiments suggest that the decay of the neutron population is often not best described by the pseudo-fundamental mode. Hopefully, this work will aid in determining if this is due to the effects of a constantly varying velocity spectrum, or to the persistence of higher spatial modes of the time response remaining for longer periods of time.



## II. THE SLOWING DOWN EQUATION

### A. THE BOLTZMAN TRANSPORT EQUATION

Starting with the general form of the time dependent transport equation, reasonable assumptions which will permit this intergro-differential equation to be restated as the slowing down equation are:

- (1) The neutron density  $N(\underline{r}, v, \underline{\Omega}, t)$  is isotropically distributed in  $\underline{\Omega}$  at all points in the physical media.
- (2) All sources and scattering kernels are isotropic.
- (3) No delayed neutron sources.
- (4) The physical media is composed of locally homogeneous isotropic materials.
- (5) Spatial dependence is included within the bounds of the diffusion approximation.
- (6) Boundary conditions are that the neutron density is zero at the extrapolated boundary and that the extrapolation distance is constant for all energies.

Conditions 1, 2 and 4 permit the equivalence of the expression for the local streaming or leakage of neutrons from a volume element to be described as:

$$-v\underline{\Omega} \cdot \nabla N(\underline{r}, v, \underline{\Omega}, t) = D(v) \nabla^2 N(\underline{r}, v, \underline{\Omega}, t) = -B^2 D(v) N(\underline{r}, v, \underline{\Omega}, t)$$

where  $B^2$  is called the geometric buckling.

Within this analysis, a source that is a delta function in time will be considered as the mathematical equation of



the pulsed neutron source with a very small time duration, which can be written as:

$$S(\underline{r}, v, \underline{\Omega}, t) = S'(\underline{r}, v, \delta(t))/4\pi.$$

Adjusting all constant terms to account for the isotropic conditions of the geometry under study, and to match the primary condition that the neutron density at a time ( $t=0$ ) will be equivalent to the value of the source strength at that time we have:

$$N(\underline{r}, v, 0) = S(\underline{r}, v, 0) = \sum_n R_n(\underline{r}) F_n(v, 0),$$

where solutions of the form  $N(\underline{r}, v, t)$  serve as Green's functions for the pulsed conditions with an arbitrary time and spatial parameter.

Starting with the basic assumption that the neutron density at any time can be described as an infinite sum of solutions of the form:

$$N(\underline{r}, v, t) = \sum_n R_n(\underline{r}) F_n(v, t),$$

it is possible to transform each of the terms of the general transport equation to the slowing down equation in the following manner:

$$(1) \frac{\partial N(\underline{r}, v, \underline{\Omega}, t)}{\partial t} = \sum_n R_n(\underline{r}) \frac{\partial F_n(v, t)}{\partial t}.$$



$$(2) -v \underline{\Omega} \cdot \nabla N(\underline{r}, v, \underline{\Omega}, t) = -D(v) \nabla^2 N(\underline{r}, v, t)$$

$$= -D(v) \sum_n \nabla^2 R_n(\underline{r}) F_n(v, t).$$

$$(3) \int v' \Sigma_t(v') \int N(\underline{r}, v', \underline{\Omega}', t) C_t, f(v', \underline{\Omega}' \rightarrow v, \Omega) dv' d\underline{\Omega}' =$$

$$\int v' \Sigma_t(v' \rightarrow v) \sum_n R_n(\underline{r}) F(v', t) dv'$$

which reduces to the now modified form of the slowing down equation:

$$\sum_n R_n(\underline{r}) \left\{ \frac{\partial}{\partial t} F_n(v, t) + \Sigma_t(v) F_n(v, t) - \int v' \Sigma_t(v' \rightarrow v) \cdot F_n(v', t) dv' \right\} + \sum_n D(v) F_n(v, t) \nabla^2 R_n(\underline{r}) = S'(\underline{r}, v, t)$$

substituting for the spatial-velocity dependent delta function source, the eigenfunctional expression of the slowing down equation is obtained:

$$\sum_n \left\{ R_n(\underline{r}) \left\{ \frac{\partial}{\partial t} F_n(v, t) + \Sigma_t(v) F_n(v, t) + \int v' \Sigma_t(v' \rightarrow v) \cdot F_n(v', t) dv' \right\} + D(v) F_n(v, t) \nabla^2 R_n(\underline{r}) \right\} = \sum_n R_n(\underline{r}) F_n(v, 0).$$



Consolidating all expansion functions of the index- $n$  in the single summation,

$$\sum_n \left[ R_n(\underline{r}) \left\{ \frac{\partial F_n(v, t)}{\partial t} + \sum_t(v) F_n(v, t) + \int v' \sum_t(v' - v) \cdot \right. \right. \\ \left. \left. F_n(v', t) dv' - F_n(v, 0) \right\} + D(v) F_n(v, t) \nabla^2 R_n(\underline{r}) \right] = 0$$

it is now possible to separate the spatial dependence from the energy time relations via the harmonic buckling factor,  $B_n^2$ , to two expressions:

$$\frac{\nabla^2 R_n(\underline{r})}{R_n(\underline{r})} = + B_n^2$$

and:

$$-D(v) B_n^2 F_n(v, t) = \frac{\partial}{\partial t} F_n(v, t) + \sum_t(v) F_n(v, t) \\ + \int v' \sum_s(v' - v) F_n(v, t) dv dt - F_n(v, 0).$$

Final clearing and rearrangement of all terms yield the decoupled expressions:

$$\nabla^2 R_n(\underline{r}) - B_n^2 R_n(\underline{r}) = 0$$

$$\frac{\partial}{\partial t} F_n(v, t) + B_n^2 D(v) F_n(v, t) + \sum_t(v) F_n(v, t)$$

$$+ \int v' \sum_s(v' - v) F_n(v', t) dv' = 0$$

with  $R_n(\underline{r})$  a time independent function and  $F_n(v, t)$  a space independent function.



## B. MOD-5: THE DISCRETE STATE APPROACH

In the separated form of the slowing down equation, the velocity-time harmonic term must satisfy the same form of the intergro-differential equation as the space independent neutron density function.

$$\frac{\partial}{\partial t} N(v, t) = -v D(v) B^2 N(v, t) - v \Sigma_t(v) N(v, t) \\ + \int v' \Sigma_s(v' \rightarrow v) N(v', t) dv' + S(v, t).$$

Any method that provides a solution to the space independent slowing down equation can be applied to solve the velocity-time harmonic equation.

Williamson and Albrecht [11,13] have developed a stochastic model for neutron moderation that provides a numerical solution of the slowing down of fast neutrons in a finite media. Since the slowing down process is a continuous time process in which collisions may occur at any time, having a continuous range of energy transition possibilities, it may be classified as a continuous time, continuous state Markov process. If the neutron cross sections remain constant with time, the slowing down may be classed as a stationary Markov system.

Application of the formalism of the Markov process calculations provides the capability to perform straightforward computer solutions for a discrete time-velocity model of the space independent equation. The probability that a neutron will be in a finite width velocity state,  $v_i$ , at a time,  $t_j$ ,



is calculated by a computer program, MOD-5, which follows the slowing down of neutrons by determining a probability density parameter  $F(v_i, t_j)$  which is the probability that a neutron will be in the velocity state  $v_i$  at the discrete time  $t_j = n\Delta t$ .

The Markov process describes the variation of a probability density vector,  $s(n\Delta t)$ , which describes the probability that a neutron will be located at some state in the system during its entire lifetime from birth to death. This state vector is given as:

$$\bar{s}(n\Delta t) = s_1, s_2, s_3, s_4, \dots, s_N$$

where  $s_i$  is the probability that the neutron is in the state bounded by velocities  $v_i$  and  $v_{i-1}$  (where  $v_i < v_{i-1}$ ) at the time  $n\Delta t$ . This method presumes that the initial state vector, neutron source,  $\bar{s}(0)$ , can be described for time  $t=0$ .

The problem description necessary for application of a MOD-5 type of approximation are:

- (a) The velocity range of interest is divided into "M" discrete states.
- (b) Collision physics is applied to construct an array,  $P(i,j)$ , of the discrete Markov transition probabilities of a neutron going from velocity state-i to state-j in the discrete time period  $\Delta t$ .

While the details of the Markov transition matrix are discussed in several sources [14], the main considerations are that the array for the slowing down approximation is



upper triangular with only positive elements, with the elements of a given row summing to unity, which states that all possible outcomes of the process are considered.

If the initial Markov transition matrix describes the probability for a transition within a specified time interval,  $\Delta t$ , the probabilities for a transition at a time  $\{n\Delta t\}$  can be determined as:

$$P_{n\Delta t} = (P_{\Delta t})^n.$$

The probability of a neutron being in any state at a time  $(-n\Delta t)$  can be determined from the  $(n-1)$ th step in the matrix product as:

$$\bar{s}(n\Delta t) = \bar{s}((n-1)\Delta t) \cdot P\Delta t.$$

Both the program MOD-5 and the harmonic expansion method employ several series of notations and can result in some conflict with existing literature. In this work, the following definitions and terms will be used:

$\bar{s}(n\Delta t)$  the probability density vector of a neutron being in any state in the system at the time  $n\Delta t$ .

$s_{i,j}$  the probability that a neutron will be in a state "i" at a time  $t_j$ ,

where the time and total density vector can be written as:

$$t_j = n\Delta t$$

$$\bar{s}(t_j) = (s_{1,j}, s_{2,j}, \dots, s_{I,j}).$$



While this form of notation applies to the case of a single decay mode, the expansion to a multiple harmonic expansion method requires the addition of another indexing parameter-n to indicate the particular harmonic mode. This results in the probability density expressions being written as:

$s_{n,i,j}$	- probability of the neutron fraction of the n-th harmonic mode being in a velocity state $v_i$ at the time $t_j$ .
$F_n(v_i, t_j)$	- n-th harmonic term of the time velocity dependence of the neutron probability density.

The contribution of each harmonic term of the neutron density is the product of the spatial component and the probability density component:

$$N_n(x, v_i, t_j) = R_n(x) F_n(v_i, t_j)$$

where the space component,  $R_n(x)$ , is a measure of the fractional contribution to the neutron population of the n-th harmonic mode, and the term,  $F_n(v_i, t_j)$  is the probability of a neutron being in the velocity state  $v_i$  at time  $t_j$  for the n-th mode of the time decay harmonic.

In the following discussions, the equivalence of the terms is described as:

$$s_{n,i,j} = F_n(v_i, t_j) \Delta v_i$$

where  $\Delta v_i$  is the width of the velocity state about the value  $v_i$ .



### III. GEOMETRICAL MODELS

#### A. SOURCE FUNCTIONS

The main objective of this thesis is the investigation of the harmonic mode expansion method for the space, time and energy dependence of the neutron density in a physical system. Simple source functions representative of actual systems are desired to demonstrate this method. Three categories of source type were considered: a finite volume source and a first collision source as interior sources in the system, and a wide beam exterior source condition.

The test geometries are:

(a) Finite Volume Source - a finite volume element, located on an axis of symmetry and at an off axis position is considered to represent a hole or cavity in the slab. This might be envisioned as a beam port or localized neutron generator capable of providing a pulse of monoenergetic neutrons in a small volume in the slab. This is represented as a fourier square pulse on the x-axis of the system.

(b) First Collision Source - this exponential distribution of neutrons in the system is considered as a more realistic approach for the response of a physical system than the fixed source in volume condition. Rather than assume that the initial burst of neutrons would remain in the fixed volume element, it can be argued that they would initially travel throughout the system with an exponential form.



As with the finite volume type source, this source type is studied for both the symmetric and unsymmetric positions.

(c) Exterior Source Beam - a higher level of complexity in the source geometry selection process is the physical situation where the assembly under study is subject to a wide beam of fast neutrons from an external point. This can be considered as a very simple approximation to a first collision response from a plane source of neutrons on the surface of the slab. While this may be considered to follow an exponential type distribution through the slab, it is approximated as a straight line or ramp function within the confines of the slab.

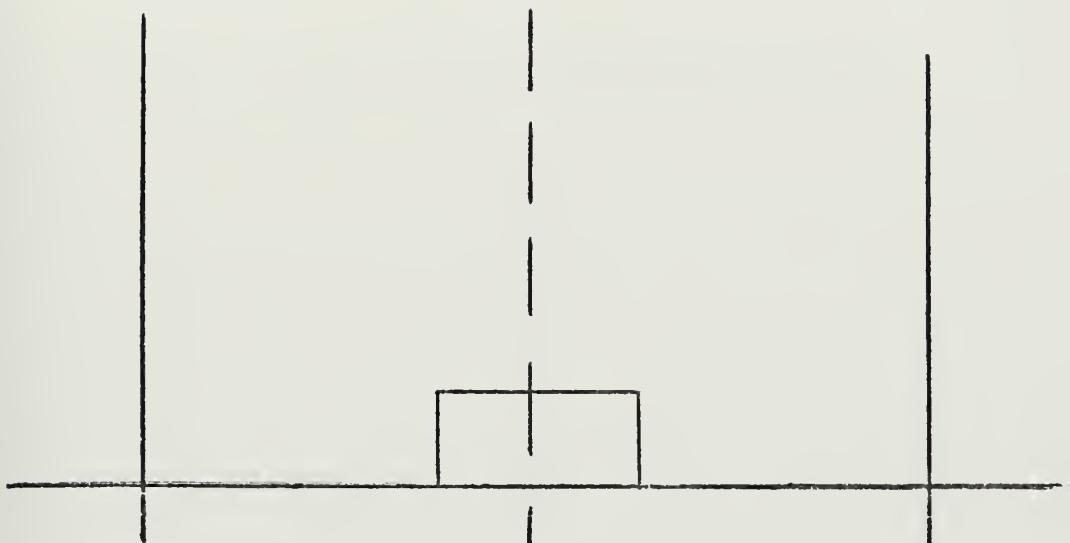
#### B. CHOICE OF DIMENSIONALITY

The principal goal of this work is the development of a useful mathematical model to determine the spatial dependent response of a physical system to a burst of fast neutrons. In this determination, it will be necessary to refer to the neutron density in terms of three main parameters - time, space and energy. The derivation of all equations and principal relationships will be written for the time, position and velocity of the neutrons. Equivalent expressions can easily be obtained using either the neutron energy (E) or lethargy (u), using the relationships:

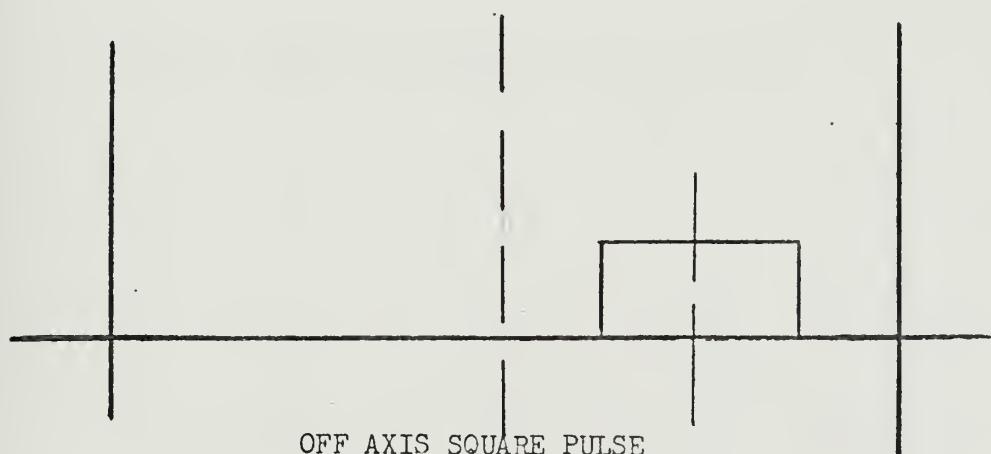
$$E = \frac{1}{2} m_n v^2$$

$$u = -\ln(E_0/E) = -2 \ln(v_0/v).$$





CENTRAL AXIS SQUARE PULSE

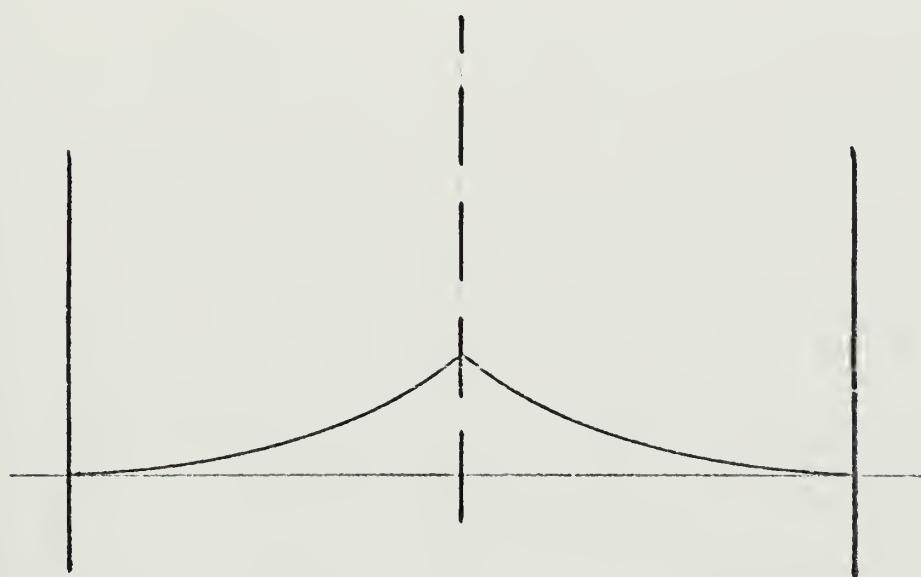


OFF AXIS SQUARE PULSE

FIGURE 1 SOURCE GEOMETRY SQUARE PULSE



CENTRAL AXIS FIRST COLLISION-EXPONENTIAL SOURCE



OFF AXIS FIRST COLLISION-EXPONENTIAL SOURCE

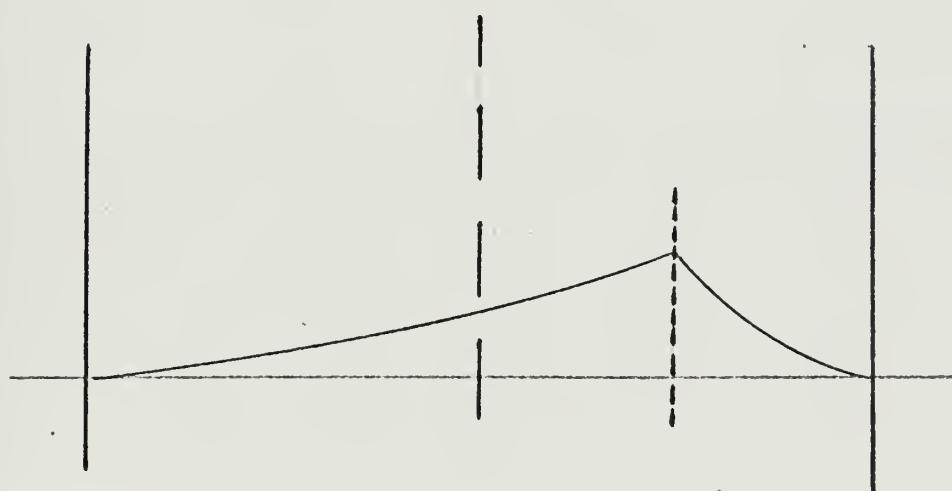


FIGURE 2 SOURCE GEOMETRY FIRST COLLISION SOURCE





FIGURE 3: SOURCE GEOMETRY EXTERIOR SOURCE RAMP FUNCTION



Employing the appropriate expressions, it can be easily shown that equivalent terms for the neutron density,  $N(\underline{r}, v, t)$ , the time-space-velocity form can also be modified to give the time-space-energy or time-space-lethargy form. While the velocity dependent form is used for the purposes of demonstration within this paper, all calculations performed by the programs MOD-5 and MIL-6 employ the lethargy dependent forms.

In the determination of numerical solutions for the slowing down equation, reasonable methods and simple systems are desired to demonstrate the harmonic expansion method for the time, velocity and spatial expressions, where the method is general enough to be easily revised to handle more difficult and complex problems. As shown previously, the neutron density can be determined by the summation of a harmonic series whose terms are the product of a space function and a time-velocity function.

$$N(\underline{r}, v, t) = \sum_n R_n(\underline{r}) F_n(v, t)$$

where the spatial dependent function,  $R(\underline{r})$ , satisfies the differential equation:

$$\nabla^2 R_n(\underline{r}) - B_n^2 R(\underline{r}) = 0.$$

In order to retain the flexibility to study and compare the response of physical systems to various source types, the selection of the dimensionality for problem selection posed some difficult choices.



The spatial functions must form an orthogonal set, a condition easily met by the Fourier expansion methods for the one dimensional problem using trigonometric functions for the cartesian geometry, but requires the use of regular Bessel functions for cylindrical geometry, and the spherical Bessel and spherical harmonic functions for the spherical geometry. A practical consideration was to obtain a balance demonstrating the harmonic expansion application with reasonable computing requirements. The more complex the spatial functions, the longer computing time and large core storage space required. The trigonometric functions can be determined quickly, while the existing methods available for the Bessel function calculations require approximately three times as much computing time.

In all three of the principal geometrical coordinate systems (cartesian, cylindrical and spherical), studies of sources located at an axis of symmetry resolve themselves into some form of the one dimensional problem. However, the study of sources located off the axis of symmetry of the system, while seeming reasonable conditions to analyze, can lead to some very messy mathematics, such that the utility of the method can become lost in the details. While these problems can be considered, the one dimensional case would allow easy comparison of all symmetrical and non-symmetrical cases without initially confusing the reader or user of this technique prior to demonstrating the method. After accomplishing that aim, the multi-dimensional cases (two or three dimensions) can be investigated.



A practical matter in the consideration of the harmonic expansion method involves the retention of data within the computing machinery while running the basic problem or one of the many variations that can be considered by this method.

An important consideration while developing the computer programming for these calculations involved the selection of methods to transfer data between the two computational routines. The program MIL-6 was written to process data input from punch cards, nine-track magnetic tape, or the IBM-2314 series of magnetic disk. The program MOD-5 was run sequentially to determine the velocity and time characteristics of the physical system for each harmonic mode of the buckling constant ( $B_n^2$ ). This generally requires one minute of computing time for the IBM-360/67 with 70 velocity states, requiring 200k bytes of core storage for each harmonic mode. Thus, a six-mode expansion would require approximately six minutes of central processor time.

Assuming that the problem under study is going to deal with 150 discrete velocity states, 20 time step values and a minimum of six harmonic modes of the functional data,  $F_n(v, t)$ , as produced by program - MOD-5, one is faced with the retention of 18000 numbers which are required either in core storage or in an on-line direct access device. In the IBM-360/67, all real numbers require a 4 byte location, and a minimum problem would require 72k bytes (12 per mode)



available during the running of the problem. Once the primary calculation of the density vector  $N(\underline{r},v,t)$  is done, data retention is required for the total system  $N(\underline{r},v,t)$  which would be estimated as 4.0 (position points) (dimensions) (energy states) (time intervals) which can easily become a very large number when considering 40 or 50 points in a physical system, for cases in two or three dimensions.

The crux of the matter is that the one dimensional model will allow the utility of this approach to be demonstrated easily with relatively simple, easily compared Fourier series approximations in sine and cosine terms, rather than the functional forms required for solution in three dimensional cylindrical or spherical coordinate systems for sources not located at the principal axis symmetry.

### C. SOURCE FUNCTION APPROXIMATIONS

Within the bounds of the one dimensional spatial model, the five source functions can be described by comparatively simple Fourier series approximations:

$$S(x, v_i, 0) = \sum_n (A_n \cos (k_n x) + C_n \sin (k_n x))$$

for a representation of the source geometry within the confines of the slab.

The following labeling conventions are used in defining the terms of the Fourier series representation:

$S_0$  source strength-total neutrons in the pulse.

$x_0$  half width of the slab.



$x_1$  half width of the finite volume source.

$x_2$  mid-point of the source for the off axis conditions

$L$  diffusion length for the source velocity neutrons.

$\lambda$  the extrapolation distance for the source neutrons.  
The source function approximations are subject to the conditions:

Square Pulse

On-Axis  $S(x, v_j, 0) = \begin{cases} 0 & x < -x_1 \\ S_0 & -x_1 \leq x \leq x_1 \\ 0 & x > x_1 \end{cases}$

Off-Axis  $S(x, v_j, 0) = \begin{cases} 0 & x < x_2 - x_1 \\ S_0 & x_2 + x_1 \leq x \leq x_2 - x_1 \\ 0 & x > x_2 + x_1 \end{cases}$

First Collision Source

On-Axis  $S(x, v_j, 0) = \begin{cases} S_0 e^{+ax} & x \leq 0 \\ S_0 e^{-ax} & x \geq 0 \\ 0 & x > x_1 \end{cases}$

Off-Axis  $S(x, v_j, 0) = \begin{cases} S_0 e^{-a|x-x_2|} & -x_0 \leq x \leq x_0 \\ 0 & x < -x_1 \end{cases}$

Exterior Source

$$S(x, v_j, 0) = \begin{cases} mx + b & -x_0 \leq x \leq +x_0 \\ m = \frac{S_0}{2x_0^2} & b = \frac{S_0}{2x_0} \end{cases}$$



Source Type	Series Representation	Fourier Coefficients
(1) Square Pulse on axis n = positive odd integers	$A_n \cos k_n x$	$k_n = \frac{n\pi}{2(x_0 + \lambda)}$ $A_n = \frac{S_0 \sin (k_n x_1)}{k_n x_0}$
(2) Square Pulse off axis n = positive odd integers	$A_n \cos k_n (x - x_2)$	$k_n = \frac{n\pi}{2(x_0 + x_1 + \lambda)}$ $A_n = \frac{S_0 \sin k_n x_1}{2n x_1}$
(3) First Collision Source on axis n = positive odd integers	$A_n \cos k_n x$	$k_n = \frac{n\pi}{2(x_0 + \lambda)}$ $A_n = \frac{S_0 (-1)^{\frac{n-1}{2}}}{n(x_0/nL)^2 (1 - \exp(-x_0/L))}$
(4) First Collision off axis source n = positive integers	$A_n \cos k_n x + C_n \sin k_n x$	$k_n = \frac{n\pi}{2(x_0 + x_2 + \lambda)}$ $A_n = \frac{S_0 (\cos k_n x_1 - k_n L \sin k_n x_1 - (-1)^{n+1} \sin(x_1/L))}{L(1 + (k_n L)^2)(1 - \exp(-(x_0 + x_1)/L))}$

TABLE I. Fourier Expansion Coefficients.



TABLE I. - Continued

$$C_n = \frac{S_o (\operatorname{Sink}_n x_1 - k_n L \operatorname{Cosk}_n x_1 - (-1)^n \sinh(x_1/L))}{L(1 + (k_n L)^2)(1 - \exp(-\frac{(x_o + x_1)/L}{n}))}$$

(5)      Exterior Ramp  
 Function Source  
 n = positive integers

$$A_n \operatorname{Sink}_n x$$

$$k_n = \frac{n\pi}{4(x_o + \lambda)}$$

$$A_n = \frac{(-1)^{n+1} 2 S_o}{n\pi x_o}$$



#### D. ENERGY-VELOCITY-LETHARGY

The slowing down equation can be written in several forms to describe the variation of the neutron energy with time. For computational purposes, it is often convenient to describe the kinetic parameter in terms of the neutron velocity, while at other times, the kinetic energy of the neutron would provide a better description. There are three main terms that can be used interchangeably to describe the neutron during the slowing down process: velocity, energy and lethargy. Any expression in one of these parameters can be transformed to the appropriate form in the others by the relationships:

$$\text{Energy-velocity} \quad E = m_n v^2 / 2$$

$$\text{Energy-lethargy} \quad - \ln(E/E_0) = u$$

$$\text{Velocity-lethargy} \quad - 2 \ln(v/v_0) = u.$$

By these expressions, the parameters,  $E_0$  and  $v_0$ , describe the neutron reference or source condition with the added convenience that as the neutron slows down, its lethargy increases, while for the velocity and energy expressions, the slowing down process obviously results in the neutron going to a state with a smaller value.

Simple relations between the derivatives of these expressions will allow the slowing down equation to be easily revised with the desired terms:

$$dE = m_n v \, dv$$



$$du = \frac{-2 dv}{v} = \frac{-dE}{E} .$$

In terms of a neutron density function,  $N_1(v)$  and the differential relations, all three forms of the neutron density are:

$$N(v) dv = N(E) dE \quad N(v) = m_n v N(E) = (2m_n E)^{\frac{1}{2}} N(E)$$

$$N(E) dE = N(u) du \quad -E N(E) = N(u)$$

$$N(v) du = N(u) du \quad N(u) = \frac{-v}{2} N(v) .$$



#### IV. APPLICATIONS

##### A. NEUTRON DENSITY

The neutron density can be determined in the numerical model by the infinite series expression:

$$N(\underline{r}, v, t) = \sum_{n=1}^N R_n(\underline{r}) F_n(v, t)$$

such that the series representation converges to the analytical value as  $N$  increases to infinity. In the numerical method for this calculation, the series representation will be determined as a finite series, truncated to the first 3 to 6 terms. In this form of the approximation, the error introduced is principally that of under estimating the true value if an even number of terms are used, while slightly over estimating with an odd number of terms.

The truncation error of the finite series approximation can be estimated from the initial conditions of the system at time equal zero as the neutron density at that time must equal the initial source strength:

$$N(\underline{r}, v, 0) = S(\underline{r}, v, 0).$$

If the total source strength ( $S_o$ ) is obtained by integrating over all velocities and the source volume, the source strength can be given by the series approximation:

$$S_o = \int_v \int_r S(\underline{r}, v, 0) dv dr = \int_v \int_r \sum_n R_n(\underline{r}) F_n(v, 0) dv dr.$$



With the use of a truncated series approximation, the error introduced can be calculated as:

$$\text{Error} = S_0 - \iint \sum_{n=1}^N R_n(\underline{r}) F_n(v, 0) d\underline{r} dv.$$

This normalization error is used to obtain a first order adjustment factor,  $\Delta_0$ , by the simple approximation:

$$\Delta_0 = \frac{\text{Error}}{S_0}$$

and a normalization adjustment is determined as:

$$\text{Source Normalization} = (1 + \Delta_0) S_0$$

such that the truncated series will give a numerical value equal to the initially defined parameter- $S_0$ . This now gives the series approximation as:

$$S(\underline{r}, v, 0) = \sum_{n=1}^N (1 + \Delta_0) R_n(\underline{r}) F_n(v, 0)$$

where the artifice of the source strength normalization factor ( $\Delta_0$ ) would be a successingly smaller correction as more terms are included in the finite series representation.

This adjustment of the numerical value of the source strength has no effect on the velocity-time functions,  $F_n(v, t)$ , as this is a constant multiplier included in the spatial dependent expressions.

In the discrete state, discrete time calculations of MOD-5, the probability that a neutron will be within a



specified velocity interval ( $v_i$ ) at the finite time step ( $t_j$ ), transforms the neutron density expression to the form:

$$N(\underline{r}, v_i, t_j) = \sum_n R_n(\underline{r}) F_n(v_i, t_j)$$

and the total neutron population at any fixed time can be found by integration over the spatial coordinate and summation over the velocity intervals:

$$N(t_j) = \int_r \sum_i \sum_n R_n(\underline{r}) F_n(v_i, t_j) d\underline{r}.$$

#### B. NEUTRON FLUX

The neutron flux, a scalar quantity to describe the net flow of neutrons per unit time and unit area, is calculated as the product of the speed of the neutron and the number density of neutrons having that speed:

$$\phi(\underline{r}, v, t) = v N(\underline{r}, v, t)$$

and is approximated in this work as the discrete value:

$$\phi(\underline{r}, v_i, t_j) = v_i N(\underline{r}, v_i, t_j).$$

The total integrated flux can be analytically determined by an integration over all speeds, and in this work, is therefore approximated by a summation over all discrete states of the density parameter  $N(\underline{r}, v_i, t_j)$  to yield the result:

$$\phi(\underline{r}, t_j) = \sum_i v_i N(\underline{r}, v_i, t_j) = \sum_i \sum_n v_i R(\underline{r}) F_n(v_i, t_j).$$



### C. MEAN ENERGY OF NEUTRONS

The spatial variation of the neutron mean energy provides valuable information on their migration and diffusion following the initial burst. This additional insight into the phenomena of "diffusion cooling" considers the general trend of the higher velocity components of the neutron population to leak out the boundaries of the finite system. As such, the remaining neutrons, while still undergoing interactions which remove them or lower their velocity, have their total number decreased by this additional effect.

In the discrete velocity state, discrete time model, the mean energy of the neutrons at a point  $\underline{r}$  in the slab is determined as:

$$\bar{E}(\underline{r}, t_j) = \frac{\sum_i E(v_i) N(\underline{r}, v_i, t_j)}{\sum_i N(\underline{r}, v_i, t_j)} .$$

In comparing the results of this mean energy determination, two basic methods can be used: first, to study the variation in time of the mean energy at a fixed position, and secondly, to consider the ratio of the mean energy at all positions in the system, for all time periods, compared to a single reference position. Within this work, the mean energy is compared at each finite time period to the current value of the mean energy at the center of the slab, while the time variation of the mean energy at the reference point is followed in detail.



In the discrete time-velocity-space model, this is determined as:

$$\frac{\bar{E}(\underline{r}, t_j)}{\bar{E}(\underline{r}_o, t_j)} = \frac{\frac{\sum_i E(v_i) N(\underline{r}', v_i, t_j)}{\sum_i N(\underline{r}', v_i, t_j)}}{\frac{\sum_i E(v_i) N(\underline{r}_o, v_i, t_j)}{\sum_i N(\underline{r}_o, v_i, t_j)}}$$

where  $\underline{r}_o$  is the reference point,  $x=0.0$ , and  $\underline{r}'$  is a position within the slab.

#### D. DETECTOR RESPONSE

Numerous methods [15] are currently utilized to determine the neutron spectra and flux in fast reactor test assemblies, for internal and external measurements. Fast neutron leakage is examined in time of flight tests to evaluate the time-varying spectra and neutron resonance absorption foils to study the internal spatial dependent spectra.

For this analytical model, no single existing device appeared to be compatible to cover the continuous time and energy response as calculated by the discrete state method. An analytical detector is therefore constructed that would respond to the time-velocity variations of the neutron density. For future applications of this model, a one dimensional version of an existing devices characteristics can be introduced.



The analytical detector is modeled as a proton recoil scintillation device capable of measuring the proton recoil of hydrogen-neutron collisions. The detector is considered to be filled with hydrogen at a molecular concentration as  $H_2$  equivalent to an ideal gas at STP., with physical dimensions of a width  $2 X_4$  about a midpoint  $\underline{r}_D = X_3$ .

The detector response (DR) is determined using the ABN-26 group cross section [16] to be:

$$DR(\underline{r}_D, v_i, t_j) = \int \Sigma_s(v_i) v_i N(\underline{r}', v_i, t_j) d\underline{r}'$$

where the integral over the spatial coordinate  $d\underline{r}'$  refers to the total detector volume.

The total response of the detector both for the total integrated flux over all velocities and over all times can be determined in this discrete velocity-time model by numerical integration and summation:

$$DRV(\underline{r}_D, t_j) = \sum_i DR(\underline{r}_D, v_i, t_j)$$

$$DRT(\underline{r}_D, t_j) = \sum_j DRV(\underline{r}_D, t_j)$$

where the final summation over all times from  $t=0$  to  $t=t_j$  provides the neutron fluence at the detector.

#### E. MIL-SIX COMPUTER PROGRAM

A general purpose computer program, MIL-SIX, was written in FORTRAN IV to process the time-energy (lethargy) response



date produced by MOD-5 to determine within this version of the transport approximation, the neutron density function- $N(\underline{r},u,t)$  for the one dimensional infinite slab geometry. Several minor modifications were made in MOD-5 to provide the ability to calculate the time-energy response of a system with multiple harmonic modes and to provide the neutron probability density for each lethargy state at the same time steps. The only form of output system tested involved the use of punch card output from MOD-5, however, provisions have been made to employ nine-track magnetic tape and/or on-line disk storage capabilities (IBM-2314 units). Program MIL-SIX provides the general user the ability to accept any of the three forms output from MOD-5, for up to six (6) harmonic modes of data. The program can be modified easily to adjust for a higher order approximation using more than six modes with a simple change in the size of the storage arrays in the program.

Program MIL-SIX is composed of the following subprograms and routines:

<u>Subroutine</u>	<u>Function</u>
Main	Principal control portion of the program; handle all logical decisions on which routines the program will perform
INCON1	Establishes all principal default parameters for the program; reads all problem definition statements (punch cards) which determine the parameters that the program



	will need to calculate the desired functions; provides the user with instructions and information concerning actions the program will perform for the stated problem run.
READ10	Processes all punch card output from program MOD-5.
READ11	Processes all tape (9-track) from program MOD-5.
ESTIMATE1	Provides guidance information of the execution time to be expected for the running of a particular problem.
MODEL 1	Calculates the non-harmonic terms of the fourier space dependent functional expression for the source geometry specified for the particular problem; estimates the error introduced by the use of a finite number of terms in the fourier infinite series approximation; adjusts the values of the fourier non-harmonic coefficients to obtain a closer approximation to the infinite series representation; and provide a simple space dependent plot of the fourier expansion of the space dependent function on the on-line printer.
PHI1	Principal function sub-program of the program: to calculate the fourier space dependent harmonic terms in the calculation



of the terms of  $R_n(x)$  ( $n$ =the harmonic mode,  $x$ =the space point) for the one dimensional model.

FOUT1( $x_i$ ) Principal determination subprogram to calculate the series of values for the function -  $N(x,u,t) = R_n(x_k)F_n(u_i,t_j)$  for summed response due to all harmonic mode to give the neutron density for the single point -  $x_i$ .

FOUT2 This subprogram will calculate the two principal space dependent relationships:  $\phi(x_k,t_j)$  and  $\bar{E}(x_k,t_j)$  the integrated (total) neutron flux and the mean energy of the neutrons at the point ( $x_k$ ) for a profile of 100 points across the slab for each time steps ( $t_j$ ), and will provide a graphical plot of the functional forms:

$\phi(x_k,t_j)$  vs  $x_k$

$\bar{E}(x_k,t_j) / \bar{E}(0,t_j)$  vs  $x_k$

via a on-line printing of a rough plot of the results or via the Cal-Comp Plotter for each set of discrete time step date ( $t_j$ ); also provides tabulated data on neutron flux, mean energy and mean energy ratio.



FOUT3

Detector response: this routine considers the simple model of a hydrogen recoil detector located inside the slab. This is not a true device, but one which determines the function:

$$\phi(x_k, t_j) = \sum_u v(u) * \Sigma_t * \sum_n R_n(x_k) F_n(u_i, t_j)$$

using the integrated position dependent flux (calculated at 5 points within the assumed detector position); a simple average is performed to obtain the flux at the detector, where the energy dependent microscopic cross section ( $\Sigma_t(u_i)$ ) for STP hydrogen is based on the ABN-26 group cross section tables. Starting with  $\phi(x_k, t_j)$  - flux data, a numerical integration is performed to determine the integrated detector response. Output of the two functional forms of the detector response are given both by the on-line printer, and the cal-comp plotter.

FOUT4

Provides a simple spectral plot of the data input from program MOD-5, to show the distribution  $F_n(u_i, t_j)$  for each harmonic mode versus the lethargy state, showing the movement of the neutron distribution to states of lower energy (higher



lethargy) with time for each mode of harmonic buckling.

FOUT5 Provides a more detailed spectral response plot of the MOD-5 input data, by consideration of the fourier harmonic coefficient  $R_n(x_k)$  for the individual position point -  $x_k$  weighting the harmonic data - to give the response spectra:  $N_n(x_k, u_i, t_j) = R_n(x_i) F_n(u_i, t_j)$  as a single plot for each harmonic mode of input data, and a total spectral response function  $N(x_k, u_i, t_j) = \sum_n R_n(x_k) F_n(u_i, t_j)$ , summed over all modes and plotted versus lethargy state -  $u_k$  on the Cal-Comp plotter.



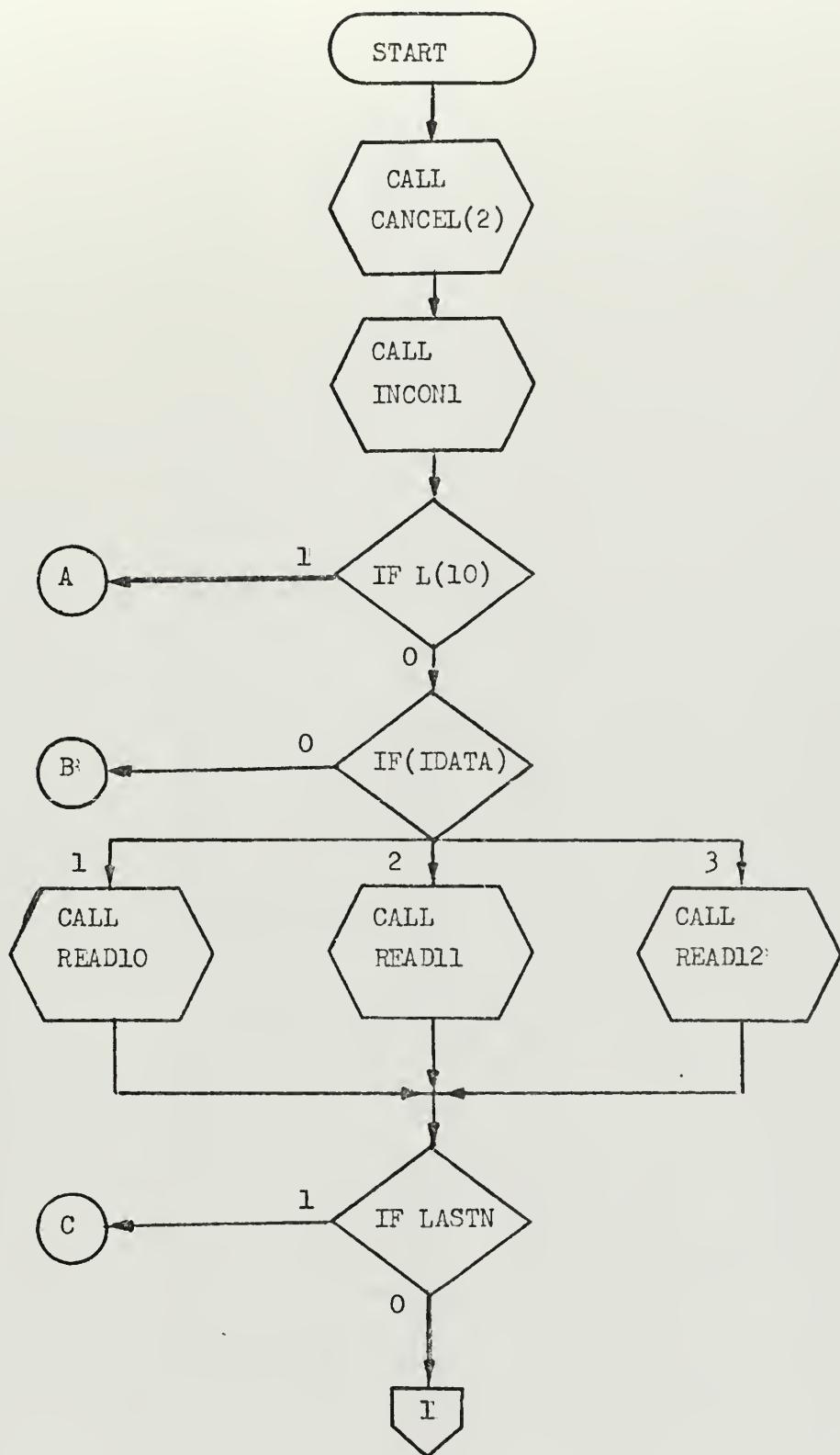
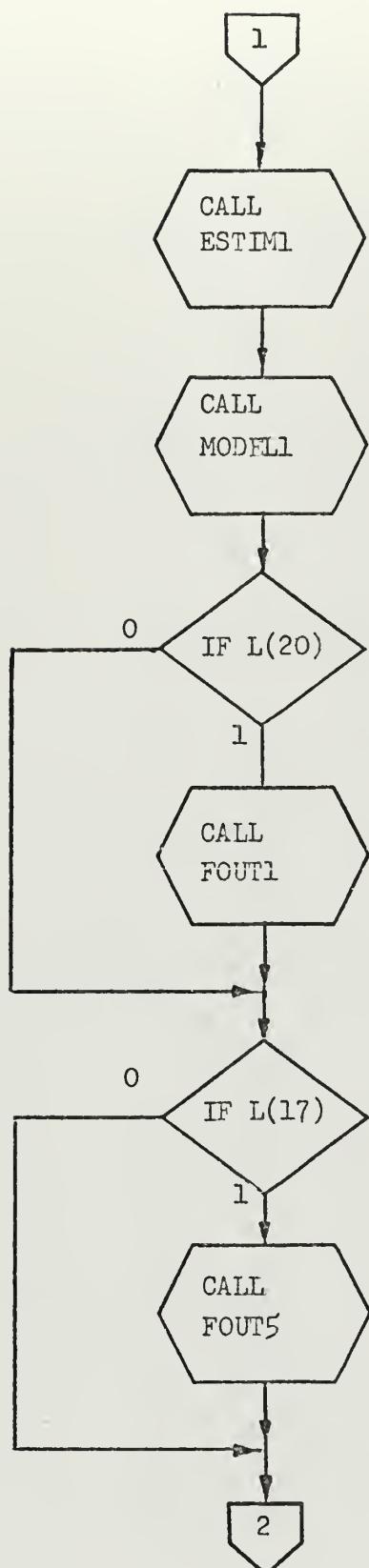
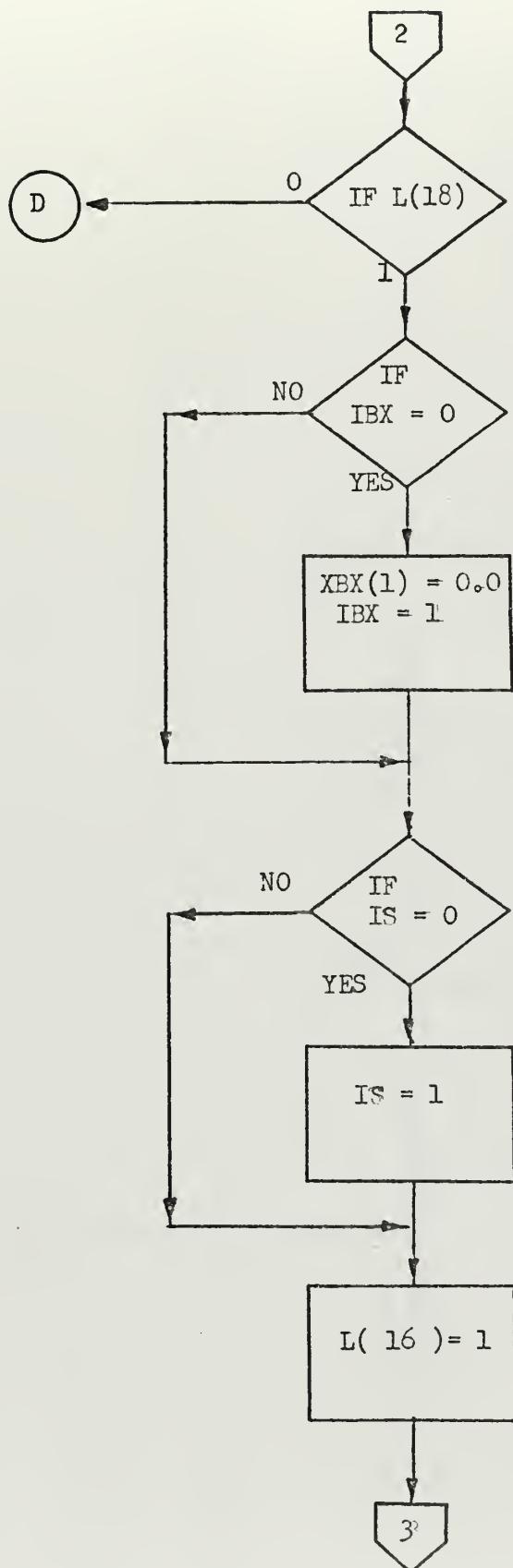


FIGURE 4 FLOW CHART COMPUTER PROGRAM MIL-SIX

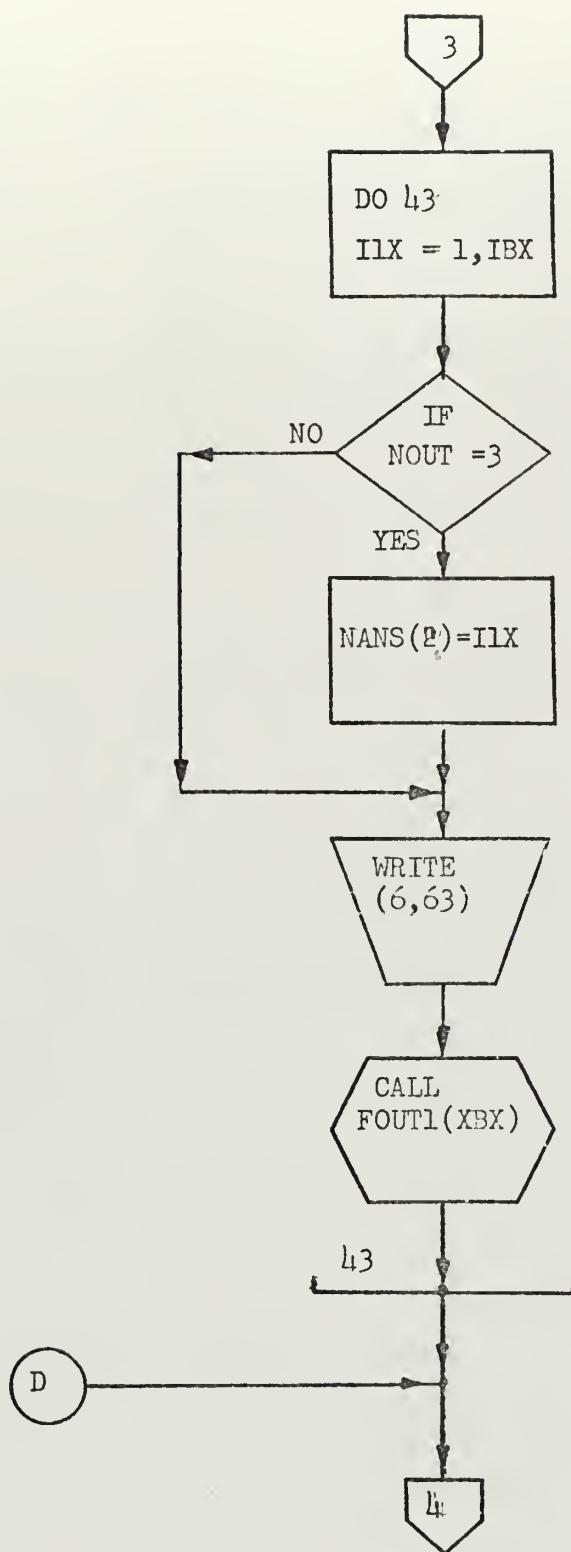




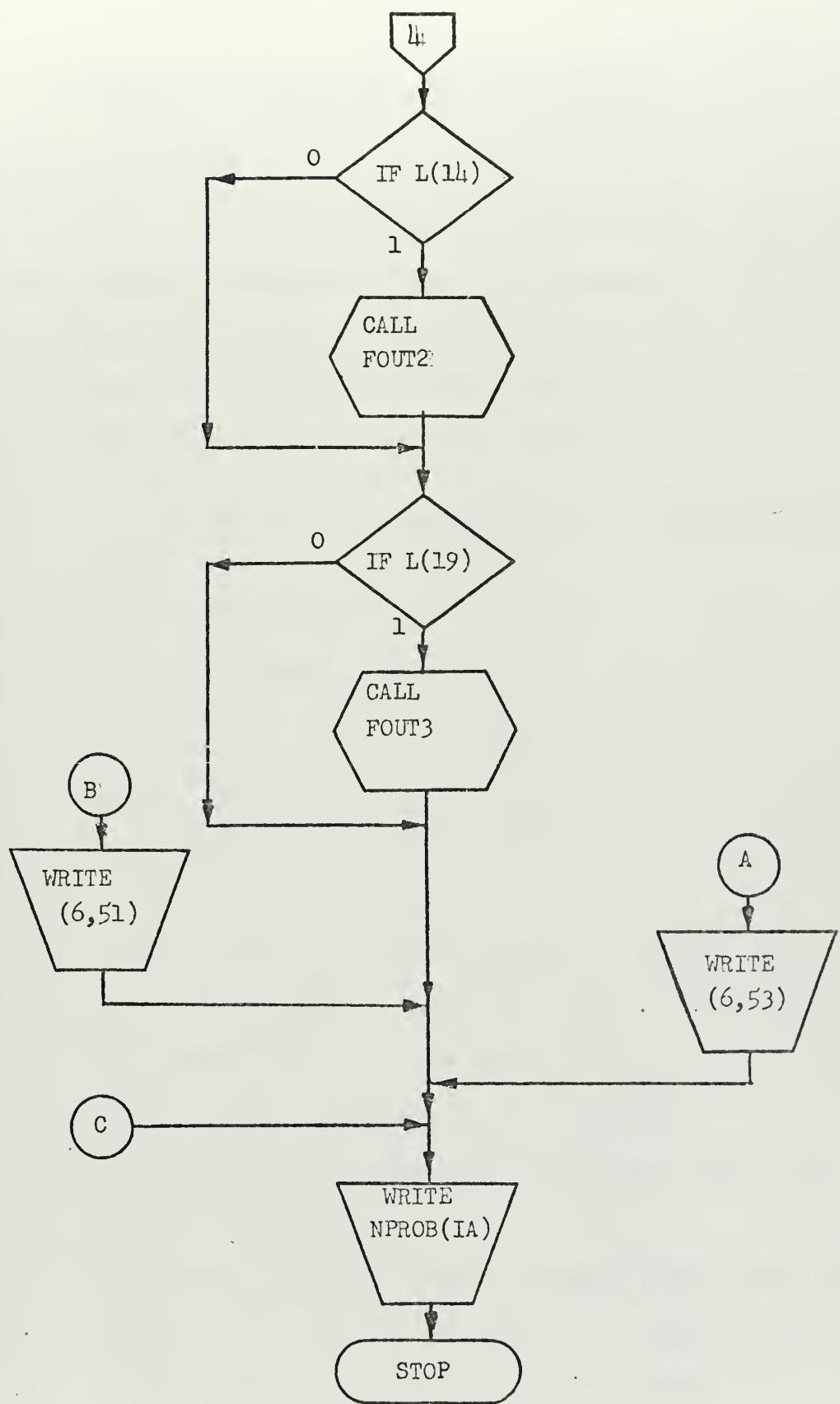














## V. DISCUSSION OF RESULTS AND CONCLUSIONS

### A. PHYSICAL SYSTEMS

The harmonic expansion method was examined for application to multiplying and non-multiplying systems, a beryllium slab and a modified form of the ZPR-3 assembly 6F. The physical parameters of the two systems are given in Table II. The nuclear characteristics were assumed to be:

System	Constituents	Mass Density (gm/cc)	Total Modes
Beryllium Slab	Beryllium	1.84	6
ZPR-3 (6F)	Aluminum	0.848	3
	Iron*	0.965	
	U-235	2.622	
	U-238	3.016	

(\*Iron substituted for stainless steel.)

A pulse of one 2.46 MeV neutron was fixed as a delta function in time to study the decay of each harmonic component of the discrete state velocity-time response  $F_n(v_i, t_j)$ . Two versions of the slowing down process were examined:

- (1) follow each harmonic mode to the same final time ( $t_j$ ) as the fundamental mode,
- (2) follow decay of each harmonic mode to the time step at which the probability density component for that mode was below 0.001 of the neutron remaining in the system.

Both the beryllium and the assembly 6F systems were studied to determine the variations in the neutron spectra



ZPR-3 (6F)Beryllium Slab

Slab Half Width ( $x_0$ )	17.50	25.00
Source Half Width ( $x_1$ )	2.00	2.50
Diffusion Length	9.79	7.085
Mid Point or Off-Axis Source ( $x_2$ )	12.50	12.50
Detector Mid Point ( $x_3$ )	10.00	
Detector Half Width ( $x_4$ )	1.00	12.50
Fundamental		1.00
Buckling Mode ( $B_0^{-2}$ )	.00704	.00093

All measurements in centimeters

TABLE II. Physical Parameters of Test Systems.



resulting from the inclusions of the higher harmonic modes. The persistance of the higher modes was observed in the early time history, but provided a negligible contribution toward the end of the decay time.

## B. DATA PRODUCTION

The neutron probability density data for the lethargy-time response was calculated by MOD-5 for all harmonic modes using two limits on the systems considered. MOD-5 normally selects those discrete time steps to provide the probability density data under its internal selection criteria as though all problems were a fundamental mode type calculation. This freedom in selecting the output time steps was over ridden by modification of the existing control sequency in the basic MOD-5 program.

The fundamental mode calculation was done for a system to obtain the fixed time intervals and output time steps for the higher harmonic mode calculations. With the "locking in" of the data output time steps, the higher harmonic modes were run to the same final discrete time as the fundamental mode. Unfortunately, as the higher modes decay more rapidly, this resulted in the abnormal termination of several computer runs when the probability density dropped to values smaller than the IBM-360/67 is capable of processing, i.e., below  $10^{-75}$ . As a result, all final computer runs were set to terminate when the probability density component for a mode dropped below an arbitrary limit of 0.001.



The time decay of the harmonic components is mainly affected by the leakage of neutrons from the system which is dependent on the buckling factor,  $DB_n^2$ , which appears in the time dependent diffusion equation as a term:

$$vD(v)B^2N(x, v, t).$$

A summary of the MOD-5 calculations is given in Table III for the two basic systems where all calculations were carried to the point where the probability that a neutron had leaked from the system, been absorbed or had dropped below the lowest energy was .999. In all cases, the higher harmonics were still providing a small contribution to the total population at times approaching 300ns, but the major contribution was due mainly to the fundamental.

#### C. ANALYSIS OF EXPANSION METHODS

The adequacy of the truncated Fourier series to represent the various source geometry conditions was evaluated by the comparison of the source strength parameter which was determined via the form:

$$\Delta_o = 1.0 - \int_{-x_o}^{x_o} \sum_{n=1}^N (A_n \cos knx + C_n \sin knx) dx$$

$$\Delta_o = 1.0 - \sum_{n=1}^N (-) \frac{An^2}{kn} \sin knx_o - \frac{Cn}{kn} \cos knx_o.$$

The uncorrected source strength error ( $\Delta_o$ ) was evaluated for the first 100 terms of the series expansion; which stabilized to almost constant-values after the first 30 to 40 terms.



System	Mode	Time (ns)	Final Neutron Probability Components (Percentage)			
			Slowing to Bottom of Spectrum**	Leakage	Capture non-fission	Fission
Beryllium	1	16620	83.9	7.4	7.9	-
	2	9834	44.7	48.0	6.8	-
	3	4983	14.2	80.3	5.4	-
	4	1952	3.1	92.6	3.4	-
	5	1315	.01	95.6	3.4	-
	6	91	.002	96.1	2.9	-
ZPR-3 (6F)	1	311	.95	51.1	8.7	39.2
	2	311	.001	89.2	1.2	9.5
	3	311	.0	88.6	1.2	10.2

\*\* Bottom of Spectrum is defined as:

Beryllium Slab below 2.31 eV

Assembly 6F below .4 eV

TABLE III. Summary of MOD-5 Computer Runs.



Uncorrected Source Error ( $\Delta_o$ ) 100 Term

Source Type	ZPR-3 (6F)	Beryllium Slab
1- Square Pulse	0.0652	0.0265
3- Exponential	0.0231	0.190
5- Ramp Function	0.0090	0.0090

The Fourier series expansions are defined such that with an infinite number of terms, the error should go to zero. However, calculations show that for the source function approximations, this error goes to some non-zero value. It should not be any surprise that these source terms have a non-zero error, since the Fourier approximation is based on the interval,  $x_o + \delta$ , where the distance " $\delta$ " represents the extrapolation distance correction for the boundary condition of forcing the neutron density to zero at that point.

In the source function testing, the physical limits of the slab are within a width  $2x_o$ , rather than boundaries of  $2(x_o + \delta)$ . This additional value of twice the extrapolation distance is responsible for apparent constant error. When the integration is done within the bounds of the slab, a loss occurs at the end points in the intervals

$$-(x_o + \delta) \leq x' \leq -x_o \quad \text{and} \quad +x_o \geq x \geq (x_o + \delta).$$

As the ratio  $\delta/x_o$  is greater than zero for all real (non-infinite systems), this apparent difficulty in the approximation will be almost negligible for systems with the slab width much greater than the extrapolation distance ( $\delta/x_o \ll 1.0$ ).

Each of the three main source types (the square pulse, exponential and ramp functions) were studied to the limit of



a 100 mode harmonic expansion for the spatial dependent functions. The value of the delta correction ( $\Delta_0$ ) are shown in Figures 5 and 6 for both systems. While all test calculations were done for a six mode expansion, the effect of the truncation to less than an infinite series approaches an asymptotic limit in each source type, the limits on the value of  $\Delta_0$  are a measure of the  $\delta/x_0$  ratio. In both of the systems studied, it can easily be seen that inclusion of more than 10 harmonic modes would not significantly change the spatial dependent values.

#### D. COMPARISON OF SOURCE TYPES

The development and testing of the time-space-energy-dependence for five source geometries was a task more formidable than originally anticipated. Numerous programming difficulties arose toward the end of the project for the two cases of the off-axis sources, which resulted in their not being included in the final calculations. Both of these sources require modification of the spatial harmonic functions than those proposed previously.

The three source conditions that are included in the detailed calculations were the symmetrical square pulse and the first collision source centered about the mid-point of the slab and the exterior wide beam or ramp function source.

The existing version of the computer program, MIL-6, demonstrated the effectiveness of the three main source conditions, but the off-axis square pulse and first collision geometry must be revised and retested prior to further analysis.



FIGURE 5  
SOURCE FUNCTION ERROR  
BERYLLIUM SLAB

● - SOURCE = 1  
X - SOURCE = 3  
+ - SOURCE = 5

TOTAL MODES - HARMONIC EXPANSION

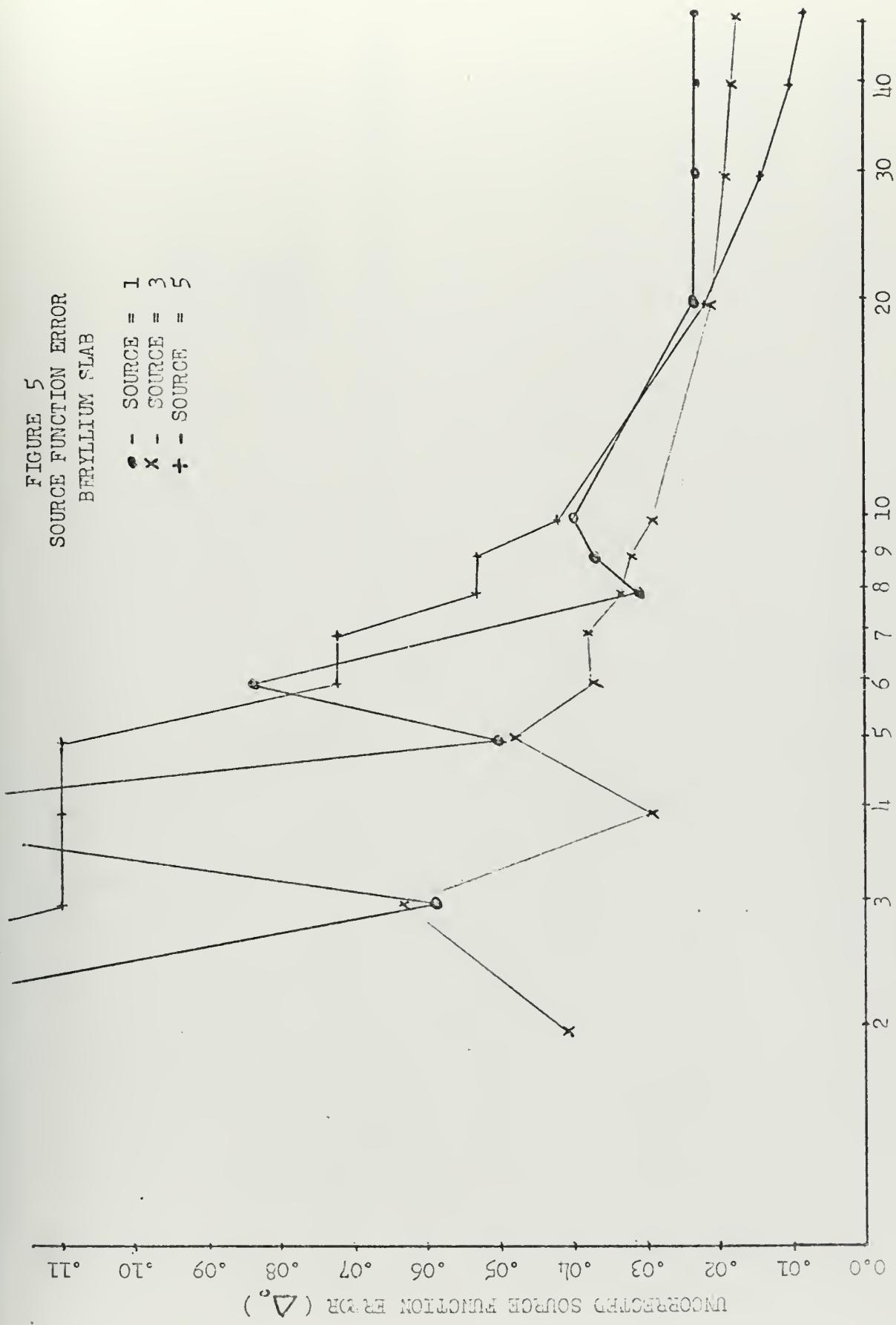
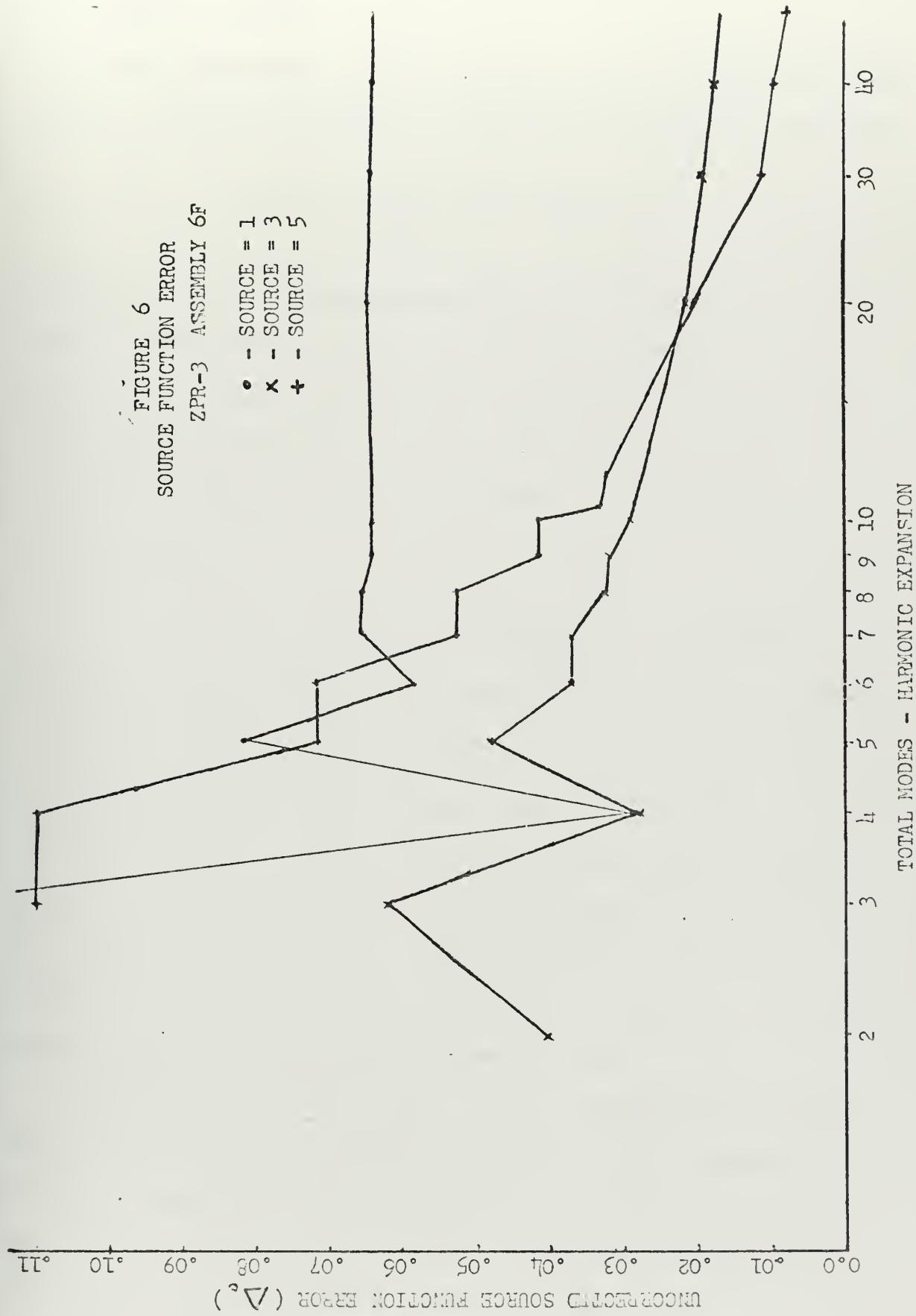




FIGURE 6  
SOURCE FUNCTION ERROR  
ZPR-3 ASSEMBLY 6F





## E. NEUTRON SPECTRA

The primary goal of this project was to demonstrate the spatial dependent response of the neutron spectral function:

$$N(x, v, t) = \sum_n R_n(x) F_n(v, t).$$

The contribution of each harmonic mode to the neutron spectra is determined by a Fourier space harmonic term,  $R_n(x)$ , and the MOD-5 time-lethargy probability density for the  $n$ -th time mode,  $F_n(v, t)$ . For each harmonic mode, the probability density function is subject to the basic condition:

$$F_n(v, t) = 1.0 - (\text{Absorption} + \text{Leakage})$$

and the initial neutron source spatial distribution provides the correct source contribution.

## F. NEUTRON DENSITY AND NEUTRON FLUX

The results of neutron density and neutron flux calculations are presented at a limited number of points to be representative of the over all processes, however, inclusion of all points in the slab would result in an inordinate number of graphs and tables of data.

The neutron density,  $N(x_k, v_i, t_j)$ , was plotted only for the point  $x_k = 0.0$  and a complete spectral plot is given in Figures 13 and 16, to be representative of the information possible to determine. All other calculations made by the program MIL-6 required the evaluation of the density at all points in the slab, and similar response curves were prepared.



The total neutron flux,

$$\phi(x_k, t_j) = \sum_i v_i N(x_k, v_i, t_j),$$

is presented for one location half way between the center and the surface of the slab, the same positions selected for the detector response calculations. Each plot is for a single system, either the Beryllium or assembly 6F, follows the complete time decay of the total flux from one nanosecond to the time at which the neutron population had decayed to 0.1% of its initial value. In both cases, the first build up of the flux indicates the initial stages of the neutron motion within the slab, to the point at 20ns at which the leakage from the free surface becomes the dominant effect.

When the leakage terms are considered, the appearance of the decay curve now looks almost identical with the pseudo decay modes previously described. When looking at the neutron density curves, it can easily be noticed that it required several nanoseconds for the initial 2.46 MeV burst to develop into a smooth distribution.

#### G. MEAN ENERGY AND MEAN ENERGY RATIO

A convenient method of comparing the time and spatial effects on the neutron density is by comparing the average or mean energy of the neutrons at all points in the slab at the same time. A complete history of the neutron population can be concisely presented by observing the absolute variation of the mean energy at a single reference point in



the slab and then comparing the relative changes in the entire slab in terms of the reference point for each discrete time step.

The basic reference point for all comparisons is the center of the slab,  $x=0.0$ , for all geometries and systems. The decay of the mean energy with time is shown for both beryllium (Figure 19) and assembly 6F (Figure 20). Starting with the two reference cases, the spatial variation from the center to the surface of the slab starting at two nanoseconds after the pulse to the final decay of the fundamental mode, (beryllium - Figures 21 to 30--assembly 6F Figures 31 to 36).

After the initial pulse, a wave-like shape is observed to develop across the slab in the direction of the free surface, as one would expect the higher energy neutrons to migrate out of the system fastest. This is shown most vividly in Figures 24 to 27.

Near the end of the decay period, the last remaining neutrons in the system have almost a constant ratio of one, as the system tends toward an equilibrium condition. The leakage at the surface allows the final stages to show this condition (Figures 30 and 36).

#### H. DETECTOR RESPONSE

This computational routine did not work as well as intended. Minor computational and programming errors were encountered very late in the project and have not been completely resolved. These will require revision and further testing as follow on work.



Comparisons of the flux calculations and the mean energy determinations referred to in Sections F and G were compared to the "proton recoil" detector (Figures 37 and 38) which follow the same basic trends for the time dependent response, but the numerical differences in the time integrated response will require improvement of the space integral and time summation numerical methods.



## VI. RECOMMENDATIONS

### A. FOLLOW ON WORK

The program MIL-6 was prepared to be a general computational tool using the harmonic mode expansion method for the one dimensional systems. Principal follow on work is recommended to correct the minor computational difficulties encountered, write and test data transfer routines for both MOD-5 and MIL-6 and review and edit the programming for a more optimized execution.

Principal areas to have future work are:

(1) re-evaluate the Fourier space functions for the two off-axis source geometries,

(2) revise and test the numerical integration techniques of the detector response program for the summation over time, and consider methods to improve the space integration over the detector volume,

(3) consider a detector response function that would be more realistic than the STP hydrogen case now included in the program,

(4) complete and test the input/output options using nine-track magnetic tape and the IBM-2314 magnetic disc systems for data transfer between the programs MOD-5 and MIL-6,

(5) revise the program MIL-6 to handle more than six harmonic modes by use of a smaller state structure of 70



to 100 states rather than the 150 states the program is currently written to process,

(6) provide more realistic source geometries for the one dimensional case than the five simple ones currently in the program,

(7) review the program for reduction in the core storage space requirements and to improve the speed of execution through program optimization.

#### B. IMPROVEMENTS TO PROGRAM MIL-6

The original program MIL-6 was written to deal with monoenergetic neutron sources and delta functions in time. This may require significant modification of the computational methods to considering the fission spectrum sources and sources that have a finite width time duration. These situations had not been tested in the MIL-6 work and would provide a more realistic model for fast neutron experiments.

A second item that appears worthy of consideration in the future will assist in modeling the wave-like motion of the neutron pulse from the finite volume square pulse and the exterior wide beam ramp function. This could be considered as a limiting correction factor to couple the time-space and energy dependence during the initial response to a neutron pulse.

#### C. THREE DIMENSIONAL SYSTEMS

The harmonic expansion method has been demonstrated to work for the one dimensional case, and definite action to



begin study of three dimensional systems is warranted. A recommended sequence of priorities would be to expand the MIL-6 work to three dimensions in the cartesian coordinates first, cylindrical geometry second and spherical geometry coordinate systems last.



## APPENDIX A: NUMERICAL RESULTS

### A. SPECTRAL RESPONSE

#### 1. Beryllium Slab

a. Neutron density harmonic component  $N_n(x, v_i, t_j)$  for harmonic modes 1 to 6 at point  $x=0.0$ . (Figures 7 to 12.)

b. Neutron density:  $N(x, v_i, t_j)$ : lethargy - time neutron population density versus lethargy for a 6 mode expansion at point  $x=0.0$ . (Figure 13.)

#### 2. ZPR-3 Assembly 6F

a. Neutron density harmonic component  $N_n(x, v_i, t_j)$  for harmonic modes - 1 and 2 at point  $x=0.0$ . (Figures 14 and 15.)

b. Neutron density  $N(x, v_i, t_j)$  - lethargy - time neutron population density versus lethargy - for a sum of 3 harmonic modes at  $x=0.0$ . (Figure 16.)

### B. MEAN ENERGY AND MEAN ENERGY RATIO

#### 1. Beryllium

a. Mean energy verus time at  $x=0.0$ . (Figure 19.)

b. Plot of mean energy ratio versus position for 8 time periods. (Figures 21 to 30)

#### 2. ZPR-3 (6F)

a. Mean energy versus time at  $x=0.0$ . (Figure 20.)

b. Plot of mean energy ratio versus position for 6 time periods. (Figures 31 to 36.)



C. DETECTOR RESPONSE

1. Beryllium - Figure 37.
2. ZPR-3 - Figure 38.



FIGURE 7  
 NEUTRON DENSITY  
 BERYLLIUM SLAB  
 FUNDAMENTAL MODE  
 $X = 0$

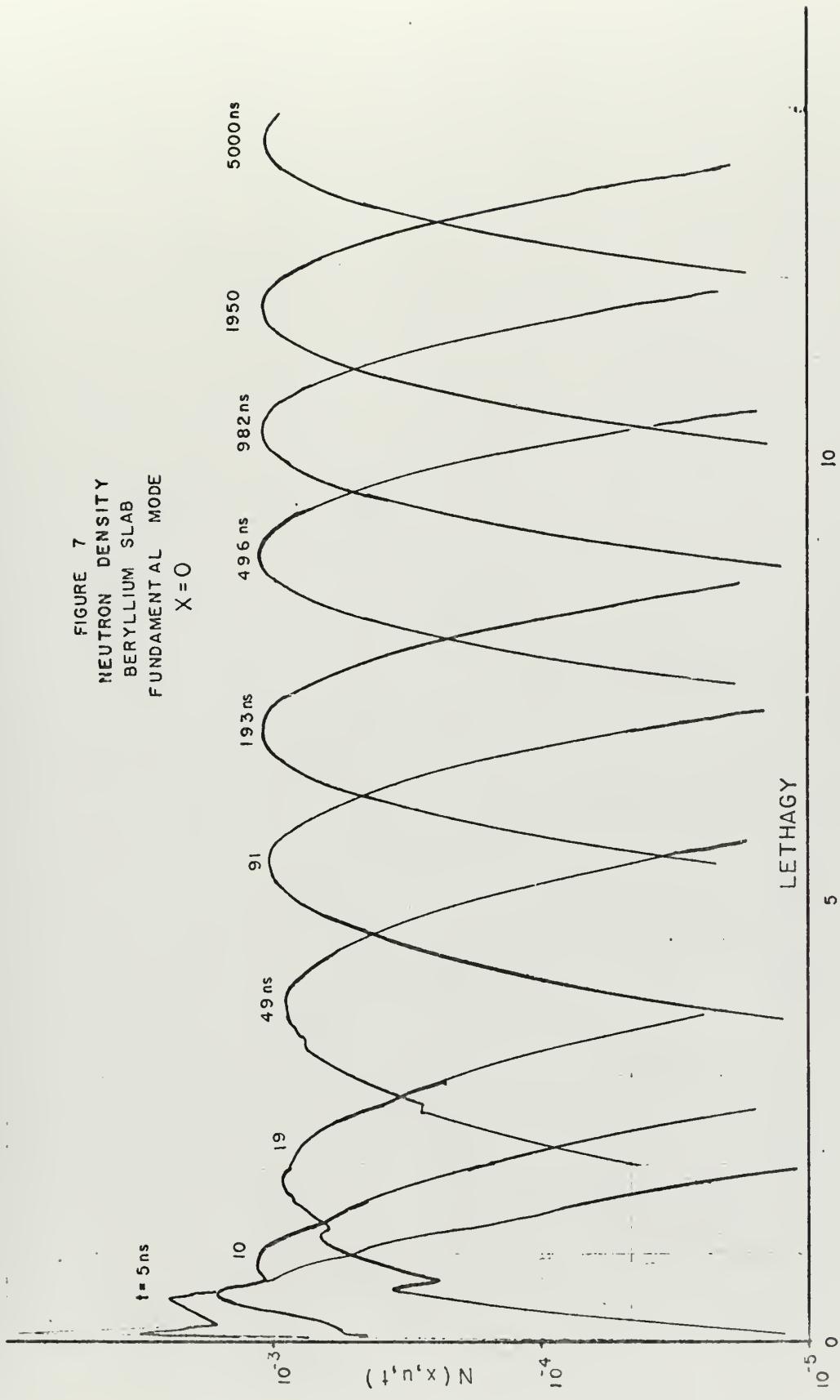




FIGURE 8  
NEUTRON DENSITY  
BERYLliUM SLAB  
FIRST HARMONIC MODE  
 $X=0$

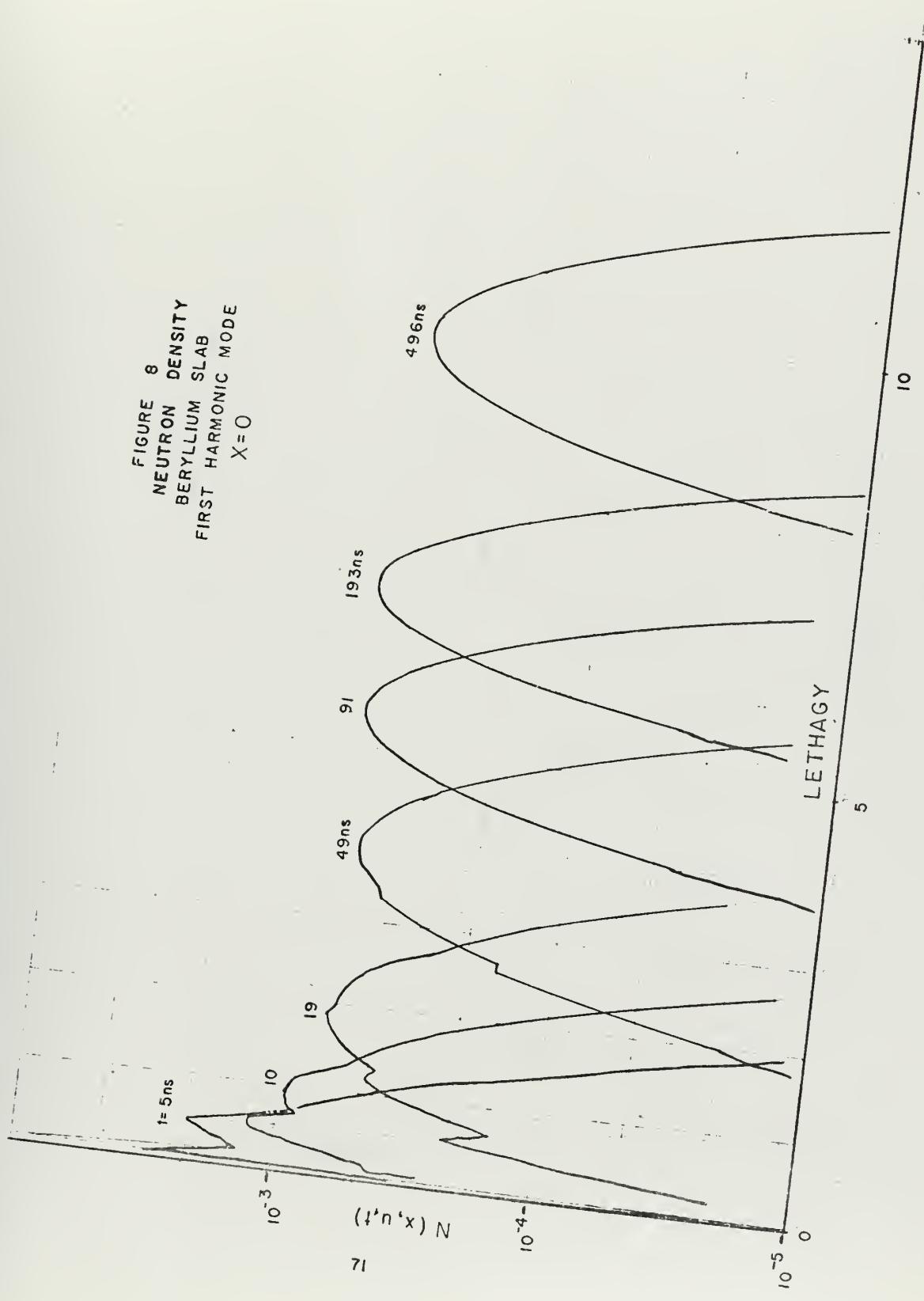




FIGURE 9  
NEUTRON DENSITY  
BERYLLIUM SLAB  
SECOND HARMONIC MODE  
 $X = 0$

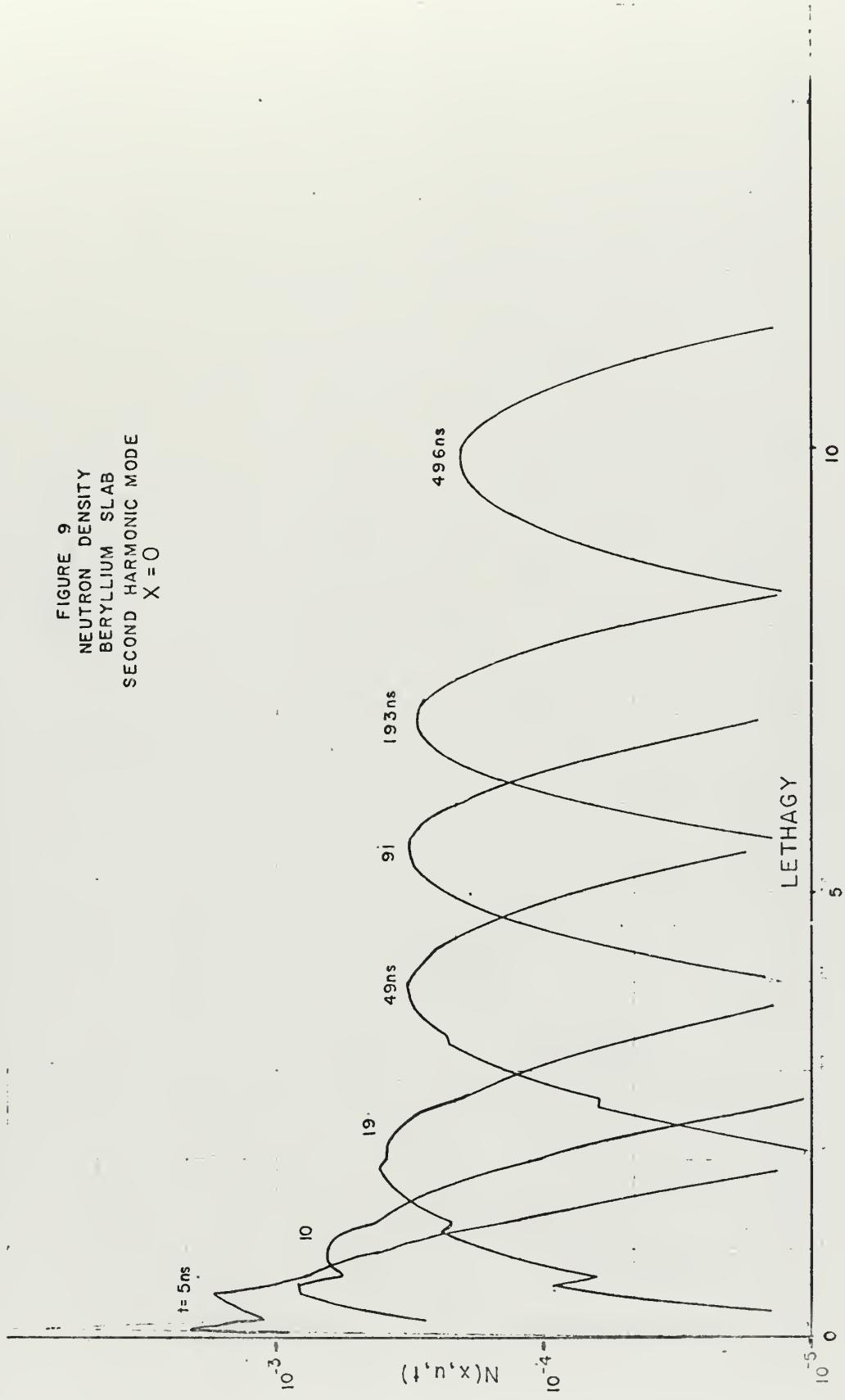




FIGURE 10  
NEUTRON DENSITY  
BERYLliUM SLAB  
THIRD HARMONIC MODE  
 $X = 0$

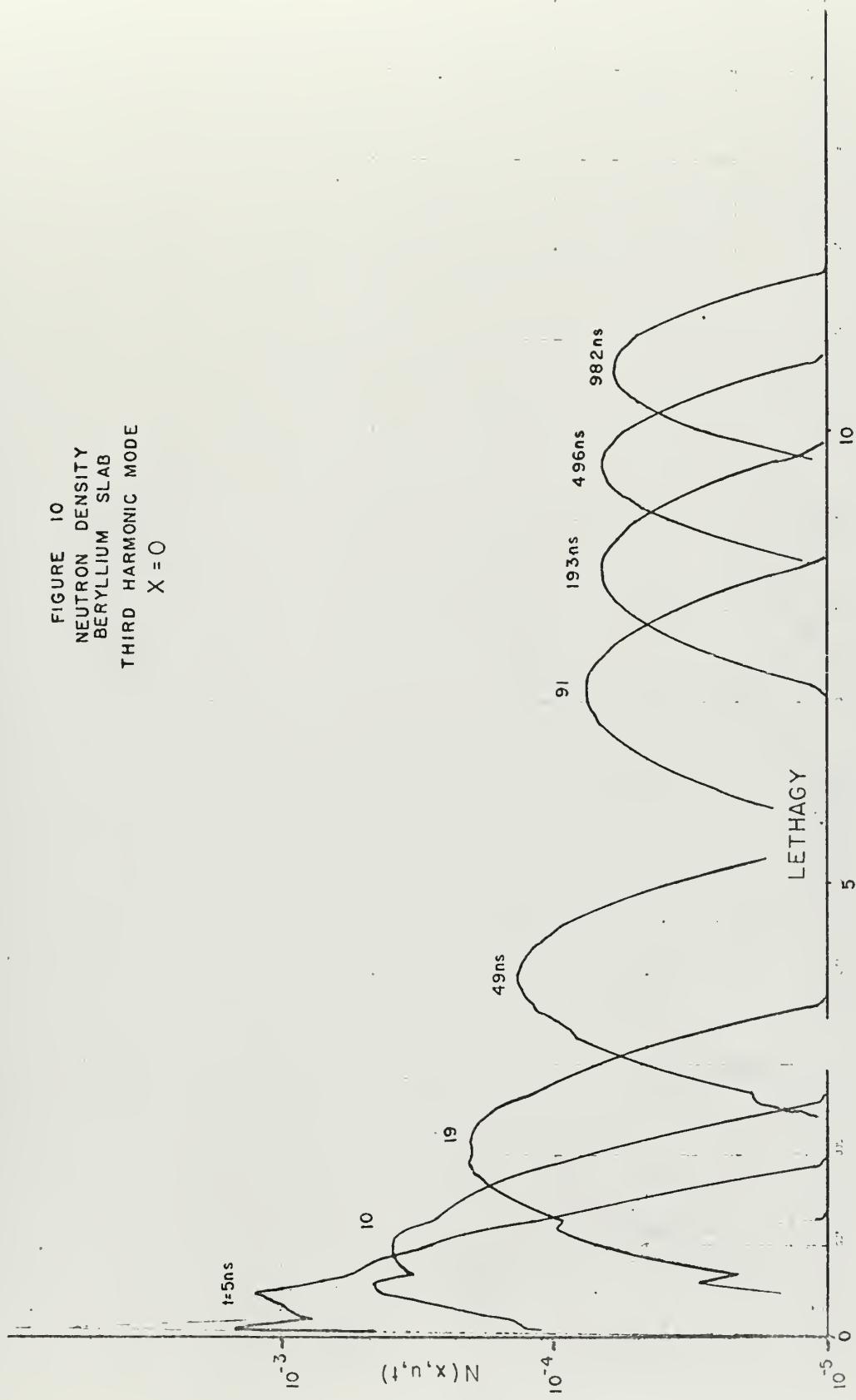




FIGURE 11  
NEUTRON DENSITY  
BERYLLIUM SLAB  
FOURTH HARMONIC MODE  
 $X = 0$

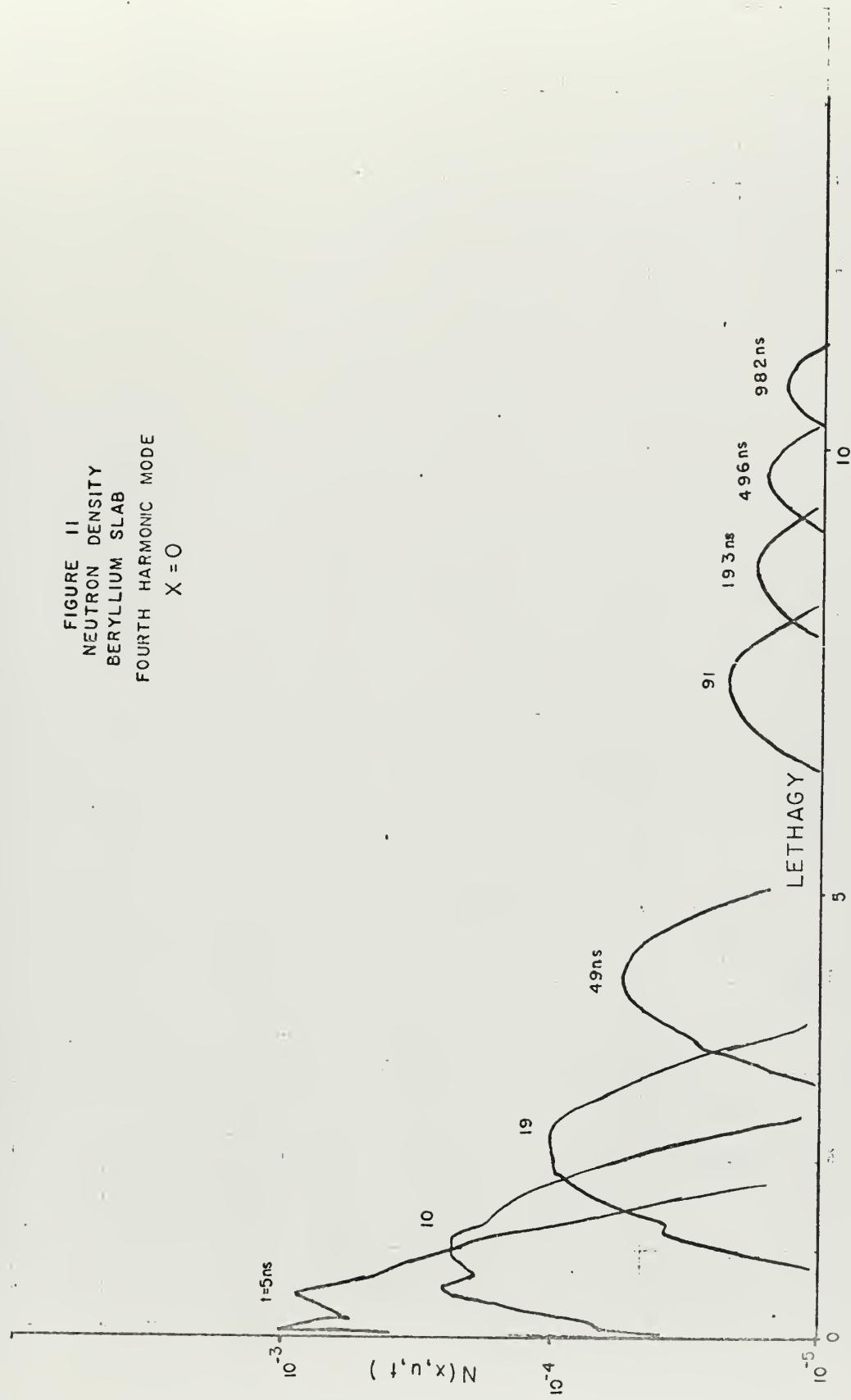




FIGURE 12  
NEUTRON DENSITY  
BERYLLOM SLAB  
FIFTH HARMONIC MODE  
 $X = 0$

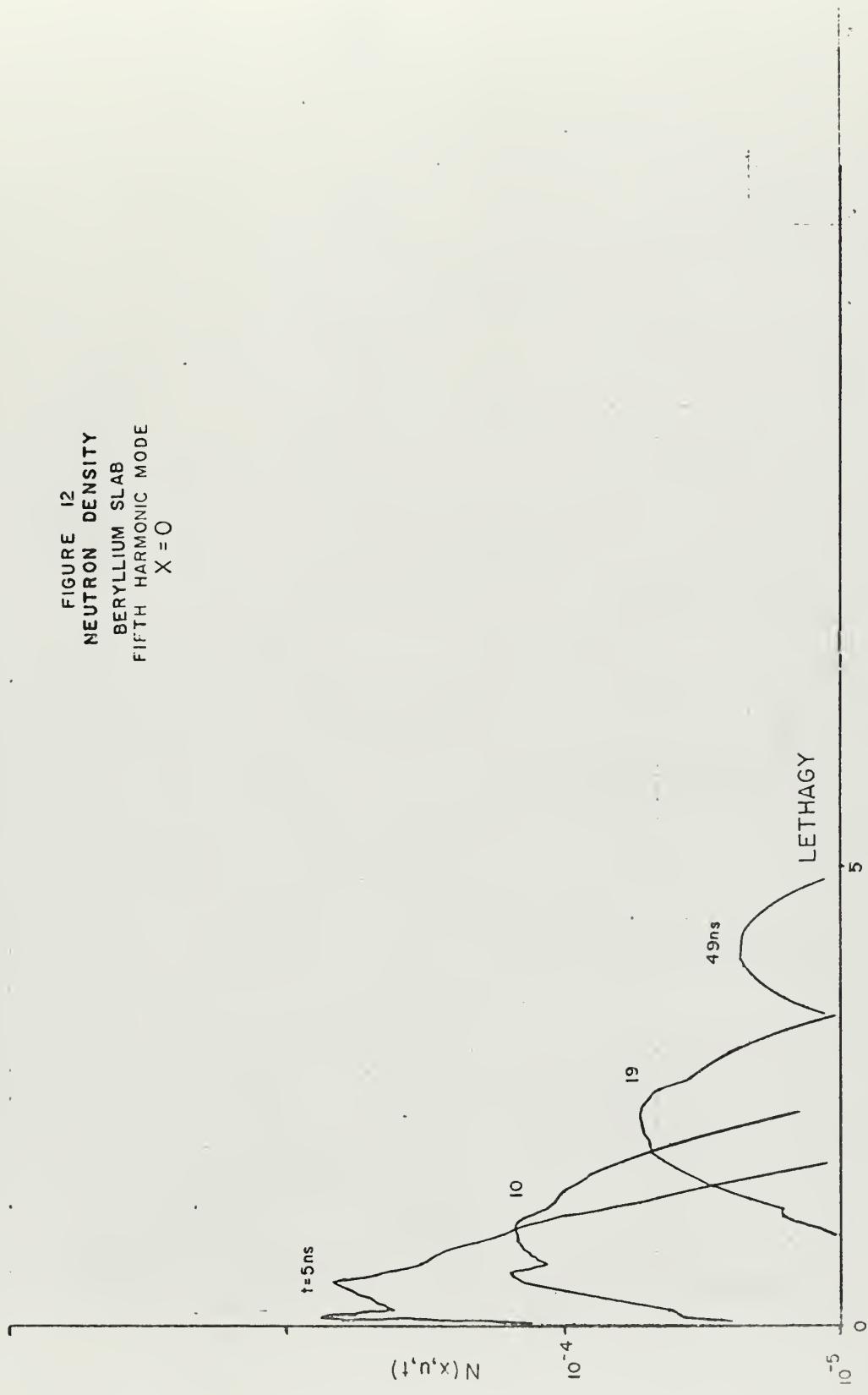




FIGURE 13  
NEUTRON DENSITY  
BERYLliUM SLAB  
SUM OF ALL MODES  
 $X=0$

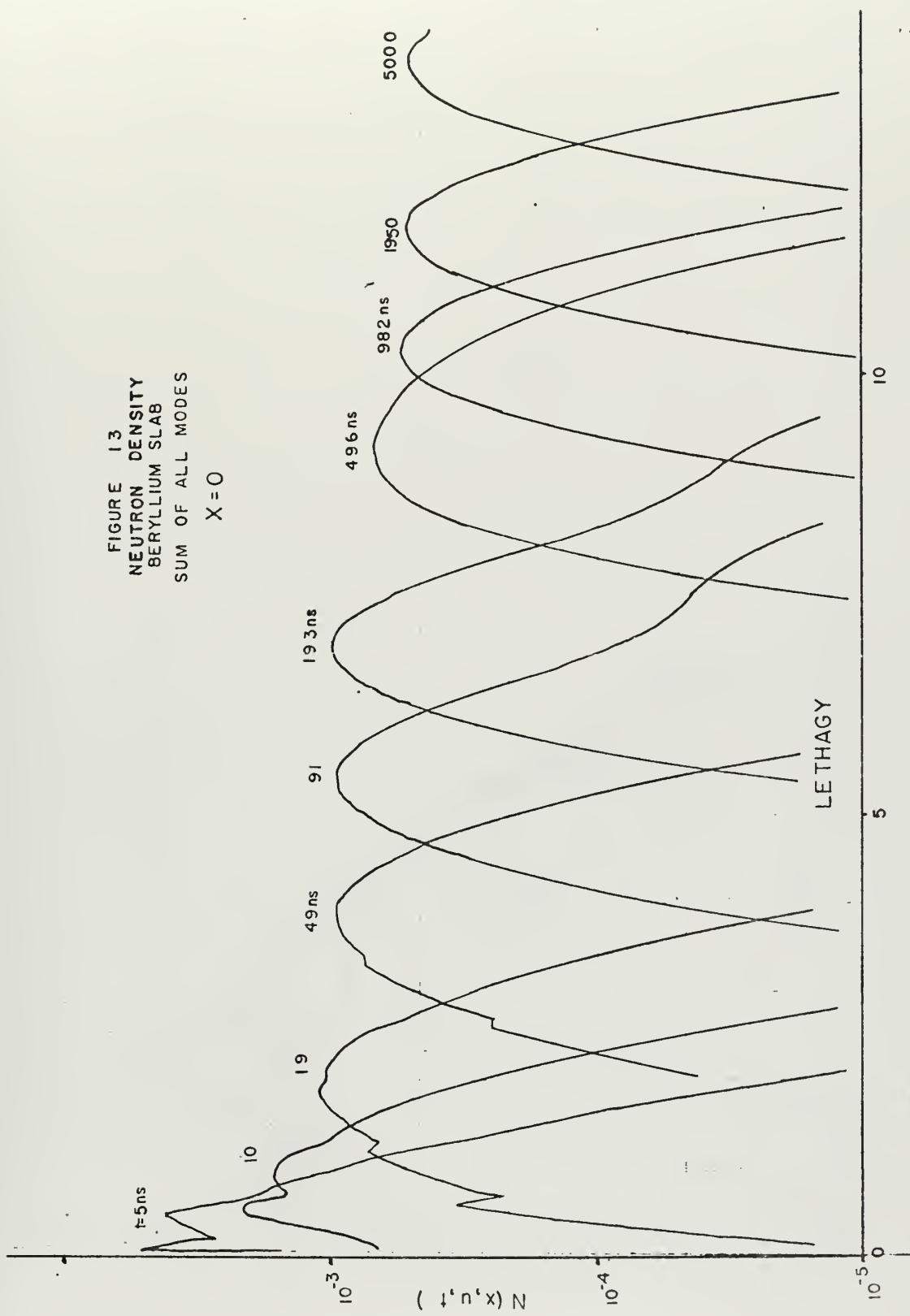




FIGURE 14  
NEUTRON DENSITY  
ASSEMBLY 6 F  
FUNDAMENTAL MODE  
 $X = 0$

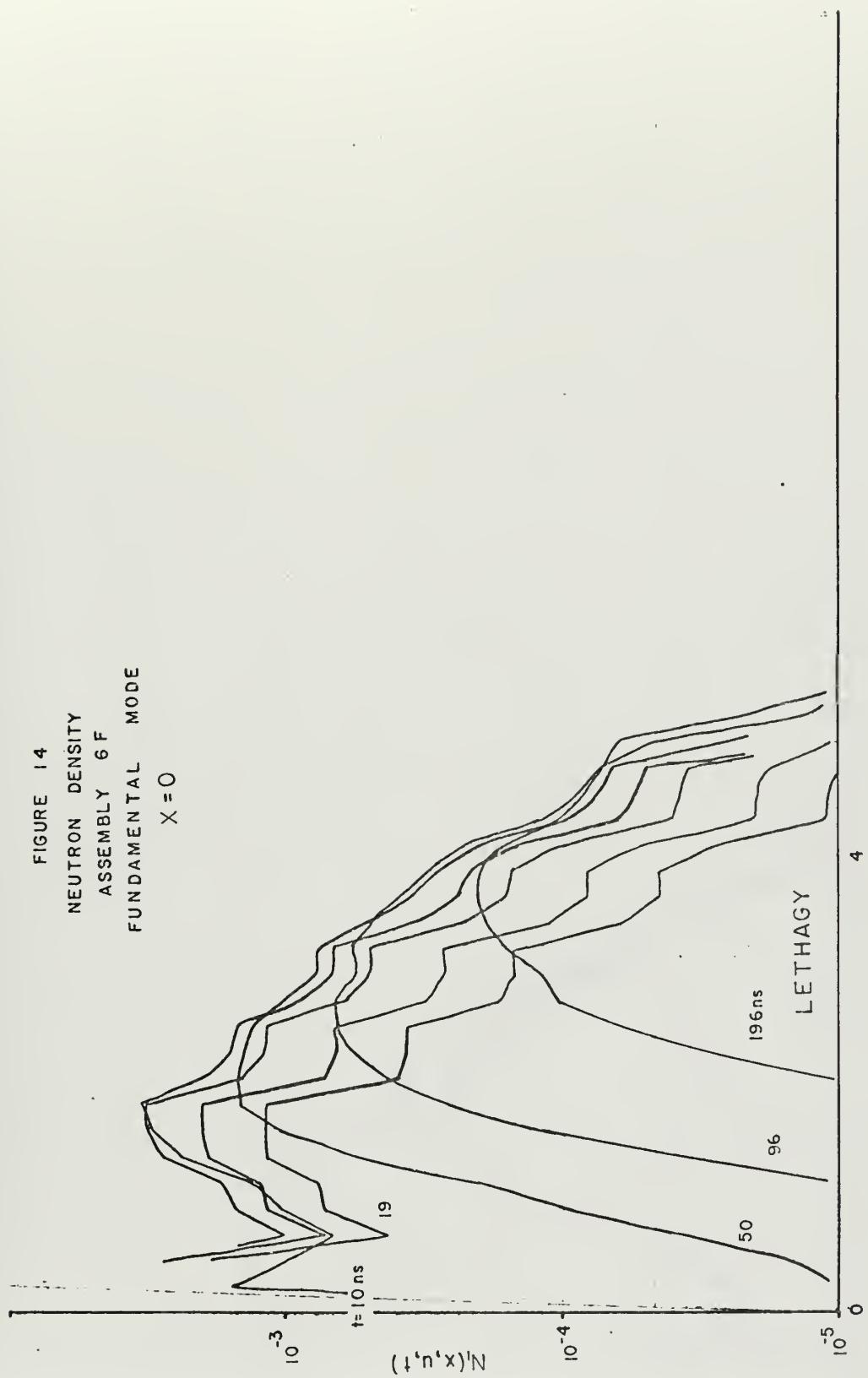




FIGURE 15  
NEUTRON DENSITY  
ASSEMBLY SF  
FIRST HARMONIC MODE  
 $X = 0$

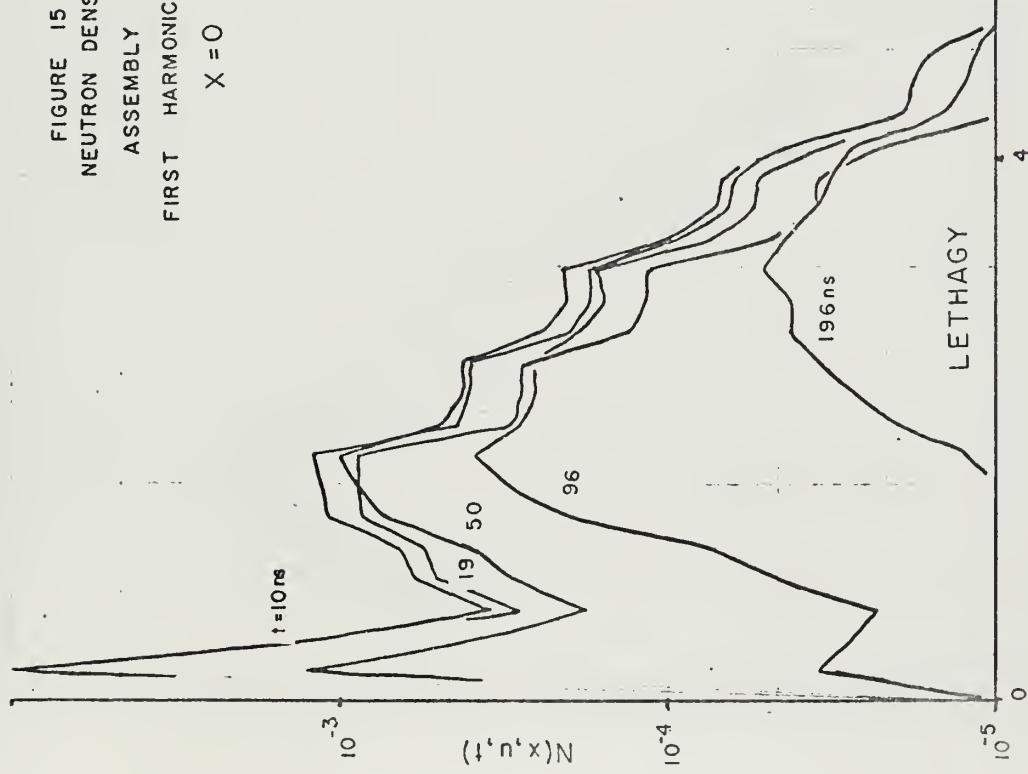




FIGURE 16  
NEUTRON DENSITY  
ASSEMBLY 6 F  
SUM OF ALL MODES  
 $X = 0$

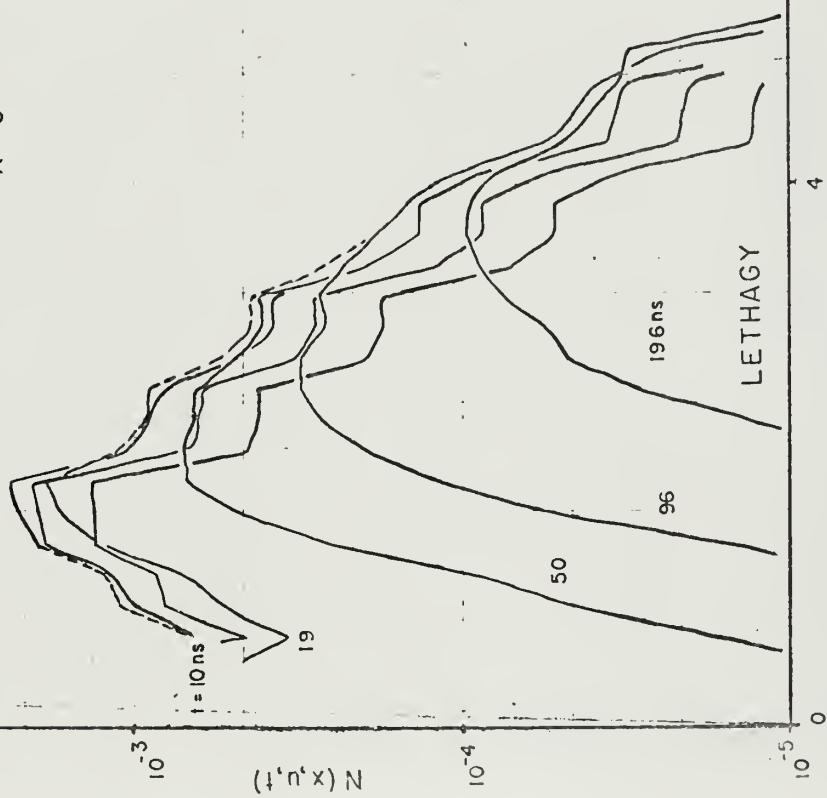




FIGURE 17  
BERYLliUM SLAB  
NEUTRON FLUX AT  $X = 0.0$

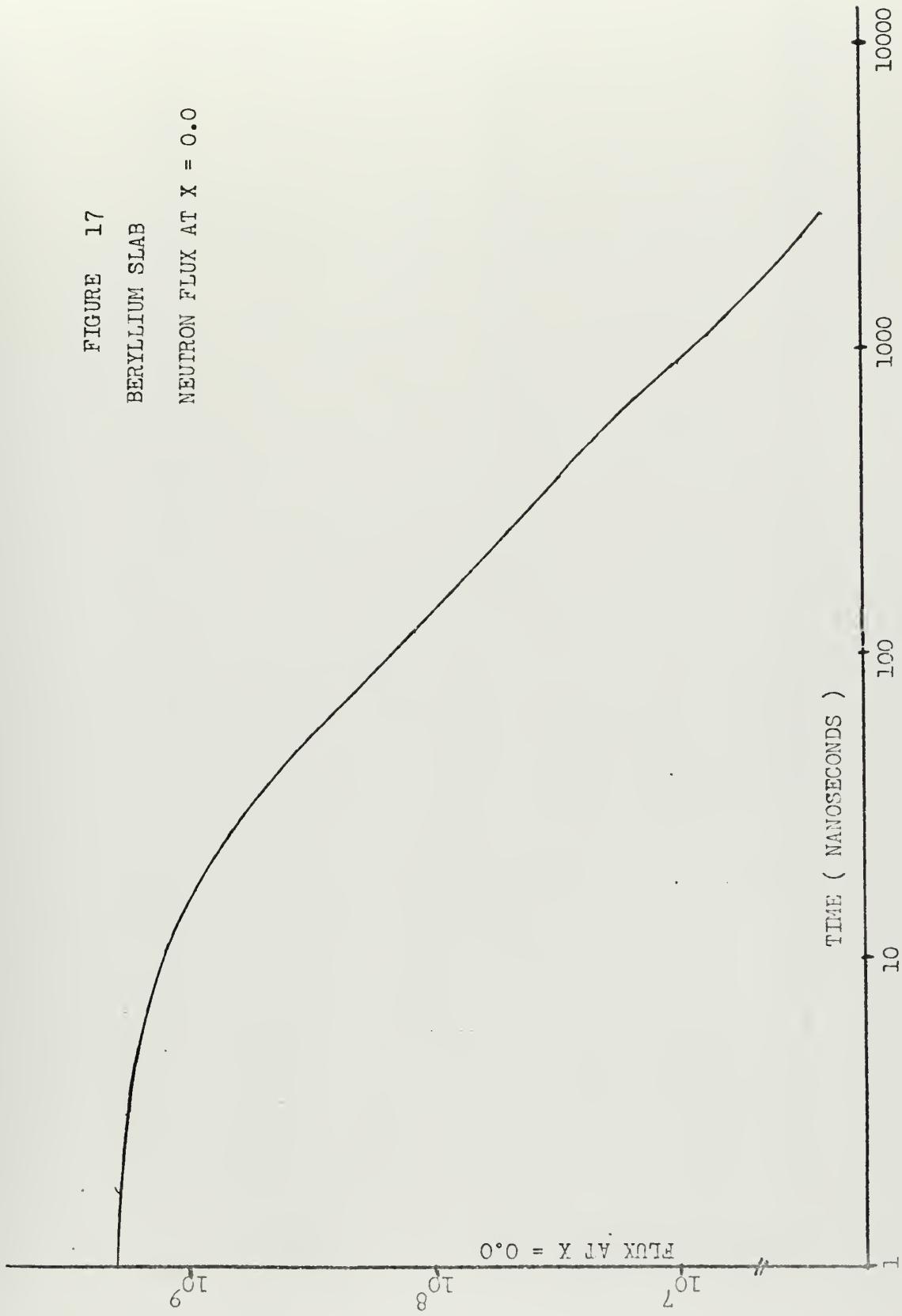




FIGURE 18  
ASSEMBLY 6 F  
NEUTRON FLUX AT  $X = 0.0$

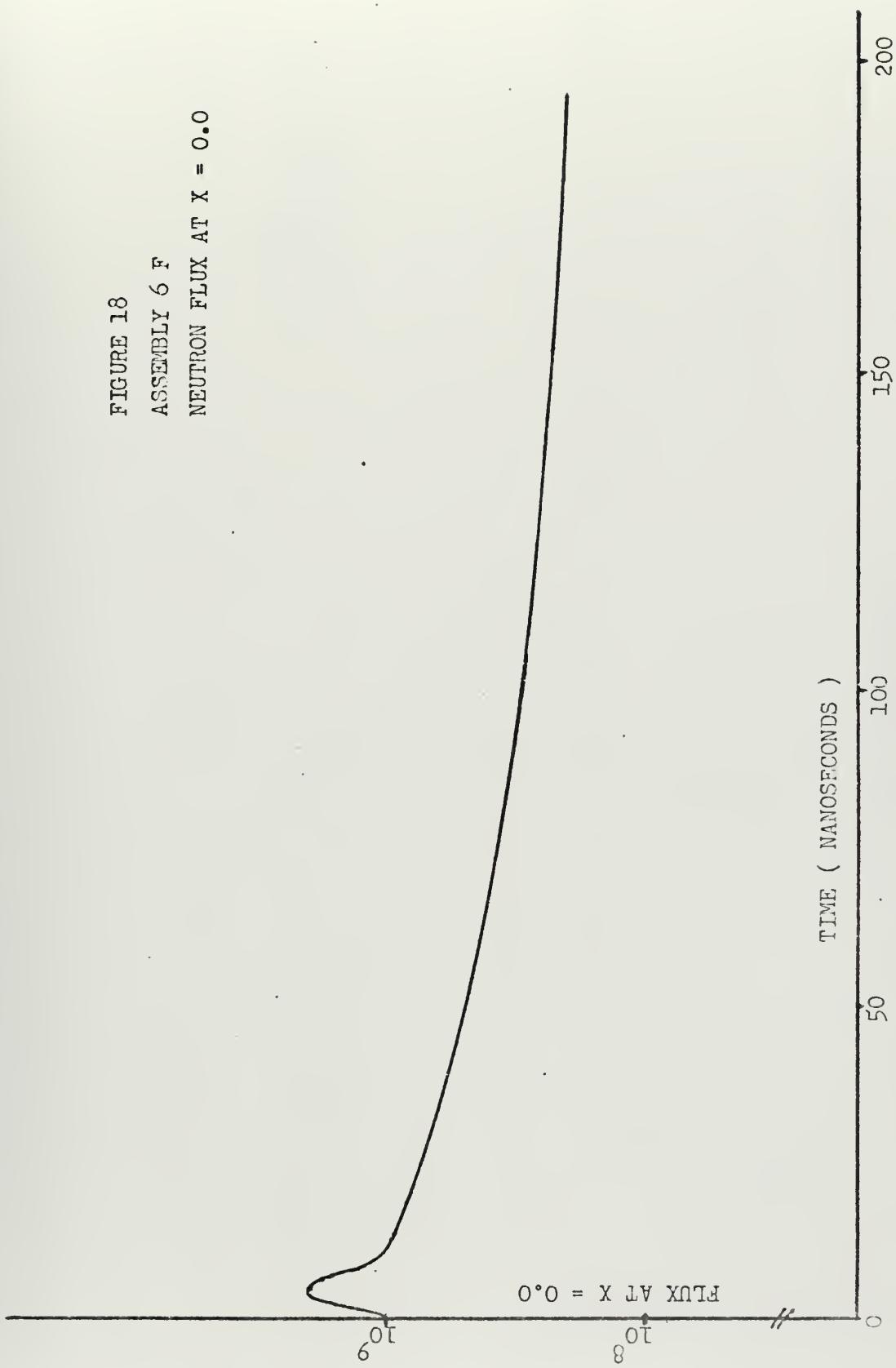




FIGURE 19  
MEAN ENERGY AT  $X = 0.0$   
BERYLLIUM SLAB

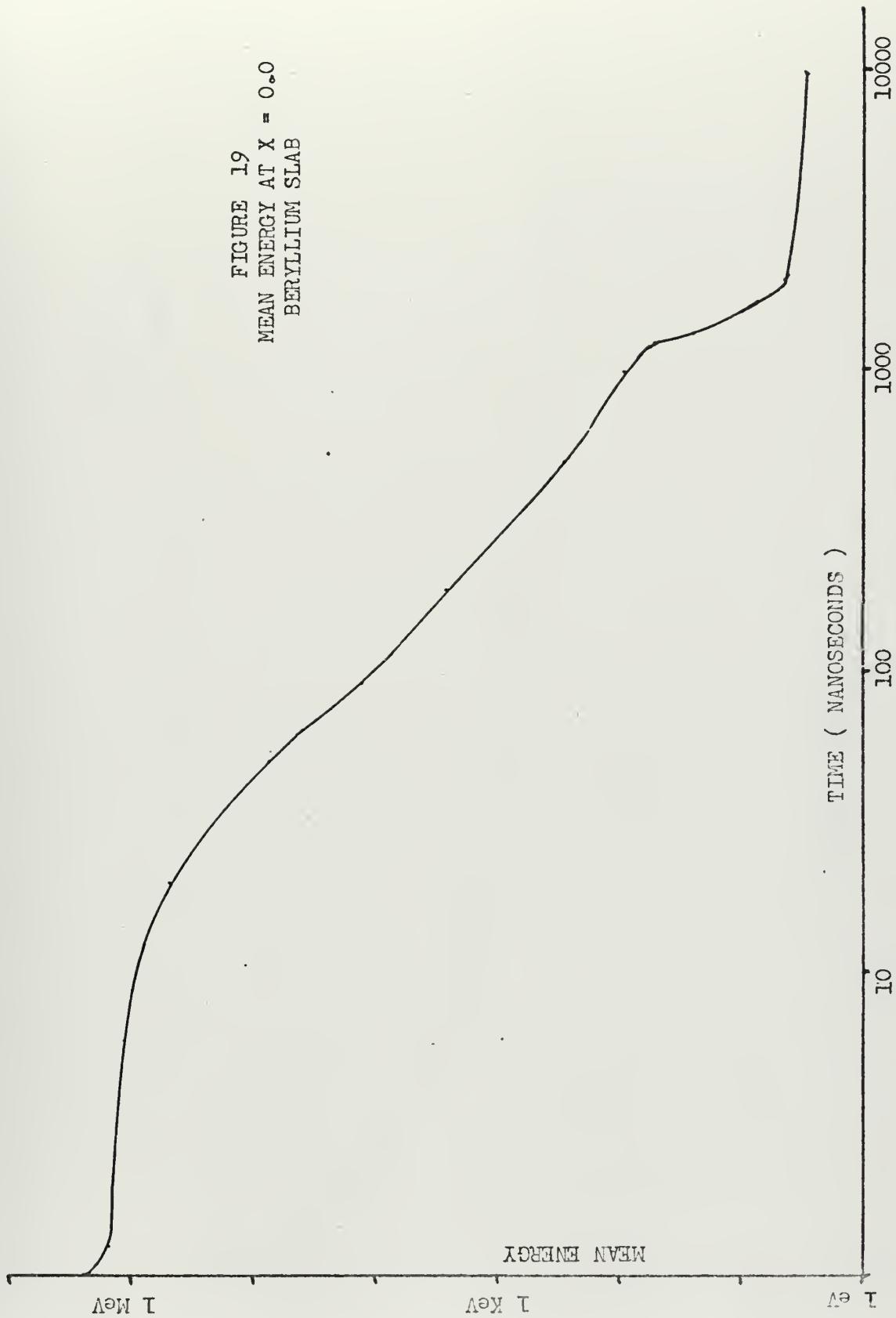
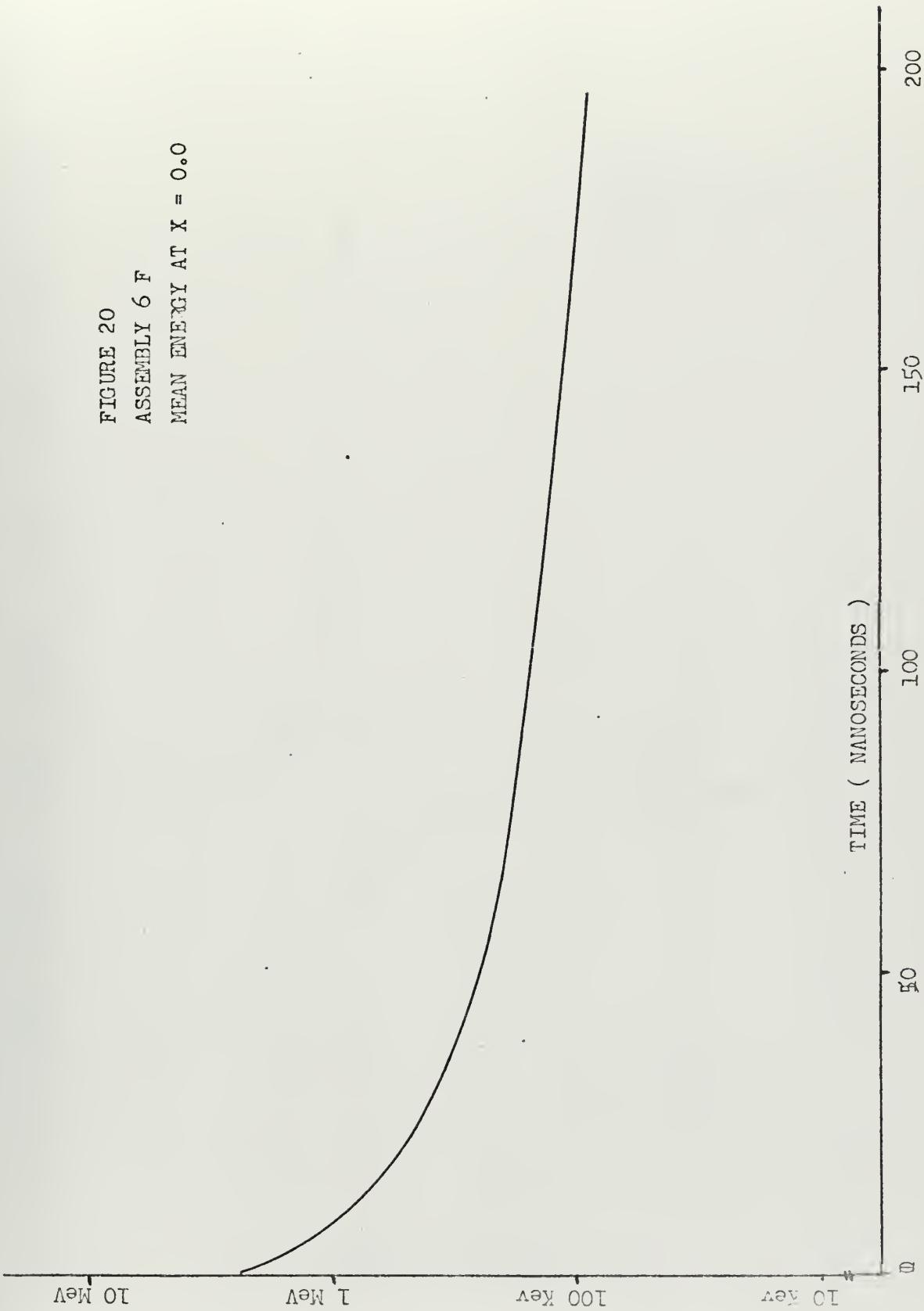




FIGURE 20  
ASSEMBLY 6 F  
MEAN ENERGY AT  $X = 0.0$





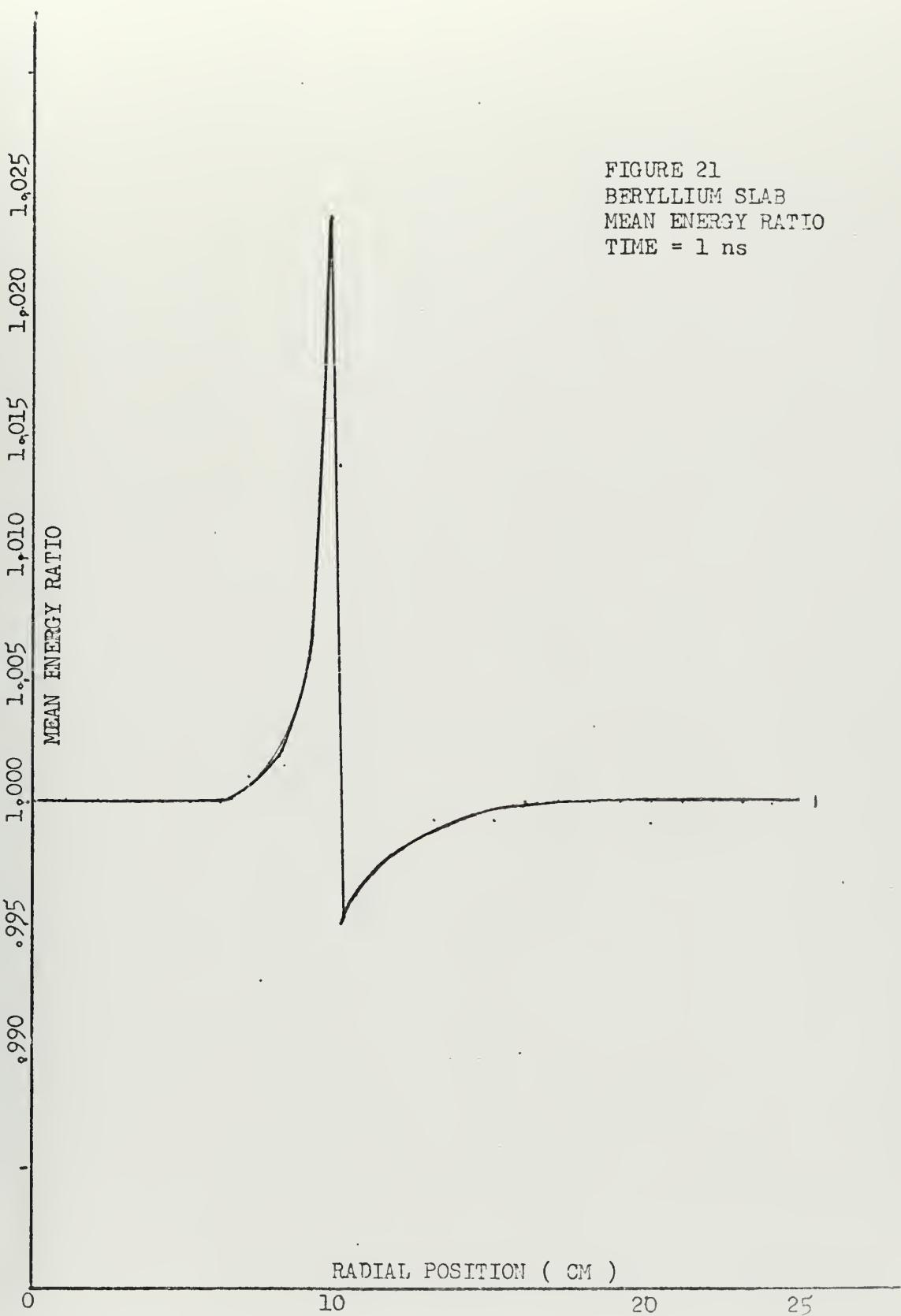
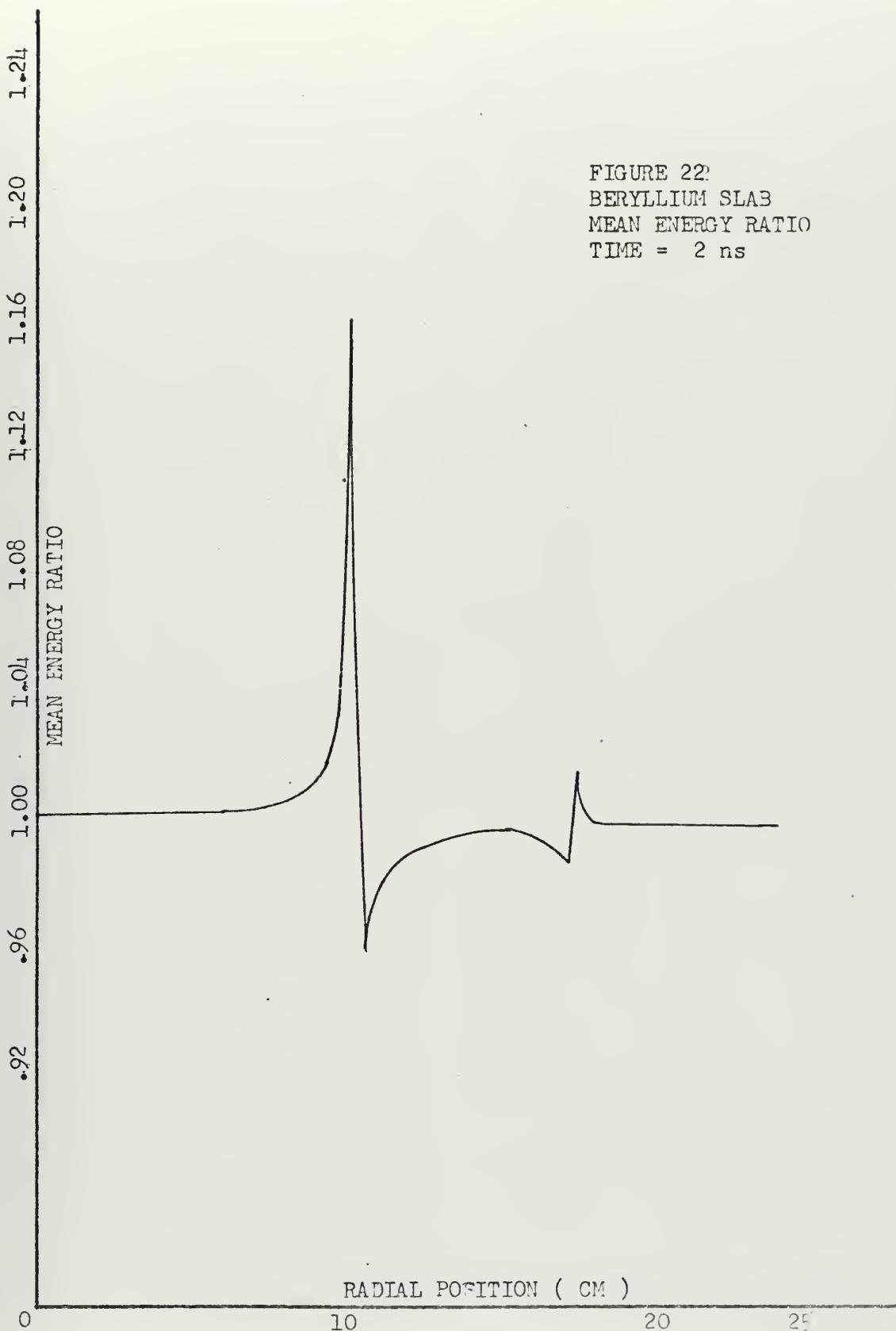




FIGURE 22  
BERYLLIUM SLAB  
MEAN ENERGY RATIO  
TIME = 2 ns





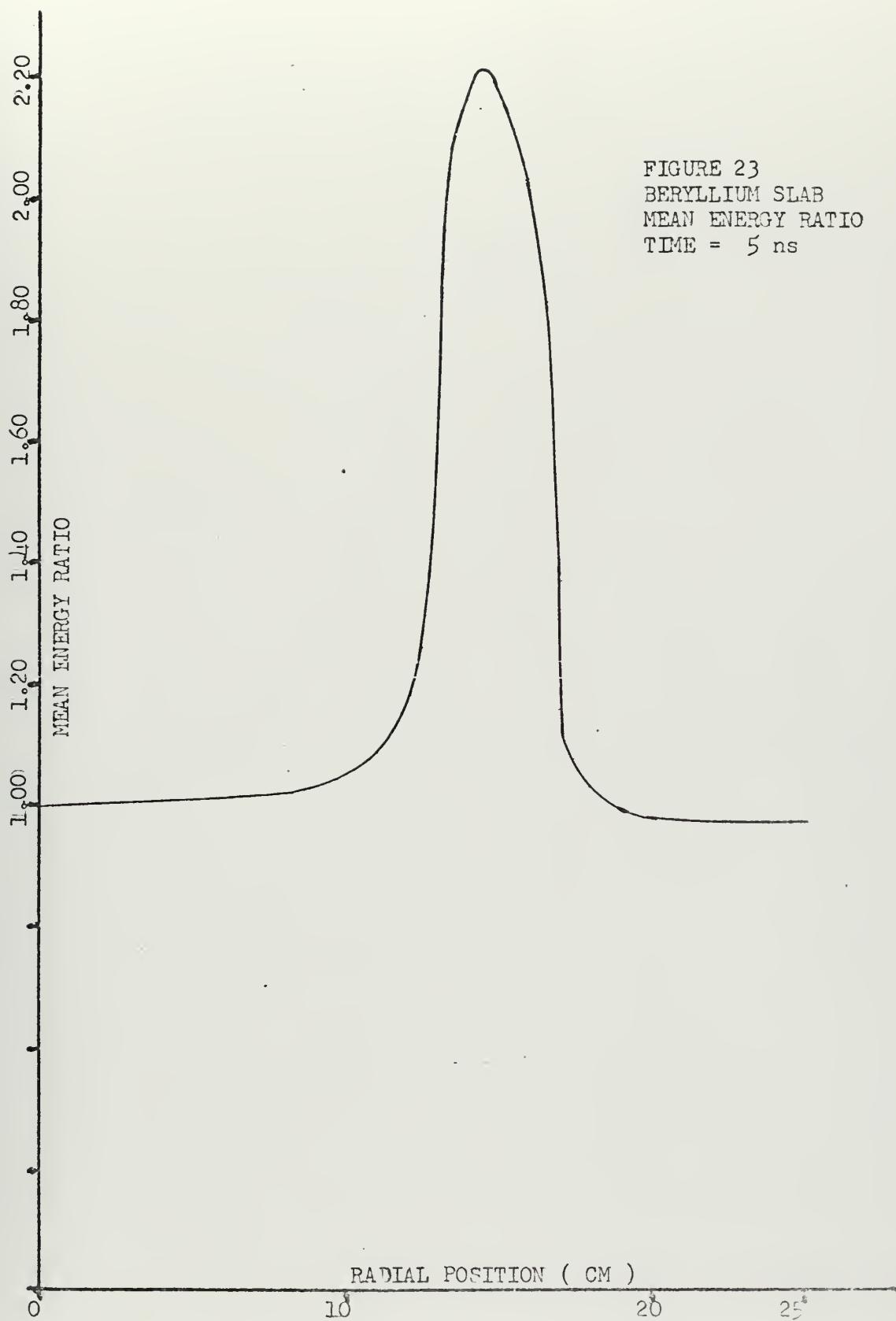


FIGURE 23  
BERYLLIUM SLAB  
MEAN ENERGY RATIO  
TIME = 5 ns



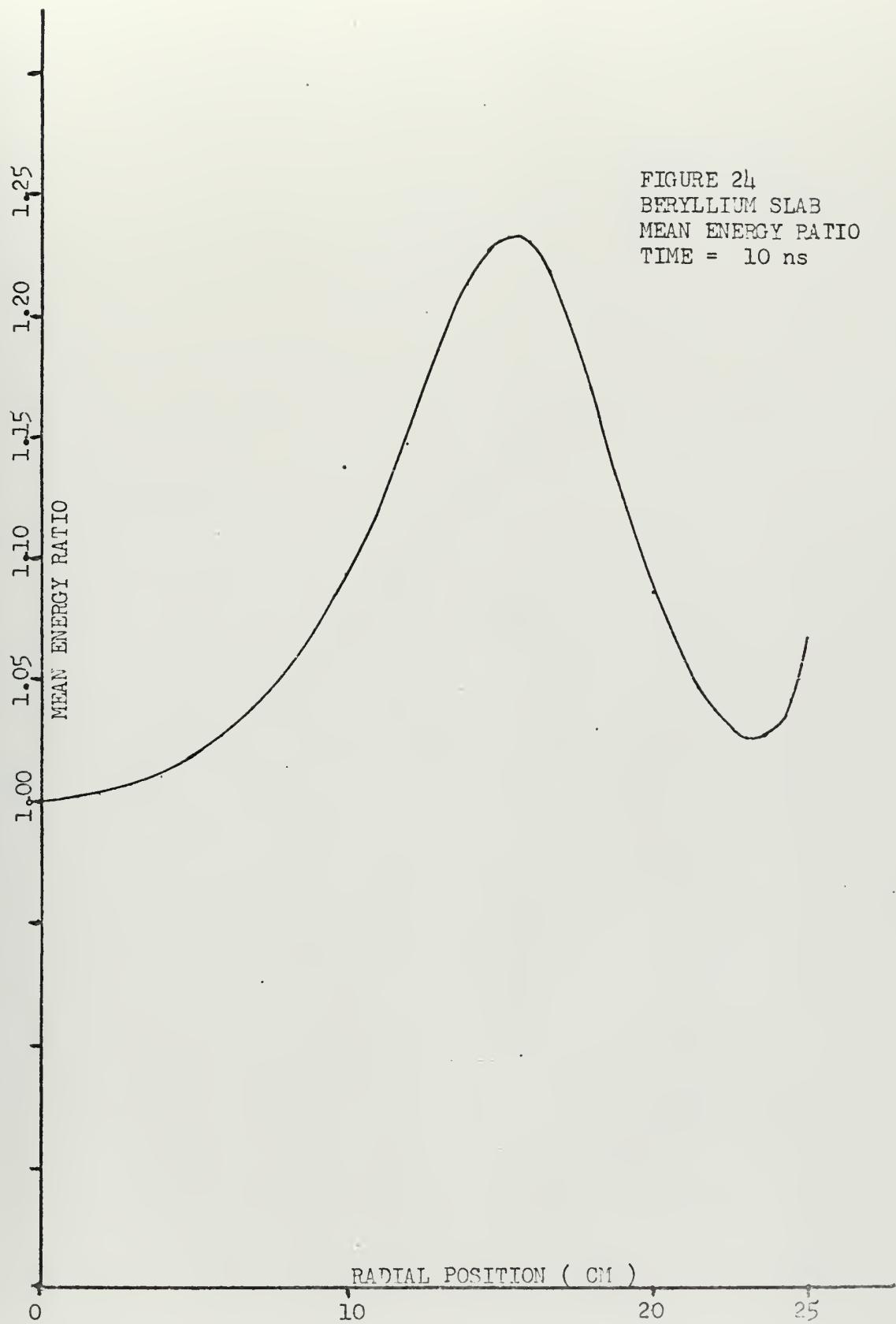




FIGURE 25  
BERYLLIUM SLAB  
MEAN ENERGY RATIO  
TIME = 19 ns

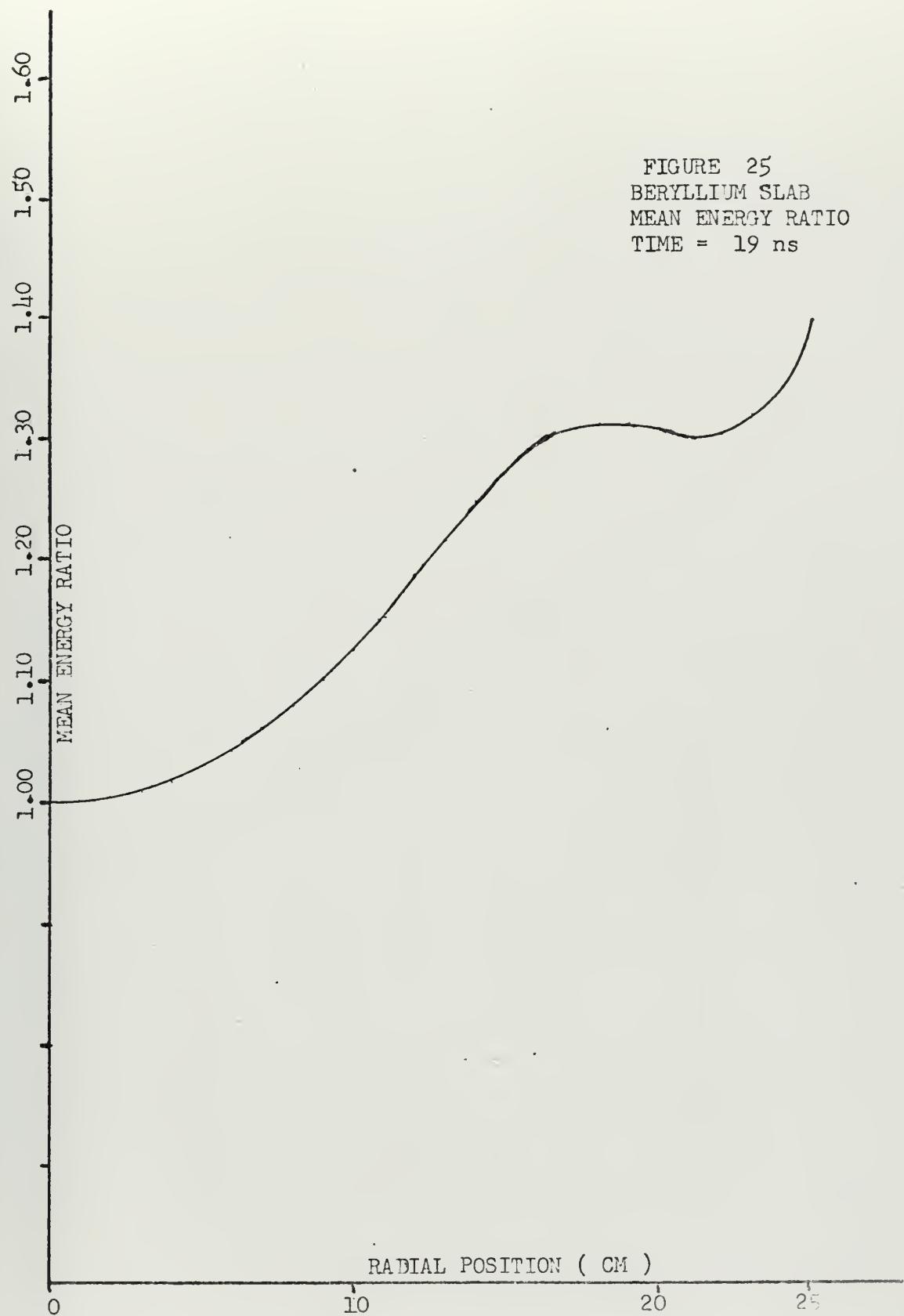
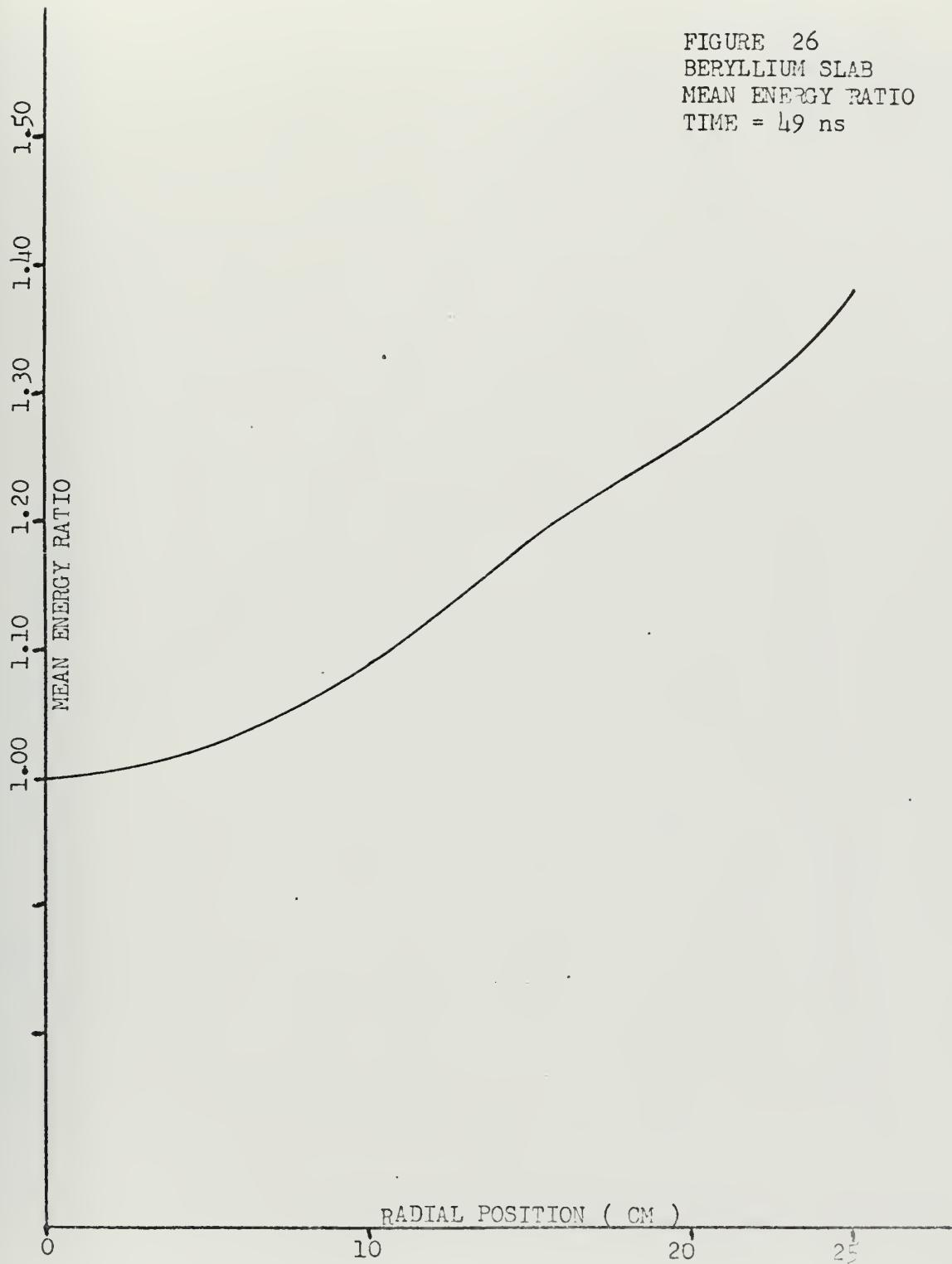




FIGURE 26  
BERYLLIUM SLAB  
MEAN ENERGY RATIO  
TIME = 49 ns





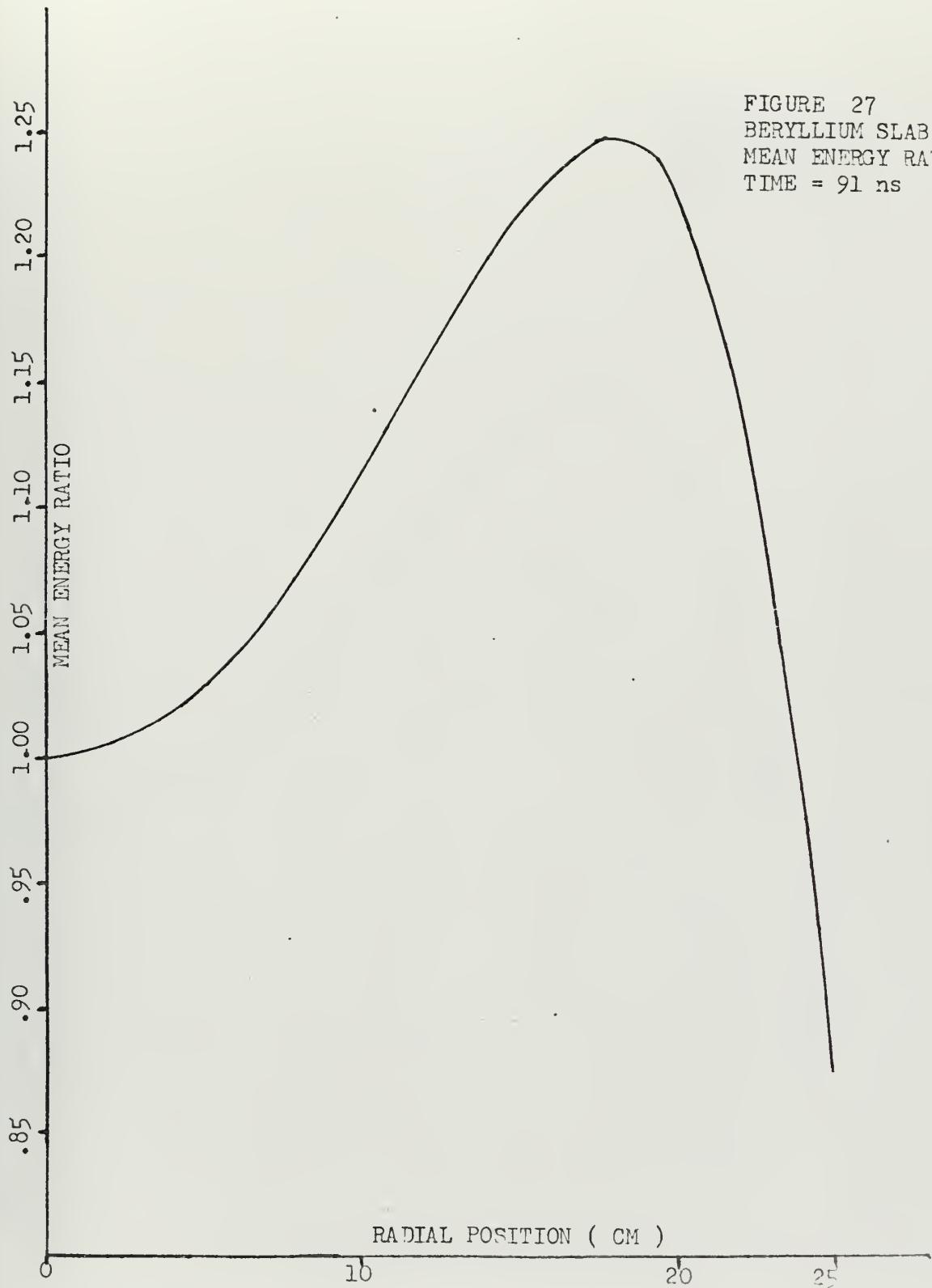
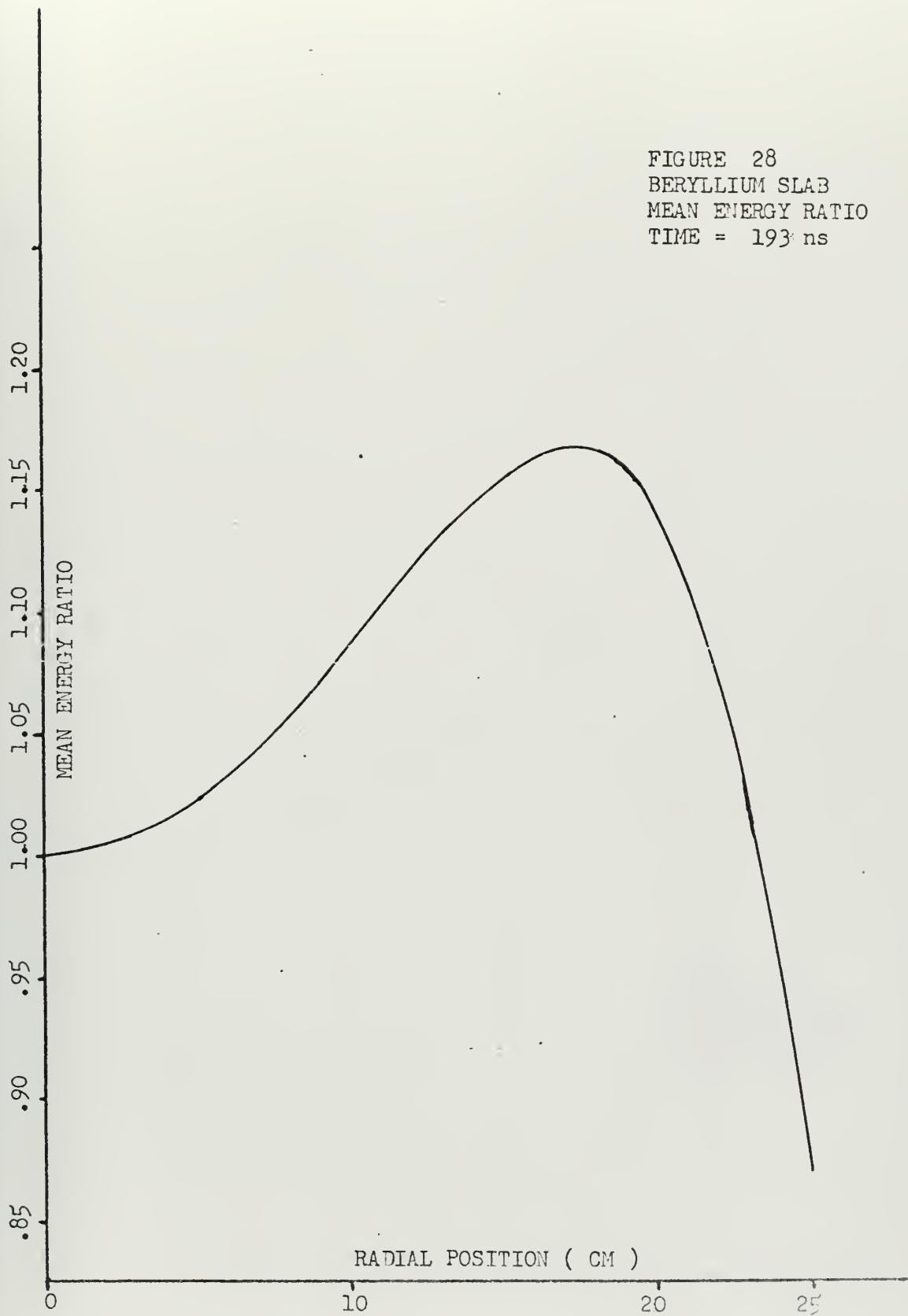




FIGURE 28  
BERYLLIUM SLAB  
MEAN ENERGY RATIO  
TIME = 193 ns





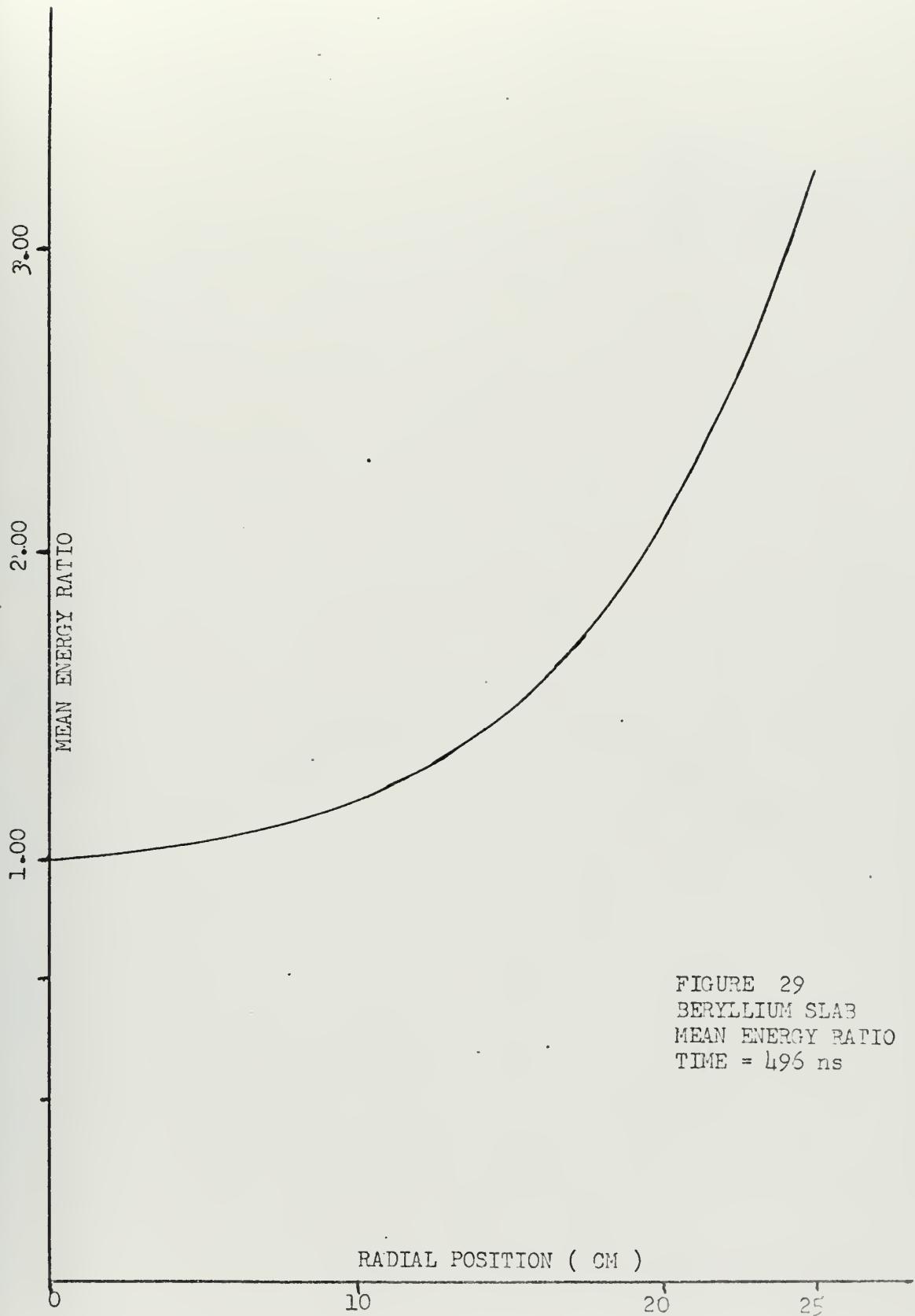


FIGURE 29  
BERYLLIUM SLAB  
MEAN ENERGY RATIO  
TIME = 496 ns



FIGURE 30  
BERYLLIUM SLAB  
MEAN ENERGY RATIO  
TIME = 982 ns

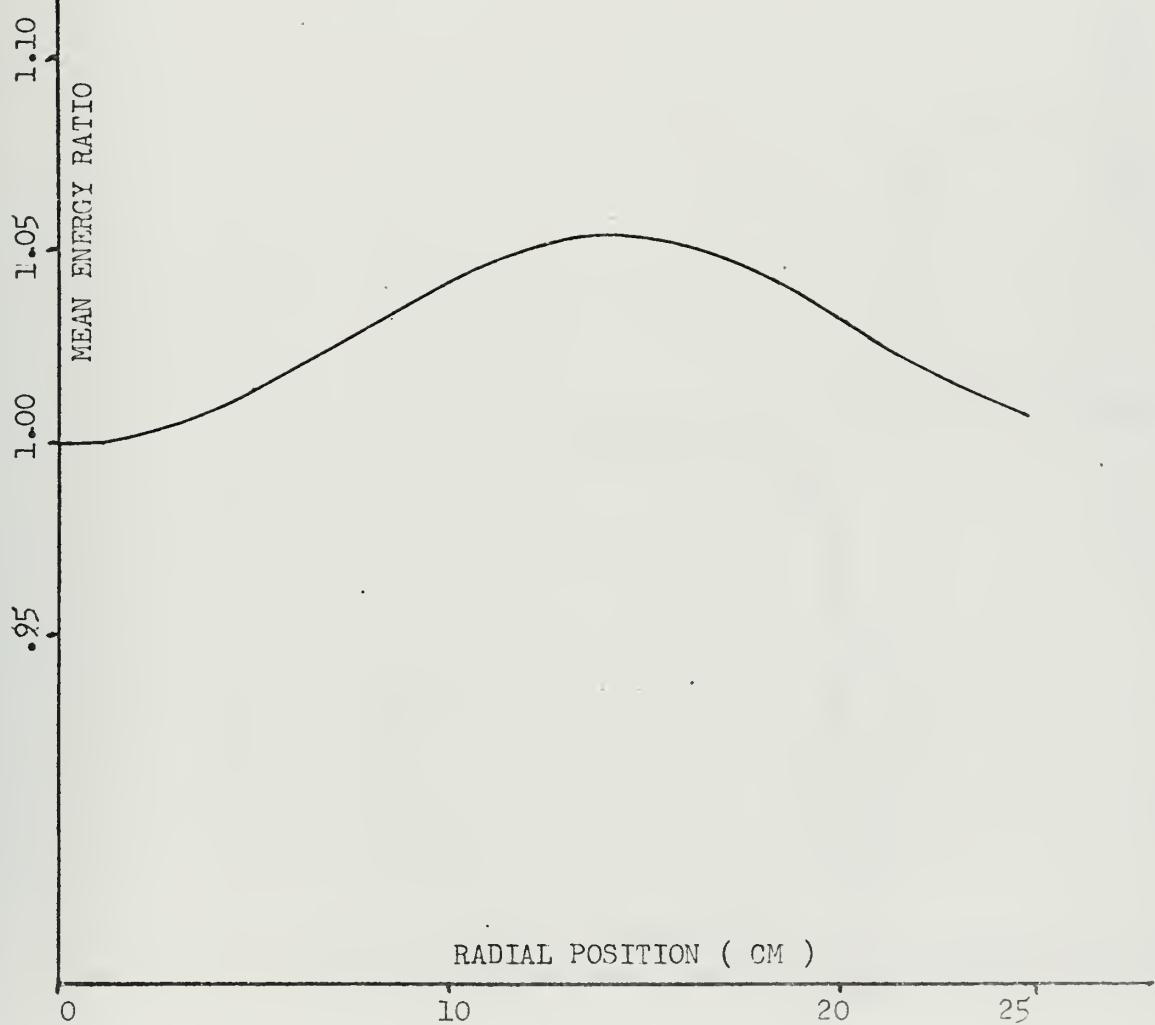




FIGURE 31  
ASSEMBLY 6 F  
MFAN ENERGY RATIO  
TIME = 2 ns

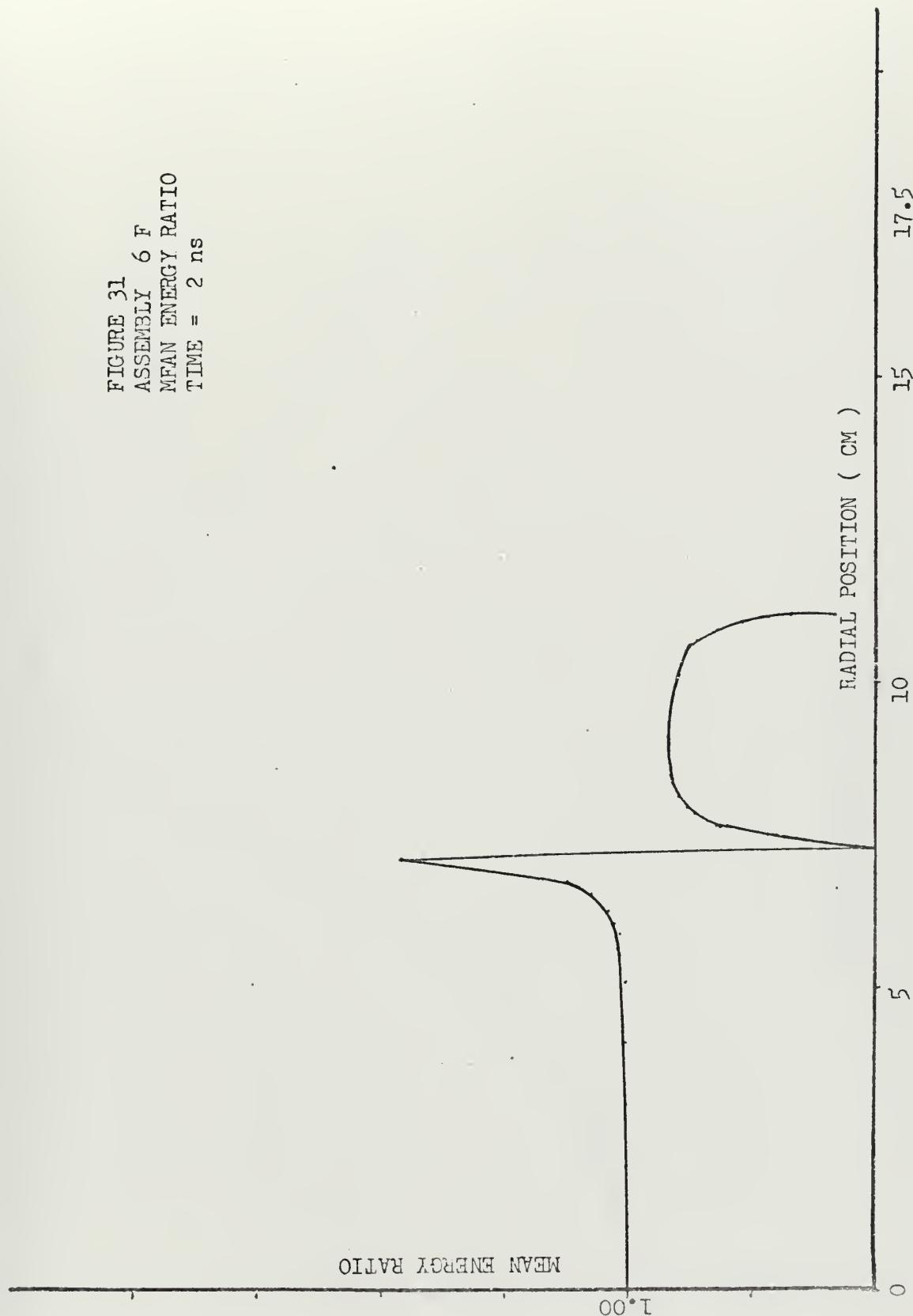




FIGURE 32  
ASSEMBLY 6 F  
MEAN ENERGY RATIO  
TIME = 10 ns

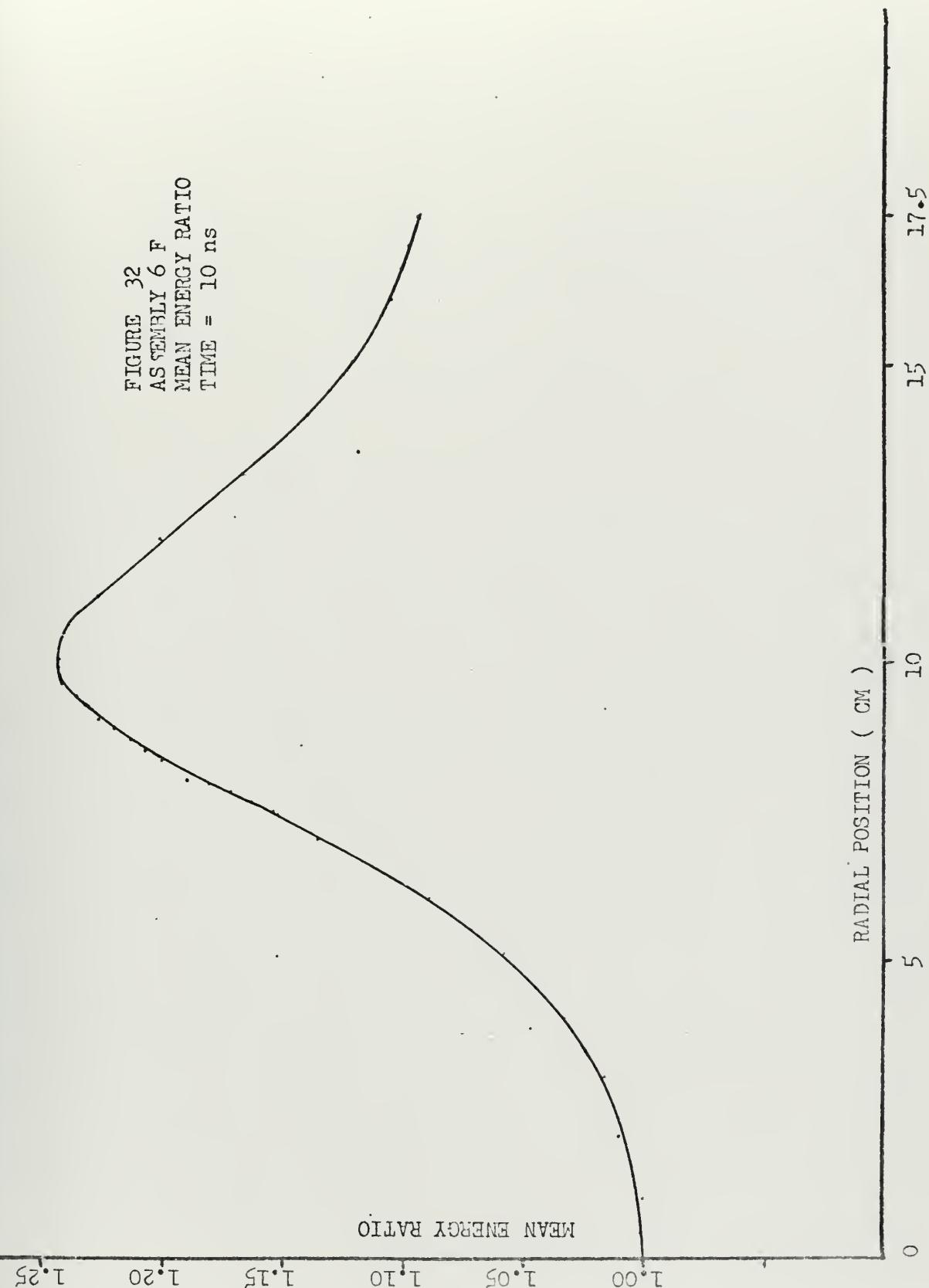




FIGURE 33  
ASSEMBLY 6 F  
MEAN ENERGY RATIO  
TIME = 19 ns

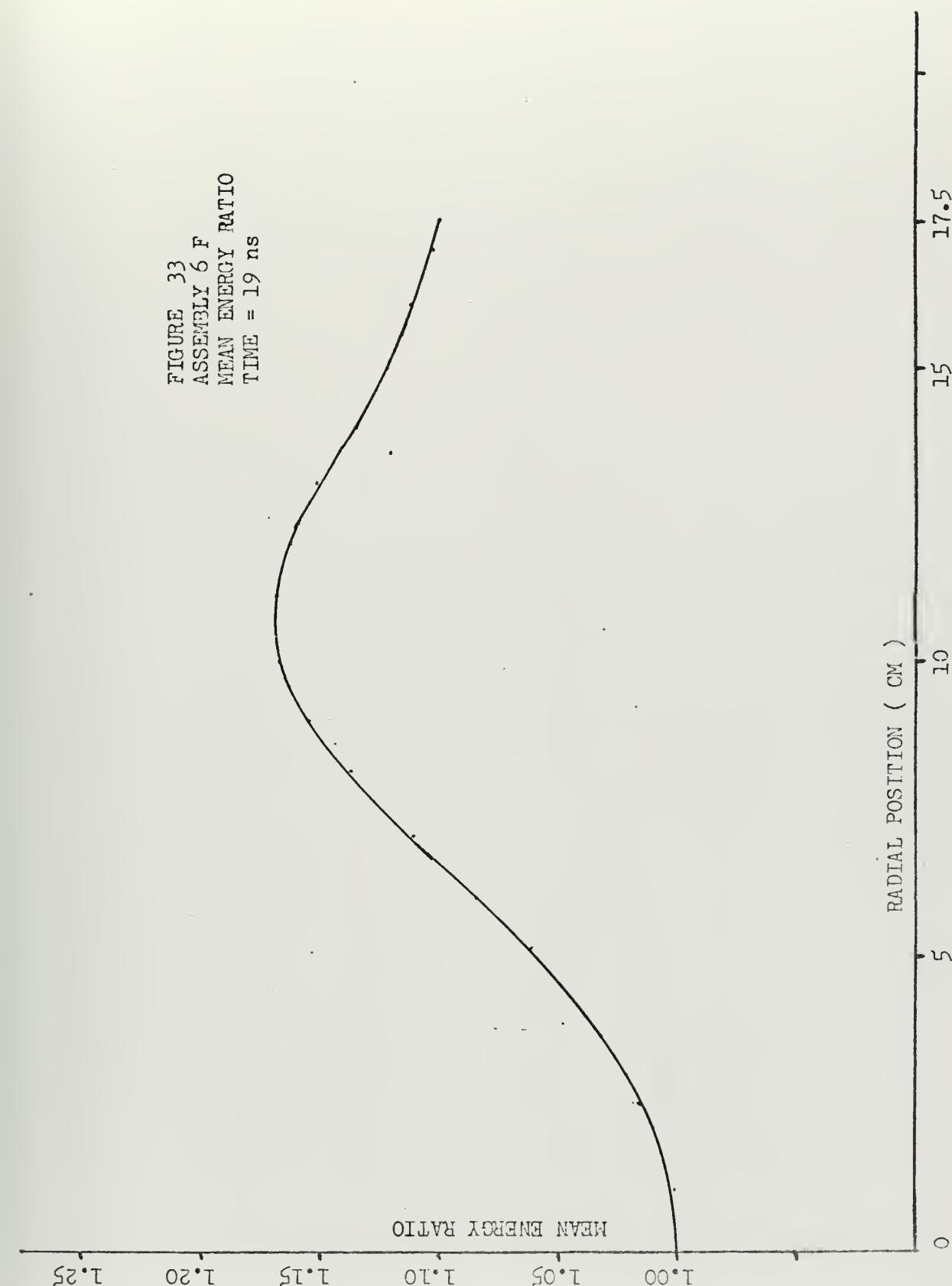




FIGURE 31<sub>4</sub>  
ASSEMBLY 6 F  
MEAN ENERGY RATIO  
TIME = 50 ns

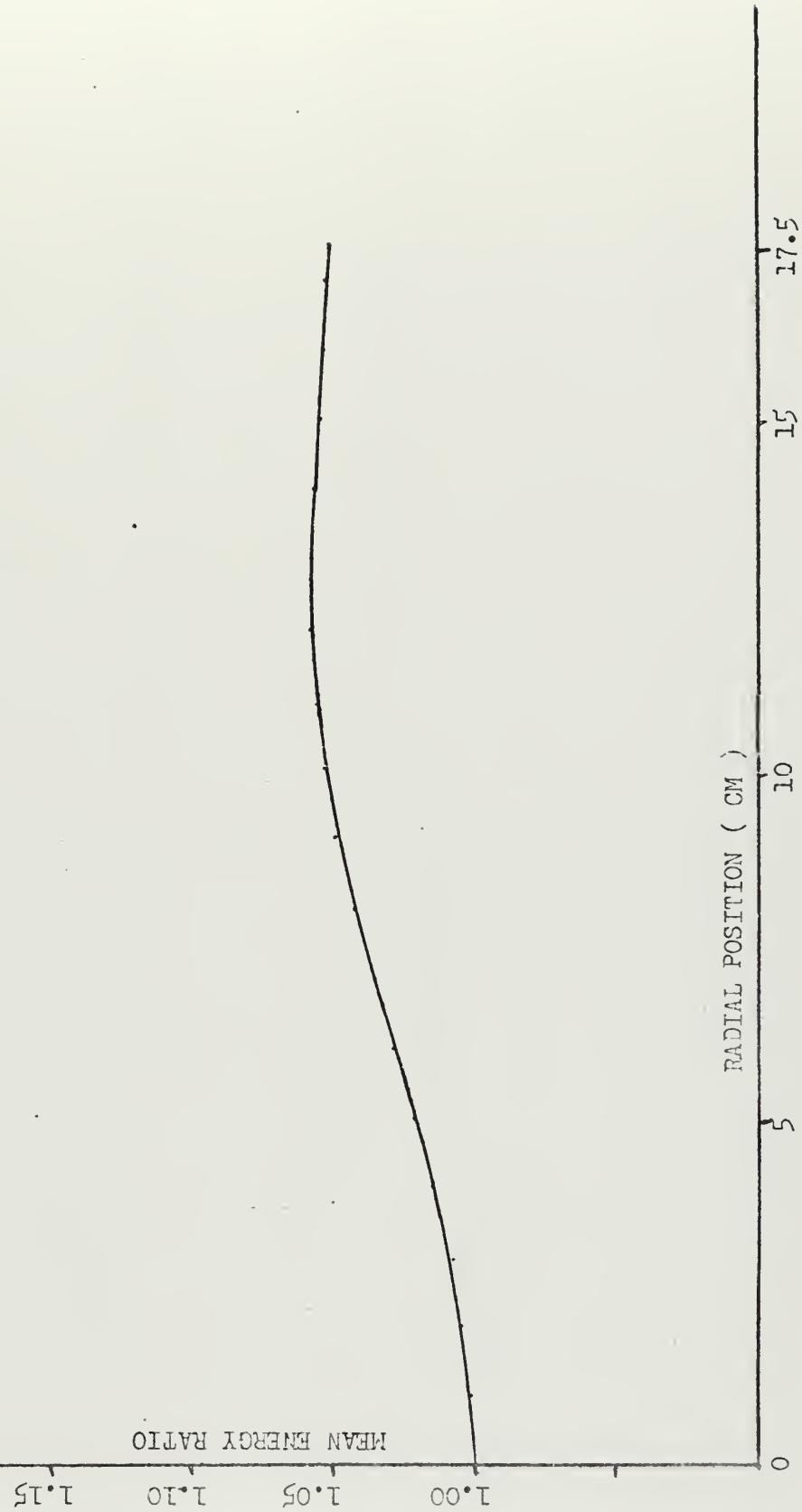




FIGURE 35  
ASSEMBLY 6 F  
MEAN ENERGY RATIO  
TIME = 96 ns

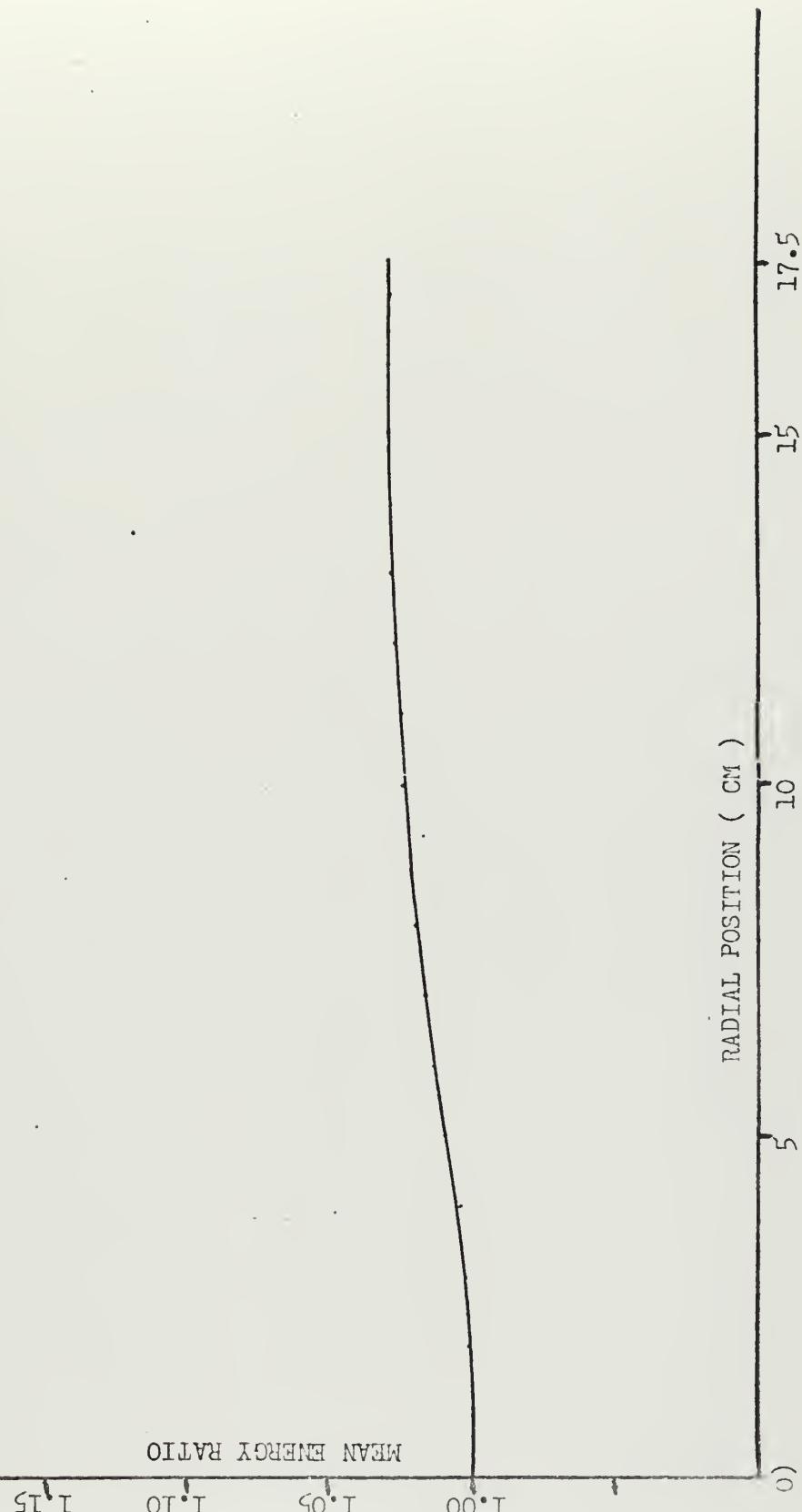




FIGURE 36  
ASSEMBLY 6 F  
MEAN ENERGY RATIO  
TIME = 196 ns

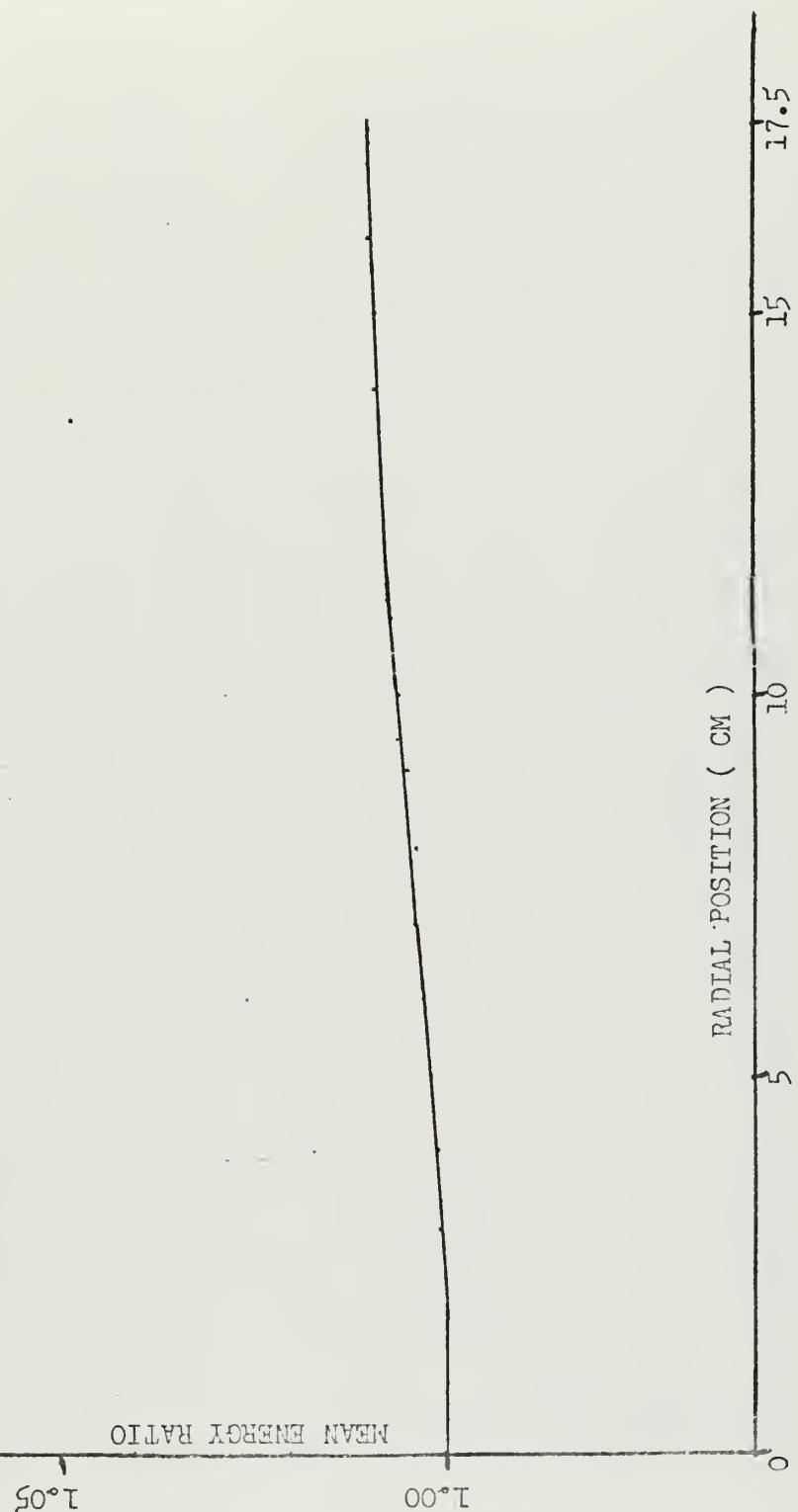




FIGURE 37  
BERYLLIUM SLAB  
DETECTOR RESPONSE  
 $X = 12.5$  CM

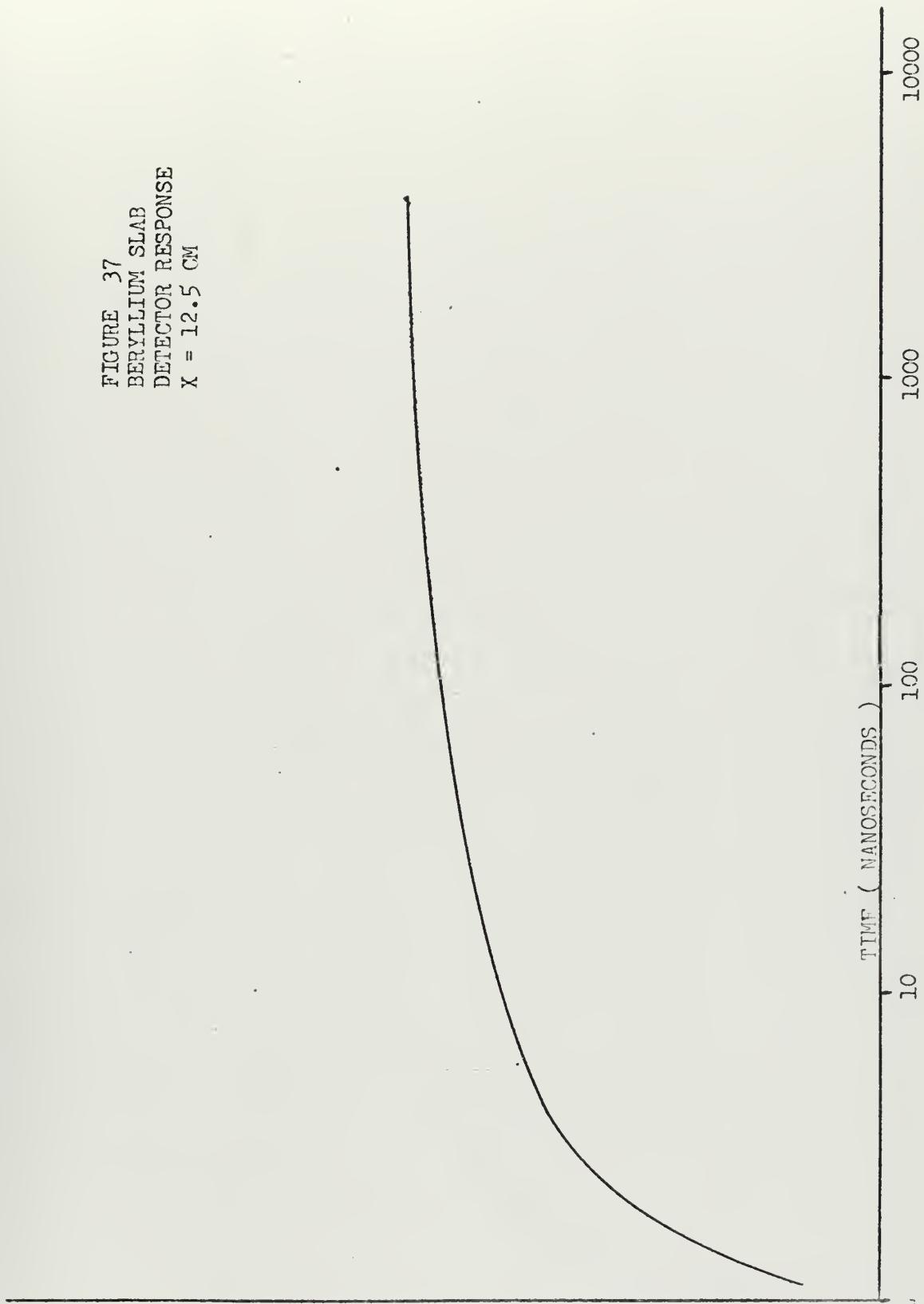
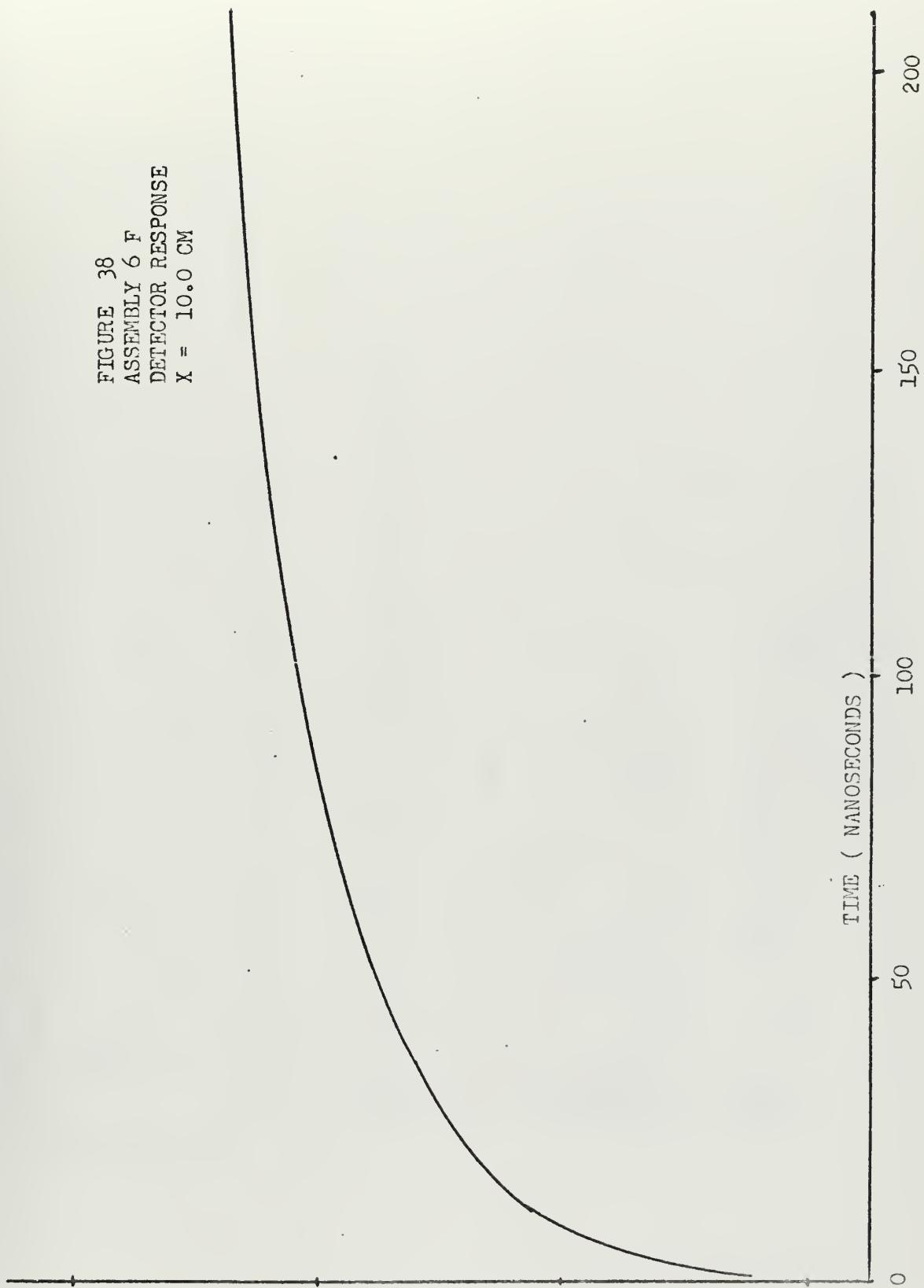




FIGURE 38  
ASSEMBLY 6 F  
DETECTOR RESPONSE  
 $X = 10.0$  CM





REVISED 2 JUNE 1971

MAIN PROGRAM INPUT DATA AND CONTROLS  
FIRST THREE CARDS DEFINE PROBLEM NAME FOR HEADINGS  
NAME LIST DATA IS INPUT TO DESCRIBE SYSTEM VARIABLES  
ONE CARDS FOLLOWS WITH SYSTEM CONTROL VARIABLES  
ONE CARD WITH LOGICAL VARIABLES

EXAMPLE DATA SET  
PROBLEM - MIL = NNN - NNN

```

NOUT = 1 PRINT OUTPUT
NOUT = 2 PLOT ON THE PRINTER USING ROUTINE -- PLOTP
NOUT = 3 CALL CAL COMP PLOTTER
NOUT = SOURCE TYPE 1 = UN AXIS SOURCE - DEFINE SOH & SA
NOUT = SOURCE TYPE 2 = OFF AXIS SOURCE -- DEFINE SOH,SA,S1
NOUT = SOURCE TYPE 3 = FIRST COLLISION SOURCE -- DEFINE SOH,DIFLEN
NOUT = SOURCE TYPE 4 = OFF-AXIS FIRST COLLISION SOURCE
NOUT = SOURCE TYPE 5 = EXTERIOR SOURCE - RAMP FUNCTION
IS = SOURCE

```



SOH = HALF WIDTH OF THE SLAB  
 SA = HALF WIDTH OF SOURCE  
 S1 = MID POINT OF OFF AXIS SOURCE  
 S2 = MID POINT LOCATION OF DETECTOR  
 S3 = HALF WIDTH OF DETECTOR  
 S4 = HALF WIDTH OF DETECTOR  
 DIFFLEN = 1.0 / SIGMAT - INVERSE TOTAL MACROSCOPIC CROSS SECTION  
 MODES = NN - TOTAL MODES FOR PROGRAMMING FOLD IN LETHARGY  
 NSTAT = M - STATE NUMBER 2  
 IBX = NUMBER OF POINTS ON X AXIS DEFINES DELTA FUNCTION 2  
 XBX(NN) INPUTS THE POSITIONS TO DO CALCULATIONS FOR  
 TO A MAXIMUM OF 100 DATA POINTS ALONG THE X-AXIS  
 EMAX = MAXIMUM ENERGY OF NEUTRON THAT A DETECTOR CAN MEASURE  
 EMIN = MINIMUM ENERGY THAT THE DETECTOR CAN MEASURE

USE OF THE LOGICAL CONTROL STATEMENTS - IN SUBROUTINE FOUT5 SPECTRAL  
 EXTERNAL CONTROL - PLOT THE SPECTRA AT POINT OTHER THAN X=0.0  
 USE INPUT VARIABLES NANS(6) FOR TOTAL POINTS TO DO.  
 USE INTERNAL CONTROL - BYPASS XBX(NN)  
 USE INTERNAL CONTROL - PRINT SEQUENCE IN FOUT2  
 USE INTERNAL CONTROL - PROFILE VS X ON CALCOMP  
 USE INTERNAL CONTROL - PLOT FLUX PROFILE VS X ON CALCOMP IN FOUT2  
 PLOT MEAN ENERGY RATIO VS X ON CALCOMP IN FOUT2  
 PLOT MEAN ENERGY - PERFORM SIMPLE PLOT OF SOURCE  
 EXTERNAL CONTROL - PRINT ON THE PINTERMINATE RUN AFTER INCON1  
 EXTERNAL CONTROL - TERMINATE RUN IN PROBLEM DEFINITION OR CONTROL MODES  
 L(5) = 1 - CALL FOUT2-CALC MEAN ENERGY-E(X,T) AND NEW F(X,T)  
 L(7) = 1 - CALL FOUT2-CALC MEAN ENERGY-E(X,T)  
 L(8) = 1 - CALL FOUT2-CALC MEAN ENERGY-E(X,T)  
 L(9) = 1 - CALL FOUT2-CALC MEAN ENERGY-E(X,T)  
 L(10) = 1 - CALL FOUT2-CALC MEAN ENERGY-E(X,T)  
 L(14) = 1 - CALL FOUT2-CALC MEAN ENERGY-E(X,T)  
 L(16) = 1 - CALL FOUT2-CALC MEAN ENERGY-E(X,T)  
 L(17) = 1 - CALL FOUT1 FOR INPUT DATA POINTS  
 L(18) = 1 - CALL FOUT1 FOR INPUT DATA POINTS - DETECTOR RESPONSE  
 L(19) = 1 - CALL SUBROUTINE FOUT3 - SPECTRAL PLOT WITHOUT  
 L(20) = 1 - CALL SUBROUTINE FOUT4 FOR SPECTRAL PLOT WITH  
 PROGRAM OUTPUT CONTROL VARIABLES  
 NANS(1) - EXTERNAL CONTROL - LIMIT TO TOTAL GRAPHS TO BE  
 PLOTTED IN SUBROUTINE FOUT1 - DEFAULT OF 1 GRAPH  
 NANS(2) - INTERNAL CONTROL - LABELING OF FOUT1 GRAPHS ON  
 CALLS FROM MAIN PROGRAM AND IN ROUTINE FOUT2  
 NANS(3) - EXTERNAL CONTROL - LIMIT TO TOTAL GRAPHS TO BE  
 PLOTTED BY FOUT2 - DEFAULT OPTION = 1  
 NANS(4) - EXTERNAL CONTROL - LIMIT TO ALLOWABLE NUMBER  
 NANS(5) - NOT USED  
 NANS(6) - NUMBER OF POINTS FOR SPECTRAL RESPONSE PLOTTING



WILL ONLY BE CONSIDERED FOR CASES WITH  $L(4) = 1$

MODIFICATION OF MIL SIX MOD 2  
COMMON BLOCKS 3 AND 6 MUST BE REVISED TO HANDLE MORE THAN  
SIX HARMONIC MODES OF THE FUNCTIONAL INPUT DATA  
COMMON/CON1/ XBX(100), U(150), E(150), TIM(20), DUMMY(150), TXBX(100)  
COMMON/CON4/ SOH, SAV, S1, S2, S3, S4, SZERO, BUCKLE, DEMAX, EMIN, DIFLEN, X  
COMMON/COM7/ S01, S02, S03, S04, S05, AK11, AK12, AK13, AK14, AN3A, AN3B,  
1AN3C, AN4, ANA4, AN2X, AK15, AK5  
COMMON/COM2/ NT(20), NPROB(54), NAN5(6), L(20), LTAB(20)  
COMMON/COM5/ MODES, IS, IDATA, N, NVIR, NI, NF, NCALL, NOUT, IBX,  
INSTANT, MODE  
COMMON/COM3/ FXT(100,20), FUTX(150,20), BUCK(6), FUT(6,71,20)  
COMMON/COM6/ NTOP(6,20), NM(6,20), IKNT(6)

C CALL CANCEL(2)

C SET INITIAL CONDITIONS OF THE PROGRAM

C 1 CALL INCON1  
C INTERNAL PROTECTION IF INPUT ERROR IN PROBLEM DEFINITION

C IF (L(1) .EQ. 1) GO TO 52

C 2 IF (IDATA .EQ. 0) GO TO 60

C IF (IDATA .EQ. 1) CALL READ10(LASTN)

C IF (IDATA .EQ. 2) CALL READ11(LASTN)

C 2 SUBROUTINE ESTIM1  
C TOTAL PROGRAM RUNNING-EXECUTION TIME DURING TESTING  
C INITIAL TIME ESTIMATING OF THE PROGRAM EXECUTION TIME  
C REQUIREMENTS WHILE STILL DOING INITIAL DEBUGGING AND CHECKOUT

C CALL ESTIM1

C CALL MODEL1

C INITIAL DEFINITION AND ERROR ANALYSIS BY ROUTINE MODEL1

C IF (LASTN .EQ. 1) GO TO 50  
C THIS WILL BY PASS ALL FOLLOWING STEPS IF THERE IS AN INPUT ERROR

C 6 CONTINUE  
C PARAMETER L(20) = 1 WILL CALL FOR A PLOT OF THE INPUT SPECTRUM  
C OF THE MOD-5 DATA WITH A PLOT OF THE DENSITY VECTOR  
C VERSUS THE LETHARGY VALUES  
C IF (L(20) .EQ. 1) CALL FOUT4

C FOUT5 WILL PLOT SPECTRUM FOR EACH MODE WEIGHTED BY HARMONIC  
C FOURIER COEFFICIENT FOR EACH MODE AND THE SUMMED SPECTRUM  
C FOR ALL MODES OF INPUT DATA  
C IF (L(17) .EQ. 1) CALL FOUT5

C CONTINUE

20



30 CONTINUE

LOGICAL PARAMETER L(18) = 1 RESULTS IN THE POSITION FUNCTION  
F(U,T,X) BEING CALCULATED FOR THE POSITIONS XBX(I)  
AS STATED IN THE INPUT DATA  
IF (L(18).EQ.0) GO TO 65

CALLING ROUTINE TO CALCULATE THE FUNCTION--F(U,T,X) FOR POSITIONS  
DESIRED AS LISTED IN THE INPUT DATA FOR THAT POINT  
IF ANY OUTPUT IS DESIRED FROM ROUTINE FOUT1 VIA INITIAL  
CALLING SEQUENCE FROM THE MAIN PROGRAM -- DEFINE L(16) = 0  
IN THE PROBLEM INPUT STRUCTURE

40 IF (IBX.EQ.0) XBX(1) = 0.  
IF (IBX.EQ.0) IS=1  
IF (IS.EQ.0) IBX=1  
PROTECTION STEP TO ALLOW FOUT1 TO PRINT OUT THE VALUES  
OF THE FUNCTION -- OF THE PROGRAM

L(16) = 0  
DO 43 IX=1,IBX  
IF (NOUT.EQ.3) NANS(2) = 11X  
WRITE (6,63)  
CALL FOUT1(XBX(11X))  
CONTINUE  
43 CONTINUE  
IF (L(14).EQ.1) CALL FOUT2

1 SUBROUTINE FOUT2 CALC MEAN ENERGY AND PLOT SUM  
OVER ALL LETHARGY STATES OF F(U,T,X) AND PLOT  
BOTH FUNCTIONS ON THE SAME PLOT FOR EACH TIME STEP

IF (L(19).EQ.0) GO TO 48  
PARAMETER L(19) = 1 AS DEFINED IN THE PROBLEM DEFINITION  
TO CALCULATE THE DECTECTOR RESPONSE IN SUBROUTINE FOUT3  
IF L(19) = 0 --- WE WILL BY-PASS THIS ROUTINE

48 WRITE (6,64)  
CALL FOUT3  
CONTINUE  
49 GO TO 60

50 THIS WILL ALLOW A BYPASS AND STOP IF ERROR  
WRITE (6,51)  
GO TO 60  
52 WRITE (6,53)  
60 CONTINUE  
WRITE (6,61) (NPROB(LA),LA=1,54)  
51 FORMAT (//,10X,INPUT ERROR IN DATA FILE - F(U,T) FROM MOD5',/



```

110X, 'PROGRAM TERMINATED'
53 FORMAT (//, 'ERROR IN PROBLEM STRUCTURING', /,
110X, 'RUN TERMINATED')
61 FORMAT (//, '3(10X,18A4,/) , 1CX,4(3X,'***END***',3X),
63 FORMAT ('1',10X,'EXECUTION BEGINS WITH CALL TO FOUT1')
64 FORMAT ('1',10X,'EXECUTION CONTINUES WITH CALL TO FOUT3')
STOP
END

```

SUBROUTINE INCON1  
REVISED 2 JUNE 1971

1 THIS ROUTINE DOES THE FOLLOWING  
2 SETS ALL ARRAY LOCATIONS TO ZERO  
3 DEFINES ALL PRINCIPAL DEFAULT OPTIONS, I.O.E.  
4 SZOH = 100.0 CM  
5 SA = 0.0 CM  
6 S2 = S3 = S4 = 0.0  
7 DIFLEN = 25.0  
8 DIS = 1  
9 IDATA = 1  
10 MODES = 1  
11 NSTAT = 2  
12 TIBX=0

1 READS ALL PROBLEM DEFINITION STATEMENTS FROM PUNCH CARDS  
2 AND ECHO PRINTS ALL INPUT DATA  
3 AND INFORMS THE USER WHAT ACTIONS THE PROGRAM WILL EXECUTE  
4 REAL LTAB(20), T1, T2, T3, T4, T5, T6, T7, T8, T9,  
1 T10, T11, T12, T13, T14, T15, T16, T17, T18, T19, T20 /  
1 DIMENSION INDATA(18)  
NAMELEN, MODES, NSTAT, EMIN, EMAX, SA, S1, S2, S3, S4, IDATA, NOUT, IBX, XBX,  
1 DIFLEN, ROUTINE WILL HAVE TO CONSIDER  
1 WITHIN INCON1 WE WILL HAVE TO ESTABLISH SOME FORM OF DEFAULT  
OPTION FOR THE MAX AND MIN ENERGYS THAT THE DETECTOR RESPONSE  
ROUTINE WILL HAVE TO CONSIDER  
COMMON/COM1/ XBX(100), C(150), E(150), TIM(20), DUMMY(150), TXBX(100)  
COMMON/COM4/ SOH, SA, S1, S2, S3, S4, SZERO, BUCKLE, EMAX, EMIN, DIFLEN, X  
COMMON/COM7/ S01, S02, S03, S04, AK11, AK12, AK13, AK14, AN3A, AN3B,  
1 AN3C, AN4, AN4, ANB4, AN2X, AK15, AK5X  
1 COMMON/COM2/ NT(20), INPROB(54), NANS(6), L(20), LTAB  
COMMON/COM5/ MODES, IS, IDATA, N, NVIR, NI, NF, NCALL, NOUT, IBX,  
1 NSTAT, MODE  
1 COMMON/CCM3/ FXT(100,20), FUTX(150,20), BUCK(6), FUT(6,71,20)  
COMMON/CCM6/ NTUP(6,20), NM(6,20), IKNT(6)







```

      WRITE (6,28) NCARD
      READ (5,29) NCARD
      DO 27 1IM=1,NCARD
      READ (5,30) (NDATAC(IA),IA=1,18)
      WRITE (6,31) (NDATAC(IA),IA=1,18)
      CONTINUE
      END OF DEBUG DATA PRINT OUT
      START BY READING THE PROBLEM TITLE
      READ (5,9) (NPROB(IA),IA=1,54)
      READ (5,DATAIN)
      INPUT OF PROGRAM LOGICAL CONTROLS
      READ (5,11) (NANS(I),I=1,6)
      READ (5,12) (L(I),I=1,20)
      DUPLICATE DATA PRINT OUT OF THE PROBLEM DEFINITION CARDS
      OUTPUT OF PROBLEM STATE STRUCTURE
      START PROBLEM LISTING ON A NEW PAGE
      WRITE (6,28)
      WRITE (6,64) (NPROB(IA),IA=1,54)
      WRITE (6,61) IS, IDATA, NOUT, IBX, SZERO, DIFLEN, SOH, SA, S1, S2, S3, S4
      WRITE (6,63) MODES, NSTAT
      IF (IBX.EQ.0) GO TO 68

```

```

      C PROTECTION STEPS TO TERMINATE PROGRAM EARLY IN EXECUTION IF
      C SOME ERRORS EXIST IN THE BASIC PROBLEM STATEMENT
      C COMPARE ALL PERTAINING DATA DEPENDING ON THE SUBROUTINES TO
      C BE DONE TO INSURE THAT ALL NECESSARY INPUT HAS BEEN PROVIDED
      C LOGICAL L(16)=0 INITIALLY-- IF INPUT DATA IS IN ERROR,
      C L(10) SET = TO 1 PROBLEM WILL BE TERMINATED IN MAIN PROGRAM
      C ON RETURN FROM THIS SUBROUTINE.
      C
      C IF ((IBX.LT.0) .OR. (IBX.GT.100) .OR. (IBX.GT.100)) L(10)=1
      C IF ((IBX.GT.100) .OR. (IBX.GT.0)) WRITE (6,62) (XBX(J),J=1,IBX)
      C GO TO 69
      C
      C 68 WRITE (6,67) XBX(1)
      C WRITE (6,70) (NANS(IA),IA=1,6)
      C WRITE (6,72) (L(I),I=1,20)
      C IBX=1

```

```

      C COMPLETE TEST SERIES TO EXAMINE MAIN INPUT CONTRAL DATA
      C
      C IF ((NOUT.GT.3) .OR. (NOUT.LE.0) .OR. (L(10)=1
      C IF ((IS.GT.5) .OR. (IS.LE.0) .OR. (L(10)=1
      C IF ((SOH.LE.SA) .OR. (SOH.LE.S1) .OR. (SOH.LE.S3) ) L(10)=1
      C IF ((NSTAT.GE.4) .OR. (NSTAT.LE.1) .OR. (L(10)=1
      C IF ((MODES.GT.1) .OR. (MODES.GT.6) .OR. (L(10)=1
      C TEST FURTHER MODES MUST BE REVISED IF PROGRAM IS MODIFIED
      C TO PROCESS MORE THAN 6 MODES OF DATA

```



C THIS TEST QUESTION MUST BE CHANGED TO REFLECT THE STORAGE  
C ARRAY CAPABILITIES OF COMMON BLOCKS COM3 & COM6

C PROGRAM WILL WRITE THE USER A SERIES OF MESSAGES TO DESCRIBE  
C WHAT ACTIONS THE PROGRAM WILL EXECUTE IN THE PERFORMANCE  
C OF THIS SPECIFIC JOB  
C WRITE (6,222) 1, IS  
C GO TO (201,202,203,204,2055), IS  
C WRITE (6,221) SOH, SA  
C GO TO 2041  
C WRITE (6,222) S1, SA, SOH  
C GO TO 2041  
C WRITE (6,223) SOH, DIFLEN  
C GO TO 2041  
C WRITE (6,2231) S1, SOH, DIFLEN  
C GO TO 2041  
C WRITE (6,2250)  
C 2055 WRITE (6,2250)  
C 2041 CONTINUE  
C GO TO (205,206,207), IDATA  
C 205 WRITE (6,224)  
C 206 WRITE (6,225)  
C 207 WRITE (6,226)  
C 208 WRITE (209,210,211), NOUT  
C 209 WRITE (6,227)  
C GO TO 212  
C 210 WRITE (6,228)  
C GO TO 212  
C 211 WRITE (6,229)  
C 212 CONTINUE  
C 213 WRITE (6,230) MODES  
C PROGRAM EXECUTION LISTING TO ADVISE THE USER THE ROUTINES  
C THAT WILL BE EXECUTED AND THE PRINCIPAL ITEMS OF DATA THAT  
C WILL BE USED BY EACH OF THEM  
C THIS WILL ASSIST THE USER IN ANY CASES WHERE SOME UNEXPECTED  
C EXECUTION ERROR MAY OCCUR  
C STATEMENTS ARE LISTED IN THE ORDER IN WHICH THE PROGRAM WILL  
C EXECUTE THEM  
C FOUT4 -- SPECTRAL DATA PLOT  
C 213 IF (L (20) \* EQ. 0) GO TO 3C1  
C WRITE (6,32C) MODES  
C FOUT5 -- FOURIER ADJUSTED SPECTRAL DATA PLOT  
C IF (L (17) \* EQ. 0) GO TO 3C2  
C WRITE (6,321) MODES  
C NPXT5=NANS(6)  
C IF (L (4) \* EQ. 1) WRITE (6,3211) NANS(6), (XBX(IJ), IJ=1, NPXT5)  
C FOUT1 - CALCULATE F(Y, T, X); = F(U, T)\*PHI(X) FOR INPUT DATA  
C







```

225 FORMAT ('/10X, "INPUT DATA FROM MAGNETIC TAPE")  

226 FORMAT ('/10X, "INPUT DATA FROM DISK")  

227 FORMAT ('/10X, "OUTPUT WILL BE LISTED ON THE PRINTER")  

228 FORMAT ('/10X, "OUTPUT PLOTS WILL BE PLOTTED ON THE PRINTER")  

229 FORMAT ('/10X, "GRAPHICAL OUTPUT WILL BE PLOTTED USING THE",  

230 1      'COMP PLOTTERS OF DATA TO PROCESS = ',2X,15)  

230 FORMAT ('/10X, "L(2U)=1",T25,"SUBROUTINE-FOUT4 WILL BE",  

230 1      'TORRECT SPECTRAL PLOT OF MOD-5 OUTPUT",  

230 1      'L(1X, "EXECUTED")/T25,"/UNCORRECT SPECTRAL PLOT OF MOD-5 OUTPUT",  

230 1      'L(1X, "VS U",/1?T35?15?3X,"HARMONIC MODES OF DATA")  

230 1      'L(17)=1?T25,"SUBROUTINE-FOUT5 WILL BE EXECUTED",  

230 1      'L/T25,"SPECTRAL PLOT OF MOD-5 OUTPUT ADJUSTED BY THE FOURIER",/,  

230 1      'L(1X, "CAL COEFFICIENT AT THE POINT X = Q.U",  

230 1      'L(1X, "L(2U)=1",T25,"SUBROUTINE-FOUT1 WILL BE EXECUTED",  

230 1      'L(1X, "AND A SUM OF ALL MODES WILL BE PLOTTED")  

232 FORMAT ('/10X, "L(18)=1",T25,"SUBROUTINE-FOUT1 WILL BE EXECUTED",  

232 1      'L/T25,"CALCULATION OF THE FUNCTION -F(U,T,X)=F(U,T)*PHI(X)",  

232 1      'L/T25,"FOR 1 X 1 X 1 DATA POINTS",/1?5X,"WITH OUTPUT VIA DEVICE",  

232 1      'L/T25,"FOR 1 X 1 X 1 DATA POINTS FOR GRAPHICAL OUTPUT",/,  

232 1      'L/T25,"DATA POINT POSITIONS ARE",/10?10X,15(F6.2?4X),/)  

321 FORMAT ('/10X, "L(14)=1",T25,"SUBROUTINE-FOUT2 WILL BE EXECUTED",  

321 1      'L/T25,"CALCULATION OF THE MEAN ENERGY AND FLUX-DENSITY PROFILE",  

321 1      'L/T25,"CALCULATION OF THE MEAN ENERGY FOR EACH TIME STEP",/,  

321 1      'L/T25,"TOTAL OF",2X,1X,"PLOTS WITH BOTH FUNCTIONS WILL BE",  

321 1      'L/T25,"A DONE")  

324 FORMAT ('/10X, "L(19)=1",T25,"SUBROUTINE-FOUT3 WILL BE EXECUTED",  

324 1      'L/T25,"CALCULATION OF A DETECTOR RESPONSE AT POSITION",1X,  

324 1      'L/T25,"WITH T25,A HALF THICKNESS OF",2X,F5.2?/,  

324 1      'L/T25,"DETECTOR WILL RESPOND TO NEUTRON ENERGY-EMAX=",1X,1PE10.3,  

324 1      'L/T25,"EMIN =",1X,1PE10.3,/,T25,"OUTPUT VIA DEVICE =",1X,  

324 1      'L/T25,"A DONE")  

325 FORMAT ('/10X, "RAMP FUNCTION-EXTERIOR SOURCE")  

325 FORMAT ('/10X, "L(4)=1",T25,"A SPECTRAL PLOT WILL BE DONE FOR",  

325 1      'L/T20,"POINTS ARE",/1?5X,F8.5?2X,)  

325 1      'L/T25,"NO PRINTED LISTING OF MEAN ENERGY",  

325 1      'L/T25,"OR FLUX PROFILE WILL BE DONE",  

325 1      'L/T25,"OVER RIDE OF OUTPUT SEQUENCE",  

325 1      'L/T25,"OVER RIDE OF PLOTTED")  

325 1      'L/T25,"OVER RIDE OF OUTPUT SEQUENCE",  

325 1      'L/T25,"MEAN ENERGY RATIO WILL BE PLOTTED")  

325 1      'L/T25,"RETURN"

```



SUBROUTINE ESTIM1  
REVISED 2 JUNE 1971

INITIAL ESTIMATING PROGRAM TO PROVIDE GUIDANCE ON TOTAL  
OF SIX INITIAL PROGRAM EXECUTION TIME WHILE ERROR CHECKOUT TO INFORM THE USER OF THE  
ANTICIPATED PROGRAM EXECUTION PARAMETERS  
EXECUTE THIS ROUTINE IMMEDIATELY AFTER READ10 OR READ11  
TO BEST ESTIMATE TOTAL EXECUTION TIME AFTER WE HAVE BETTER  
KNOWLEDGE OF THE TOTAL PROGRAM OPERATING PARAMETERS

```
COMMON/COM1/ XBX(100),U(150),E(150),T(20),DUMMY(150),TXBX(100)
COMMON/COM4/ SOH,SA,S1,S2,S3,S4,SZERO,BUCKLE,EMIN,DIFLEN,X
COMMON/COM7/ SOL,SO2,SO3,SO4,AK11,AK12,AK13,AK14,AN3A,AN3B,
1AN3C,AN4,AN4,ANB4,AN2X,AK15,AK5X
COMMON/COM2/ NT(20),NPROB(54),NANS(6),L(20),LTAB(20)
COMMON/COM5/ MODES,IS,IDA,T,N,VIR,NI,NF,NCALL,NOUT,IBX,
INSTAT,MODE
COMMON/COM3/ FXT(100,20),NFUTX(150,20),BUCK(6),FUT(6,71,20)
COMMON/COM6/ NTOP(6,20),NM(6,20),IKNT(6)
REAL ST(15)*0.0
WRITE(6,76)
WRITE(6,77)
DEFINE TOTAL TIME AS A ZERO VALUE PRIOR TO CALCULATION
SUMT = 0.0
ST(1) = 0.75
STIME FOR MAIN AND INCON1
```

C EXECUTION TIME FROM READ PROGRAMS

```
IF((IDA,EQ.1).AND.(NOUT,EQ.1)) ST(2)=MODES*1.25 + 0.50
IF((IDA,EQ.1).AND.(NOUT,EQ.3)) ST(2)=MODES*1.25 + 0.50
IF((IDA,EQ.2).AND.(NOUT,EQ.1)) ST(3)=MODES*1.25 + 0.50
IF((IDA,EQ.2).AND.(NOUT,EQ.3)) ST(3)=MODES*1.25 + 0.40
FOUT1 TIME
IF((L(18),EQ.1).AND.(NOUT,EQ.1)) ST(4)=MODES*IBX*12.50
IF((L(13),EQ.1).AND.(NOUT,EQ.3)) ST(4)=MODES*IBX*2.50
FOUT2 TIME ESTIMATE
IF((L(14),EQ.1).ESTIMATE ST(5)=10.0*MODES
FOUT3 TIME ESTIMATE
IF((L(19),EQ.1).AND.(NOUT,EQ.1)) ST(6)=MODES*2C.C
FOUT4 TIME FOR EXECUTION
IF((L(20),EQ.1).ST(7)=MODES*3.0
FOUT5 TIME ESTIMATE
IF((L(17),EQ.1).ST(8)=MODES*4.0
ESTIMATED TIME - ESTIM1
EST(9)=L.25
ESTIMATED TIME - MODEL 1
EST(10)=6.25*MODES
```



```

C      SUM ALL TIME STEP ESTIMATES
C      DO 63 IX2=1,10
C      63  SUMT=SUMT+IX2
C      NOW WRITE OUT TOTAL TIME ESTIMATES
C      WRITE (6,78) (ST(IA),IA=1,10),SUMT
C      76  FORMAT (1,1,10X,3CX,ESTIMATED,PROGRAM EXECUTION,TIME,,
C      77  FORMAT (1,10X,10X,10X,ESTIMATED TIME,,
C      13CX,SECONDS)
C      78  FORMAT (1,10X,MAIN+INCON1,T30,F8.3,10X,READ10',T30,F8.3,/,,
C      11CX,READ11,T30,F8.3,10X,F0UT1,F8.3,/,,
C      21CX,F0UT2,T30,F8.3,10X,F0UT3,T30,F8.3,/,,
C      31CX,F0UT4,T30,F8.3,10X,F0UT5,T30,F8.3,/,,
C      51CX,ESTIM1,T30,F8.3,10X,MODEL1,T3C,F8.3,/,,
C      61CX,TOTAL TIME = T30,F8.5)
C      RETURN
C      END

```

```

FUNCTION PHI1(IMODE,IS,XPNT)
REVISIED 10 JUNE 1971

COMMON/COM4/ SOH,SA,S1,S2,S3,S4,SZERO,BUCKLE,EMAX,EMIN,DIFFLEN,X
COMMON/COM7/ S01,S02,S03,S04,AK11,AK12,AK13,AK14,AN3A,AN3B,
1AN3C,AN3C,AN4,AN4,AN5A,AN5B,AK15,AK5X

DEFINITIONS:
IS = NUMBER OF THE HARMONIC MODE
IS = TYPE SOURCE GEOMETRY

NOTE THAT ALL NON-HARMONIC DEPENDENT FACTORS ARE PRECALCULATED
IN SUBROUTINE - MODELL AND STORED IN COMMON BLOCK COM7

XPNT
GO TO (1,2,3,4,5),IS
1 CONTINUE
IS=1
IMOD=2*IMODE-1
AKN1=IMOD*AK11
PHI1=(S01/SOH)*SIN(AKN1*SA)*COS(AKN2*(X-SA))/AKN1
GO TO 10
CONTINUE
2 IS=2
AKN2=(2*IMODE-1)*AK12
PHI1=(AN2X)*SIN(AKN2*SA)*COS(AKN2*(X-SA))
GO TO 10
CONTINUE
3 IS=3
M66=(-1)*((IMODE+1))

```



```

AKN3=AK13*IMODE
AN3=AN3C*(AKN3*M66M*AN3B-AN3A)/(1.0+(AKN3*DIFLEN)**2)
PHI1=AN3*COS(AKN3*X)
GO TO 10
CONTINUE
4 IS=4
M44M=2*I MODE-1
M55M=(-1)**(IMODE+1)
AKN4=AK14*M44M
AK5=AKN4*S1
AK6=AKN4*DIFLEN
AK7=1.0+(AKN4*DIFLEN)**2
SI X5=SIN(AK5)
CI S6=COS(AK5)
AN4=(AN4/AK7)*(CI S6-AK6*SI X5-(M55M*AK6*ANB4))
CN4=(AN4/AK7)*(SI X5-AK6*CI S6+M55M*ANB4)
XP1=AKN4*X
PHI1=AN4*COS(XPT)+CN4*SIN(XPT)
GO TO 10
CONTINUE
5 IS=5
PHI1=AK5*X*SIN(IMODE*AK15*(XPNT+SOH))/(1.0*IMODE)
CONTINUE
RETURN
END

```

SUBROUTINE MODEL1  
REVISED 12 JUNE 1971

THIS ROUTINE WILL:  
1 CALCULATE THE CONSTANT ( NON-HARMONIC ) FACTORS IN THE  
FOURIER POSITION EXPRESSIONS FOR THE INDICATED SOURCE  
GEOMETRY( IS ) AS STATED IN THE PROBLEM  
2 CALCULATE THE SUM OF THE INTEGRATED FOURIER EXPRESSION  
TERMS OF THE FOURIER EXPANSION FOR THE REMAINING TERMS  
3 ADJUST THE VALUE OF THE SOURCE STRENGTH ( SZERO ) TO OBTAIN  
AN APPROXIMATION TO WITHIN +/- 5% OF THE ORIGINAL SOURCE  
A LIMIT OF 100 ITERATIONS WILL BE DONE.  
A DEFAULT OF 3 ITERATIONS WILL BE DONE IF THIS IS NOT  
SPECIFIED BY THE VARIABLE NANS ( 4 ) IN THE INPUT DATA  
4 IF LOGICAL ( 9 ) = 1 A SIMPLE PLOT OF THE TOTAL SUM VS  
POSITION ( X ) WILL BE DONE ON THE PRINTER  
L( Q ) EXTERNAL CONTROL - PERFORM SIMPLE PLOT OF SOURCE  
FUNCTION ON THE PRINTER IN ROUTINE MODEL1  
NANS ( 4 ) - EXTERNAL CONTROL - LIMIT TO ALLOWABLE NUMBER  
OF ITERATIONS IN ROUTINE MODEL1

CC



```

COMMON/COM1/ XBX(100),U(150),E(150),TIM(20),DUMMY(150),TXBX(100)
COMMON/COM3/ FXT(100,20),FUT(150,20),BUCK(6),L(20),FUT(6,71,20)
COMMON/COM5/ NT(20),NPROB(54),NANS(6),L(20),LTAB(20)
COMMON/COM6/ MODE,IS,IDA,NT,NVIR,NI,NF,NCALL,NOUT,IBX,
INSTAT,MODE
COMMON/COM4/ SOH,SA,S1,S2,S3,S4,SZERO,BUCKLE,EMIN,EMAX,DIFLEN,X
COMMON/COM7/ SO1,SO2,SO3,SO4,AK12,AK13,AK14,AN3A,AN3B,
1AN3C,AN3C,AN4,ANB4,AN2X,AK15,AK5X
DIMENSION FX6(100)
EQUIVALENCE (DUMMY,FX6)

C   DEFINE ALL TERMS FOR THE FOURIER EXPANSION COEFFICIENTS
C   AS INITIAL CONDITIONS TO REDUCE COMPUTATION IN THE FUNCTIONS
C   SUBPROGRAM PHI1 FOR ALL FOUR ( 5 ) SOURCE TYPE GEOMETRIES

      CALL CANCEL(2)
      IF(NANS(4).LE.0) NANS(4)=3
      BUCKLE=BUCK(1)
      WRITE(6,96)
      IF(NANS(4).GT.10) NANS(4)=10
      PROTECT=STEP TO RETAIN INITIAL VALUE OF SZERO
      SPZERO=SZERO
      SP1P1=1.15*SZERO
      SP1P9=1.90*SZERO
      NFLAG=0
      PI=3.14159
      START OF INTERATION LOOP
      CONTINUE
      NFLAG=NFLAG+1
      WRITE(6,97) NFLAG,IS
      NFLAG=IS
      BRANCHING RETURN OF INTERATIVE PROCEDURE
      C   THE ROUTINE WILL DROP OUT OF THE INTERATION SERIES IF ONE OF
      C   THE FOLLOWING CONDITIONS IS MEET:
      C   1  CORRECTION FACTOR ( CORRF ) IS BETWEEN 1.05 AND 1.95
      C   2  SUM OF ADJUSTED TRUNCATED SERIES IS BETWEEN 1.1C AND 0.90 OF
      C   PARAMETER - SZERO
      C   3  TOTAL NUMBER OF INTERATION PASSES ( NFLAG ) EQUALS CONTROL
      C   VALUE INPUT - NANS(4)
      GO TO (201,202,203,204,205),IS
      201 CONTINUE
      C   IS=1
      SO1 = SZERO/2.0*SA
      AK11 = SQRT(BUCKLE)
      GO TO 205
      202 CONTINUE
      C   IS=2
      SO2=SZERO/2.0*SA
      DELTAZ=SO1 + PI/(2.0)*SQRT(BUCKLE))

```



```

AK12= PI/(2.0*DELT A2)
AK12=AK12*0.5
GO TO 215
CONTINUE
203 IS=3
AK13= SORT(BUCKLE)
AK13=AK13*0.5
AN3A=1.0/DIFLEN
AN3B=EXP(-1.0*SOH/DIFLEN)
SO3=SZERO/(2.0*DIFLEN*(1.0**2))
AN3C=(SO3/SOH)*((DIFLEN*(1.0**2))
GO TO 215
CONTINUE
204 IS=4
SO4=SZERO/(2.0*DIFLEN)*(1.0*EXP(-1.0*(SOH+SI)/DIFLEN))
DELT A4= SI + PI/(2.0*SQR((BUCKLE))
AK14= PI/(2.0*C*DELT A4)
AK14=AK14*0.25
ANC4= EXP(-1.0*SOH/DIFLEN)
ANB4=SINH(SI/DIFLEN)
ANA4=2.0*SO4*DIFLEN
ANB4=ANB4*ANC4
GO TO 215
IS=5
CONTINUE
AK15=SORT(BUCKLE)
AK5X=2.0*SZERO/(SOH*PI)
215 CONTINUE
START APPROXIMATION ANALYSIS TO DETERMINE THE ERROR IN
FOURIER APPROXIMATION FOR THE TOTAL MODES OF HARMONIC DATA TO
CONSIDER FOR THIS PROBLEM
GOTO (301,302,303,304,305),IS
CONTINUE
IS=1 CASE
SUM1=0.0
DO 95 MODE=1, MODES
AKN1=(2*MODE-1)*AK11
SUM1=SUM1+ SIN(AKN1*SOH)*SIN(4*KN1*SA)/(AKN1)**2
95 CONTINUE
SUM1=SUM1*2.0*SOI/SOH
ER1=SZERO-SUM1
CORF=1.0+COR1
WRITE(6,98) SZERO,IS,MODES, SUM1,ER1,COR1,CORF
IF ((SUM1.LE.SP1).AND.(SUM1.GE.SP2P9)) GO TO 1015
GO TO 315
CONTINUE
302

```



C      IS=2    CASE

    SUM2 = SOH-S1  
    DEL1 = MODE-1, MODES  
    DO 295 AK12\* (2\* MODE-1)  
    AKN2 = AK12\* SIN(AKN2\* SOH)\*COS(AKN2\* S1)  
    BX2 = 2\*SIN(AKN2\* S1)/((2\* MODE-1)\*2)  
    BX3 = SIN(AKN2\* DEL1)  
    BX4 = BX3\*( BX2-BX4)  
    BXSUM=SUM2+ BXSUM  
    CONTINUE

295    SUM2 = SUM2\*S02/ (SA\*( PI\*\*2))  
      ER2= SZERO-SUM2  
      CORF=1.0+COR2  
      WRITE(6,98) SZERO, IS, MODES, SUM2, ER2, COR2, CORF  
      IF (( SUM2.LE.SP1P1).AND.(SUM2.GE.SPZP9)) GO TO 1015  
      GOTO 315

303    C      IS=3    CASE

      ABX3 = 4.0\*C\*S03\*((DIFLEN)\*\*2)/(SOH\*AK13)  
      SUM3 = 0.0  
      DO 395 MODE=1, MODES  
      M35 = (-1)\*\*(MODE+1)  
      M36 = 2\*\* MODE-1  
      AKN3 = AK13\*M36  
      BX35 = M35\*(AKN3\*EXP(-1.0\*SOH/DIFLEN))\*M35-1.0/DIFLEN)  
      BX36 = M36\*((1.0+((AKN3\*DIFLEN)\*\*2))  
      BX37 = BX35/BX36  
      SUM3 = SUM3+ BX37  
      CONTINUE

395    SUM3 = SUM3\*ABX3  
      ERR3 = SZERO-SUM3  
      CORF= ER3/SUM3

      WRITE(6,98) SZERO, IS, MODES, SUM3, ERR3, CORF  
      IF (( SUM3.LE.SP1P1).AND.(SUM3.GE.SPZP9)) GO TO 1015  
      GOTO 315

304    C      IS=4    CASE

      SUM4 = 0.0  
      DO 495 MODE=1, MODES  
      AKN4=AK14\*(2\* MODE-1)  
      AX5=AKN4\*S1  
      AX6=AKN4\*DIFLEN  
      AX7=EXP(-1.0\*SOH/DIFLEN)\*SINH(S1/DIFLEN)  
      BX5=COS(AX5)-AX6\*SIN(AX5)+((-1)\*\*MODE)\*AX6\*AX7  
      BX6=SIN(AKN4\*SOH)/(1.0+((AX6)\*\*2))



```

BX5=BXS5*BXS6
SUM4=SUM4**2.0*SD4*DIFLEN
SER4=SZERO-SUM4
COR4=ER4/SUM4
CORF=1.0+COR4
WRITE(6,98) SZERO,IS,MODES,SUM4,ER4,COR4,CORF
IF ((SUM4.LE.SP1P1).AND.(SUM4.GE.SPZP9)) GO TO 1015
GO TO 315
1S=5
CONTINUE
SUM5=0.0
AX51=4.0*SZERO/(PI**2)
DO 595 MODE=1,MODE
M51=(-1)**(MODE+1)
M52=1+M51
M53=MODE**2
SUM5=SUM5+(1.0*M51*M52)/(1.0*M53)
SER5=SPZERO-SUM5
COR5=ER5/SUM5
CORF=1.0+COR5
WRITE(6,98) SZERO,IS,MODES,SUM5,ER5,COR5,CORF
IF ((SUM5.LE.SP1P1).AND.(SUM5.GE.SPZP9)) GO TO 1015
CONTINUE
IF ((CORF.LE.1.05).AND.(CORF.GE.0.95)) GO TO 1015
IF (NANS(4).EQ.0) GO TO 1015
SZERO=SZERO*CORF
GO TO 1001
CONTINUE
315
CONTINUE
315
IF ((CORF.LE.1.05).AND.(CORF.GE.0.95)) GO TO 1015
IF (NANS(4).EQ.0) GO TO 1001
CONTINUE
C
C END OF ERROR ANALYSIS ROUTINE
C A PLOT OF THE SUM OF ALL FOURIER POSITION DEPENDENT TERMS
C WILL BE PLOTTED ON THE PRINTER IF THE LOGICAL CONTROL VARIABLE
C L(9)=1 TO GIVE A SIMPLE SOURCE REPRESENTATION
C IF (L(9).EQ.0) GO TO 25
C PROTECTION STEP TO ZERO ALL STORAGE PRIOR TO PLOTTING SUM
C DO 36 IK5=1,100
C FX6(IK5)=0.0
C NPNTS=100
C DELX=SOH/4.9*0
C TXBX(1)=-1.0*SUH
C TXBX(50)=0.0
C TXBX(100)=+SOH
C DO 4 IKX=2,49
C TXBX(IKX)=TXBX(1) + DELX*(IKX-1)
C CONTINUE
4

```



SUBROUTINE READIO(LASTN)  
REVISED 2 JUNE 1971  
INPUT ALL BASIC STATE PARAMETERS FROM MOD-5 PROBLEM OUTPUT  
ECHO PRINT OF STATE STRUCTURE

READS ALL PUNCH CARD OUTPUT FROM MOD-5  
BY READING A BLOCK OF SIX DATA SETS FUNCTION F(U,T)  
FOR UP TO SIX MODES OF BUCKLING AND TWENTY TIME STEPS  
OF DATA PER BUCKLING MODE  
FOR THE INPUT DATA - A PRINTED OUTPUT IS GIVEN OF THE  
MOD 5 STATE STRUCTURE N(U,A) AND E(I,A) FOR  
ALL ENERGY/LETHERGY STATES OF THE SYSTEM  
THIS OUTPUT IS LISTED ONCE AFTER THE FIRST SET OF INPUT DATA  
A PRINT OUT LISTING TIME STEP DATA I.E. NT AND TIM(NT)  
IS GIVEN AT THE END OF EACH DATA INPUT SEQUENCE  
AFTER ALL DATA FOR THE PARTICULAR MODE IS DONE  
ROUTINE DOES A TEST COUNT TO DETERMINE WHICH  
DATA SETS WILL BE PROCESSED BY CHECKING THOSE PROGRAMS  
WHICH GAVE MUD 5 OUTPUT AT THE SAME VALUES OF NT  
A LISTING IS GIVEN AT THE END OF THE READ SEQUENCE  
OF THE NUMBER OF TIME STEPS THAT WILL BE RETAINED.







THEN DEFAULT BY PRINTING OUT ALL REMAINING DATA THAT  
 CANNOT BE PROCESSED INTO MIL SIX  
 MOD FIVE PRODUCES THE FUNCTIONS F(U,T)  
 DUMMY PARAMETER NRUN DEFINES INTERMEDIATE OR LAST DATA SET  
 IF NRUN = 1 THIS IS INTERMEDIATE DATA SET  
 IF NRUN = 2 THIS IS LAST DATA SET

```

IA=ISET = NCALL
NCALL = C
IKNT(ISET) = 0
WRITE(6,23)
21 NCALL = NCALL + 1
  IF (NCALL.EQ.21) GOTO 8
  READ(5,1,END=13,ERR=16) NRUN, TIM(NCALL), NT(NCALL),
  1 INTOP(IA,NCALL) NM(IA,NCALL)
  C MOD-5 ALWAYS PUNCHES CARDS FROM THE DENSITY VECTOR IN THE
  FORM (POP(IA,1) IA=1,NM) IN A FORMAT 1PE10.3 FOR 4 FIGURE DATA
  MIL=NM(IA,NCALL)
  IF (NRUN.EQ.1) IKNT(ISET) = IKNT(ISET) + 1

  C TEST SERIES TO DETERMINE THE TOTAL DATA SETS WITH NRUN = 1
  THIS SHOULD REDUCE THE PROCESSING OF USELESS DATA LATER IN THE
  PROGRAM CALCULATIONS
  C OUTPUT NEWS NOTE TO WRITE AT END OF THE CARD READING ROUTINE
  READ(5,2,END=15,ERR=16) (FUT(IA,IB,NCALL),IB=1,MIL)
  WRITE(6,5) NRUN, TIM(NCALL), NT(NCALL), NM(IA,NCALL)
  IF (NRUN.EQ.1) GO TO 21
  3 CONTINUE
  IF (NRUN.EQ.2) WRITE(6,7)
  C GO TO 15
  8 CONTINUE
  C THIS FIXES DATA SET 20 AS THE MAXIMUM AMOUNT FOR THE PROGRAM

MTEST = 0
9 WRITE(6,10)
11 READ(5,1,END=13,ERR=16) NNRUN,TTMN,NTTN,NTOPN,NMM
  IF (NNRUN.EQ.1) MTEST = MTEST + 1
  READ(5,2) (DUMMY(IA),IA=1,NVIR)
12 IF (NNRUN.EQ.1) GO TO 11
13 WRITE(6,7)
  IF (NTEST.GE.NCALL) NCALL=NTEST
  15 CONTINUE
  GO TO 19
  20 CONTINUE

```



```

C 16 LASTN=1
C 19 CONTINUE
      WRITE(6,91) MODES
      DO 90 IX1=1 MODES
      WRITE(6,92) IX1,IKNT(IX1)
90 CONTINUE

C 1 FORMAT(15,15,8,315)
C 2 FORMAT(10,15,5X,1PE10.3,5X,315)
C 5 FORMAT(//,1CX,*END OF INPUT DATA FILE - F(U,T)*)
C 7 FORMAT(//,1CX,*NCALL EXCEEDS 20 - MAXIMUM SPACE IN PROGRAM*)
C 10 FORMAT(//,1CX,*TIME STEP DATA //,11X,NRUN,*)
C 23 FORMAT(NT,6X,NT,2X,INTOP,3X,NM,*)
C 1 9X,TIM(NT)
C 43 FORMAT(15)
C 44 FORMAT(18A4)
C 45 FORMAT(14I5,15,8)
C 46 FORMAT(8(1PE10.3))
C 49 FORMAT(10X,2(15,5X,2(1PE10.3,8X),6X)))
C 50 FORMAT(10X,18A4)
C 51 FORMAT(10X,NUMBER OF REAL STATES =',3X,15,/,
C 11CX,*NUMBER OF VIRTUAL STATES =',15,/,
C 21CX,*NUMBER OF ISOTOPES =',9X,15,/,
C 31CX,*NUMBER OF ISOTOPES =',2X,15,/,
C 41CX,*BUCKLING =',15X,F10.5)
C 53 FORMAT(*1*10X,*INPUT DATA FROM MOD 5*)
C 55 FORMAT(//,8X,(2(*STATE!,9X,U-LETHARGY!,6X,ENERGY--EV!,5X)))
C 56 FORMAT(//,10X,15)
C 57 FORMAT(10X,15)
C 91 FORMAT(//,1UX,15,2X,*,TOTAL USEFUL DATA SETS FROM MOD-5*,/,,
C 1 15X,DATA SET!,5X,SETS!)
C 92 FORMAT(17X,13,7X,13)
      RETURN
END

```

SUBROUTINE FOUT1(XPNT)  
REVISED 2 JUNE 1971

CCCCCCCC

SUBROUTINE FOUT1  
CALCULATES THE POSITION FUNCTION F(U,T,X) FOR A SERIES OF SIX  
BUCKLING SETS AT A TIME.  
IF ONLY SIX BUCKLING MODES ARE CONSIDERED, THIS ROUTINE CAN DO  
THE TIME-ENERGY-POSITION FUNCTION FOR AS MANY POINTS AS DESIRED  
OF THIS ONE DIMENSIONAL MODEL -- SLAB GEOMETRY



FUTURE CHANGES FOR THIS ROUTINE WILL BE ABLE TO DEAL WITH  
 ALMOST ANY QUANTITY OF INPUT DATA BY THE EARLIER CALCULATION  
 OF A NEW FUNCTION CALLED PHIX(NN,MM)  
 THIS FUNCTION WILL BE BASED ON A TOTAL NUMBER OF POINTS FOR  
 SOME MAXIMUM NUMBER OF BUCKLING MODES--MM  
 FUTURE PROCESSING WILL FOLLOW THE SEQUENCE  
 1 DEFINE TOTAL MODES OF DATA TO PROCESS  
 2 CALCULATE THE POSITION FUNCTION PHIX(NX,NMODES)  
 3 CALCULATE/UPDATE THE FUNCTION F(NU,NT)=F(NU,NT)\*PHIX(NX,IMODE)+F(NU,T,X)-OLD  
 4 TERMINATE THE CALCULATE RUN WHEN ALL DATE HAS BEEN EXHAUSTED  
 5 PROVIDE OUTPUT VIA THREE DEVICES

PHIX(NXPNT,NXMODES)  
 AFUTX(NXPNT,NU,NT)  
 CURRENT FUNCTION -- FUTX(NU,NT) DEALS WITH ONLY ONE POSITION  
 VALUE OF THE POSITION POINT-X. AT THIS TIME

```
COMMON/COM1/ XBX(100),U(150),E(150),TIM(20),DUMMY(150),TXBX(100),  

COMMON/COM4/ SOH,SA,S1,S2,S3,S4,SZERO,BUCKLE,EMAX,EMIN,DIFLEN,X  

COMMON/COM7/ S01,S02,S03,S04,AK11,AK12,AK13,AK14,AN3A,AN3B,  

1AN3C,AN4,AN4,ANB4,AN2X,AK15,AK5X  

COMMON/COM2/ NT(20),APROB(54),NANS(6),L(20),LTAB(20)  

COMMON/COM5/ MODES,IS,IDA,N,NVIR,NI,NF,NCALL,NOUT,IBX,  

1NSTAT,MODE  

COMMON/COM3/ FXT(100,20),FUTX(150,120),BUCK(6),FUT(6,71,20)  

COMMON/COM6/ NTOP(6,20),NM(6,20),IKNT(6)
```

```
REAL INMM  

REAL LABEL(20),T1,T2,T3,T4,T5,T6,T7,T8,T9,T10,T11,T12,T13,T14,T15,T16,T17,T18,T19,T20/  

1REAL*8 TIMST(4),NT=1-5,POINT-1,POINT-2,POINT-3,POINT-4,  

1REAL*8 APNT(6),POINT-5,POINT-6/  

1REAL*8 BMODES(20),1- MODE,2-MODE,3-MODES,4-MODES,  

1-5-MODES,6-MODES,7-MODES,8-MODES,9-MODES,  

1-11-MODES,12-MODES,13-MODES,14-MODES,  

3-15-MODES,16-MODES,17-MODES,18-MODES,  

4-20-MODES/  

1REAL*8 INAME(12),/MILLST-B', '0X-M', 'F(U,T,X)', '-FLUX AT',  

1'X VS U-, 'LETHRGY', 'MOD-2 CA', 'CALCULATION', 'N,  

2DIMENSION FFUTX(150)  

EQUIVALENCE (FFUTX,DUMMY)
```



```

C CALCULATE THE FUNCTION F(U,T,X) = F(U,T)*PHI1(X) FOR ALL FUT
N5=NSTAT
X1=XPTN

```

```

C CHANGE TO FOUT1 CHANGE AFTER ROUTINE IS SET TO PROCESS MORE
C THAN SIX INPUT DATA DECKS OR TAPE BY PROVIDING PROTECTED
C STORAGE FOR THE FIRST MODE OF BUCKLING

```

```

BUCK1 = BUCK(1)
DO 2 I MOD=1, MODES
NCALL1 = IKNT(IMOD)
DO 2 NCX=1, NCALL1
MIL1 = NTOP(IMOD, NCX) + 1
MIL2 = MIL1 + NM(IMOD, NCX)
DO 2 M=MIL1, MIL2
FUTX(M, NCX)=FUT(IMOD, M, NCX)*PHI1(IMOD, IS, X1)
1 + FUTX(M, NCX)
2 CONTINUE

```

```

OVER RIDE STEP TO BY PASS THE OUTPUT ROUTINES ON CALLS FROM
OTHER SUBROUTINE AND SUB PROGRAMS OF MOD 2
OVER RIDE PARAMETER IS LOGICAL -- L(16)
IF (L(16) = 1) THE PRINT ROUTINE AND ALL OTHERS ARE SKIPPED
IF (L(16) .EQ. 1) GO TO 25

```

```

WRITE OUT ALL VALUES OF FUTX(M, NCX)
BEFORE CALLING THE CAL COMP PLOTTER
NCALL = IKNT(1)
WRITE (6,90) X, NVIR, NCALL

```

```

PLOT P SEQUENCE
DO 93 I12=1, NCALL
WRITE (6,94) I12, T(I12), NT(I12)
WRITE (6,95) (FUTX(I13, I12), I13=1, NVIR)
93 CONTINUE

```

```

IF (NCUT .NE. 3) GO TO 25
IF THE TEST VARIABLE NOT EQ 3 BY PASS REST OF SEQUENCE
DEFINE SCALE PARAMETERS FOR THE DRAW PLOT ROUTINE
PAGES TO BE PLOTTED VIA THE CAL COMP
TOTAL LIMITS ON GRAPH PLOTTING FOR THIS ROUTINE IS DEFINED
VIA THE INPUT CONTROL VARIABLE NANS (1) = (( THE TOTAL GRAPH
DEFINE A DEFAULT VALUE FOR NANS(1) OF 1
IF (NANS(1) .LT. 1) NANS(1) = 1
EXSCALE = U(NVIR-1)/8.
YSCALE = FUTX(2,1)/8.
IXUP = C
IYRIGHT = C
MODXAX = C

```



```

MODYAX = 0
IWIDE = 9
IHIGH = 10
IGRID = 1
KKNT = NCALL
KTEST = NCALL/5 + 1
KK1 = 0
DO 8 IKT = 1, KTEST
  MODIFICATION(10) = BMODES(MODES)
  INAME(10) = NAMS(2)
  NPXNT = NAMS(2)
  INAME(11) = APNT(NPXNT)
  IF(IKT.LE.4) INAME(12) = TIMST(IKT)
  IF(IKT.GT.4) INAME(12) = TIMST(4)
  DO 5 IXM=1,5
    KK1 = KK1 +
    MODCUR = 2
    IF(IXM.EQ.1) MODCUR = 1
    IF(IXM.EQ.5) MODCUR = 3
    IF(KK1.EQ.NCALL) MODCUR=3
    IF(KK1.GT.2) GO TO 3
    INMM = LABEL(KK1)
    IF(KK1.GT.20) INMM = LABEL(20)
  3 CALL DRAW USING THE WORKING VECTOR FFUTX
    USE OF THE INPUT CONTROL PARAMETERS LIMITS GRAPH
    TO SOME MULTIPLE OF FIVE DATA PLOTS IF POINT PLOTTING
    PARAMETER USED IS NAMS(1)
    TEST QUESTION TO CHECK IF( NAMS(1).EQ.1 .AND. IKT.EQ.6) GT TO10
    INITIAL TEST QUESTION FORM:::THE FUTURE:::
    SECOND FORM TO USE IN THE FUTURE:::
    NTST2 = NAMS(1)*5 + 1
    IF(NTST2.EQ.IKT) GO TO 10
    PARAMETER NTST2 DEFINES THE TOTAL NUMBER OF
    TIMES MODES THAT PLOT WILL BE CALLED
  THREE DEFUALT OPTIONS TO DROP OUT OF THE CALL DRAW PROGRAM
  DO 4 ICX = 1, NVIR
  4 FFUTX(ICX) = FUTX(ICX,KK1)
    IF((NAMS(1).EQ.1).AND.(IKT.EQ.6)) GO TO 10
    IKT1 = IKT
    IF(NVIR.GE.30) IKT1 = 0
  1 CALL DRAW (NVIP,U,FFUTX,MODCUR,IKT1,INMM,I NAME,EXSCAL,YSCALE,
    IXUP,IYRIGHT,MODYAX,IWIDE,INMM,I NAME,EXSCAL,YSCALE,
    IF(LAST.EQ.2) GO TO 10
    IF(LAST.EQ.3) GO TO 10
  C

```



```

IF (KK1.EQ.NCALL) GO TO 10
ON REACHING MAXIMUM GRAPH OUTPUT-DROP OUT OF THE LOOP
5 CONTINUE
8 CONTINUE
10 CONTINUE
125 FORMAT ('1',1CX,'CALCULATION OF THE POSITION FUNCTION (U,T,X)',  

1/ '2CX, '2X, 'F10.5,2X, 'WITH',2X,15,2X,'STATES AT',2X,  

215,2X, 'TIMES')
94 FORMAT ('/','1CX,'RUN NUMBER',2X,15,2X,'TIME =',2X,E15.8,  

12X,'SECONDS',5X,'NT =',2X,15)
95 FORMAT (41(/,10X,5(1PE10.3,5X)))
END

```

SUBROUTINE FOUT2  
REVISED 16 JUNE 1971

CORRECTED ROUTINE TO CALCULATE THE MEAN ENERGY OF THE  
NEUTRONS AT A POINT X(I) AND THE SUM OVER ALL LETHARGY STATES  
OF THE NEUTRON DISTRIBUTION AT THE SAME POINT AND PLOT BOTH  
LETHARGY SUM  
F(U,T,X,MODE) = F(U,T,MODE)\*PHI(X,MODE)  
F(U,T,X) = SUM OVER MODES--F(U,T,X,MODE)  
OVER ALL STATES U(LETHARGY) TO GIVE THE RESULTANT  
FUNCTION --  
F(X,T) = SUM OVER LETHARGY - F(U,T,X)  
MEAN ENERGY CALCULATION  
EMEAN(X,T) = SUM OVER LETHARGY-F(U,T,X)\*E(U)  
COMMON BLOCKS 3 AND 6 MUST BE REVISED TO HANDLE MORE THAN  
SIX HARMONIC MODES OF THE FUNCTIONAL DATA INPUT  
DIMENSION VSIG(150)  
EQUIVALENCE (DUMMY(1),VSIG(1))  
REAL LABEL  
REAL\*8 ENERGY,'ENERGY','ENGX','RATIO','  
REAL\*8 FNGY,'NEUTRON','FNGX','DENSITY','  
REAL\*8 ARUNT(9)/'SYSTEM-1','SYSTEM-2','SYSTEM-3','SYSTEM-4',  
1\*SYSTEM-5,'SYSTEM-6','SYSTEM-7','SYSTEM-8','SYSTEM-9',  
1\*REAL\*8 MTIT(12)/'FUNCTION',F(X,T,)'  
1\*REAL\*8 TSIMX(20)/'NT-1','NT-2','NT-3','NT-4','NT-5',  
1\*NT-6,'NT-7','NT-8','NT-9','NT-11','NT-12',  
2\*NT-13,'NT-14','NT-15','NT-16','NT-17','NT-18',  
3\*NT-19,'NT-20','NT-21','NT-22',  
REAL\*8 BMODES(6)/'1-MODE','2-MODES','3-MODES','4-MODES',







```

NPNTS=99
DELX2=SOH/49.0
TXBX(1)=(-1.0)*SOH
TXBX(99)=SOH
DC3_I2=2.98
TXBX(I2)=TXBX(1)+DELX2*(I2-1)
3 TXBX(50)=0.0
CONTINUE
CP2
      WRITE (6,1002)
      C DEFINE INITIAL CONSTANTS FOR THE ROUTINE
      C NONE = 1
      BKBK=BUCK(1)
      NTMAX=IKNT(1)
      NMAX = NVIR-1
      C IMPROVEMENT TO INCREASE EXECUTION SPEED OF CALCULATION
      DO 321 NXPT=1,NPNTS
      DO 321 KMOD=1,MODES
      FUTX(NXPT,KMOD)=PH11(KMOD,IS,TXBX(NXPT))
321 CONTINUE
CP3
      WRITE (6,1003)
      C START OF CALCULATIONS
      DO 250 INT=1,NTMAX
      DO 245 IXPT=1,NPNTS
      AX=0.0
      BX=0.0
      CX=0.0
      CX5=0.0
      DO 240 IMODE=1,MODES
      C PROTECTION STEP WITHIN CALCULATION LOOP TO REDUCE UNNECESSARY
      C CALCULATIONS
      NP3=IKNT(IMODE)
      IF (INT.GT.NP3) GO TO 240
      MIL1=INT*OP(IMODE,INT)
      MIL2=NW(IMODE,INT)
      MIL3=MIL1+MIL2
      DO 235 MIL1=1,MIL2
      MIL4=MIL1+NU1
      IF (MIL4.GT.NMAX) GO TO 235
      BX1=FUT(IMODE,NU1,INT)*FUTX(IXPT,IMODE)
      C
      CX5=CX5+BX1
      AX=AX+BX1*VSIG(MIL4)
      BX=BX+BX1*E(MIL4)
325 CONTINUE
3240 CONTINUE
      C DEFAULT PROTECTION STEP FOR PARAMETER CX5

```



THIS VARIABLE IS TO REPRESENT THE NEUTRON POPULATION DENSITY  
 FOR A PARTICULAR TIME STEP VALUE  
 IN GENERAL IT SHOULD BE A VALUE BETWEEN 0.001 AND 1.000  
 FOLLOWING IS TO PROTECT AGAINST AN ACCIDENTAL ATTEMPT TO DEVIDE  
 BY ZERO OR SOME OTHER UNUSUAL VALUE THAT MIGHT ACCIDENTLY  
 COME UP.

TEST CHECK AND DEFAULT SET:  
 IF ((CX5.LE.0.001).OR.(CX5.GE.1.001)) CX5=1.000

EMEAN(IXP,INT)=BX/CX5  
 FXT(IXP,INT)=AX/CX5

245 CONTINUE  
 250 CONTINUE  
 CP4  
 WRITE(6,1004)  
 DO 253 NT4=1,NTMAX  
 IF ((IS.EQ.1).OR.(IS.EQ.3)) ZA=EMEAN(1,NT4)  
 IF ((IS.EQ.2).OR.(IS.EQ.4).OR.(IS.EQ.5)) ZA=EMEAN(50,NT4)  
 C DEFAULT TEST CHECK ICN ZA VALUE  
 WRITE(6,601)  
 IF (ZA.LE.0.0) GO TO 255  
 DO 252 IX4=2,NPNTS  
 FUTX(IX4,NT4)=EMEAN(IX4,NT4)/ZA  
 CONTINUE

252 FUTX(1,NT4)=1.0  
 GO TO 253

255 WRITE(6,271)  
 DO 256 IX5=1,NPNTS  
 FUTX(IX5,NT4)=.90

256

253 CONTINUE  
 CP5  
 WRITE(6,1005)  
 BY PASS THE CAL COMP ROUTINE PORTION IF NOUT NOT EQUAL 3  
 IF (NOUT.NE.3) GO TO 12  
 C START PREPARATION OF DATA FOR GRAPH PLOTTING ROUTINES  
 C TWO ITEMS OF DATA WILL BE PLOTTED FOR EACH TIME STEP  
 LAST=0  
 ITYPE=C  
 XCAL=SOH/8.0  
 IHIGH=10  
 IWIDE=8  
 IYRT=C  
 IGRID=1  
 FXT(I,J)=FLUX PROFILE FOR X AND T  
 EMEAN(I,J)=MEAN ENERGY DATA VERSUS X AND T  
 FUTX(I,J)=RATIC MEAN ENERGY(X,T)/MEAN ENERGY(0.0,T)



```

C L(K) = 1 EXTERNAL CONTROL - BYPASS THE PRINT SEQUENCE IN FOUT2
C L(7) = 1 EXTERNAL CONTROL - PLOT FLUX PROFILE VS X ON CAL COMP
C L(8) = 1 ADJUSTMENT OF TITLING FOR GRAPHS
C MTTIT(10) = BMODES(NM,NT) DATA
C LOAD DUMMY WIT FXT(NX,NT) DATA
C DO 155 NX1=1,NTMAX
C LABEL=LTAB(NXT1)
C MTTIT(12) = TIMX(NXT1)
C IXUP1=0
C YSCAL1=FXT(1,NXT1)/4.0
C MODDC1=1
C MODDX1=0
C IXUP1=C
C IYRT=5
C MDDY1=0
C PROTECTICN STEP FOR GRAPH PLOTTING
C PLOTTING PROTECTION STEP TO MAINTAIN ONLY POSITIVE VALUES
C IF (L(7).EQ.0) GO TO 1175
C DO 162 NP5=1,NPNTS
C DUMMY(NP5)=FXT(NP5,NXT1)
C IF (DUMMY(NP5).LE.0.001) DUMMY(NP5)=0.0
C 162 CONTINUE
C CALL DRAW TO PLOT FIRST GRAPH FOR THIS TIME STEP
C MTTIT(2)=FNGY
C MTTIT(3)=FNGX
C CALL DPAW(NPNTS,TXBX,DUMMY,MODC1,ITYPE,LABEL,MTIT,XCAL,
C 1YSCAL1,IXUP1,IYRT,MODX1,MODY1,IWIDE,IHIGH,IGRID, LAST)
C 1175 IF (L(8).EQ.0) GO TO 1155
C MODDX2=2
C YSCAL2=0.50
C IXUP2=6
C MODDC2=3
C MDDY2=0
C DU163 NP6=1,NPNTS
C DUMMY(NP6)=FUTX(NP6,NXT1)
C CALL TO DRAW TO PLOT THE MEAN ENERGY RATIO
C MTTIT(2)=FNGY
C MTTIT(3)=FNGX
C CALL DPAW(NPNTS,TXBX,DUMMY,MODC2,ITYPE,LABEL,MTIT,XCAL,
C 1YSCAL2,IXUP2,IYRT,MODX2,MODY2,IWIDE,IHIGH,IGRID, LAST),
C 1176 C16
C INTERNAL PROTECTION STEP TO LIMIT TOTAL GRAPHICAL OUTPUT TO
C THE CAL COMP PLOTTER VIA VARIABLE NANS ( 3 ) TO COMPARE
C WITH THE INDEXING VARIABLE --NXT1
C KPLOTS IS A DUMMY VARIABLE DEFINED WITH A DEFAULT VALUE OF 1
C CONTINUE

```



```

155 IF ( NNTT .GE. KPLOTS ) GO TO 156
156 CONTINUE
      WRITE (6,1006)
      PLOT OUTPUT ROUTINE USING THE PRINTER
      C 12 IF (NOUT .EQ. 3) GO TO 16
      NCALL = IKNT(1)
      DO 15 J = 2, NCALL
      DO 13 I = 1, NPTS
      DUMMY(I) = FXT(I,J)
      13 MODC = 0
      CALL PLOT (TXBX, DUMMY, NPTS, MODC)
      WRITE (6,14) J, NT(J), TIM(J)
      15 CONTINUE
      PRINT-PLOT THE MEAN ENERGY VALUES OF THE PRINTER
      DO 615 J = 2, NCALL
      DO 613 I = 1, NPTS
      DUMMY(I) = EMEAN(I,J)
      613 MODC = 0
      CALL PLOT (TXBX, DUMMY, NPTS, MODC)
      WRITE (6,617) J, NT(J), TIM(J)
      615 CONTINUE
      PRINT-PLOT THE MEAN ENERGY RATIO
      DO 715 J = 2, NCALL
      DO 713 I = 1, NPTS
      DUMMY(I) = FUTX(I,J)
      713 MODC = 0
      CALL PLOT (TXBX, DUMMY, NPTS, MODC)
      WRITE (6,717) J, NT(4), TIM(J)
      715 CONTINUE
      PRINT OUT RESULTS ON THE PRINTER ONLY
      C PRINT OUT ALL VALUES OUT IN COLUMNS OF SIX TIME STEPS
      C L(5) = 1 EXTERNAL CONTROL-BYPASS THE PRINT SEQUENCE IN FOUT2
      C 16 IF (L(5) .EQ. 1) GO TO 23
      C OVER RIDE CONTROL TO SKIP THE PRINT SEQUENCE
      C 35 CONTINUE
      NCALL = IKNT(1)
      IF (NCALL .GT. 6) NTT = 6
      IF (NCALL .LE. 6) NTT = NCALL
      WRITE (6,22) (NT(I), I = 1, NTT), (TIM(I), I = 1, NTT)
      WRITE (6,501)
      DO 17 I = 1, NPTS
      17 WRITE (6,21) TXBX(I), (FXT(I,J), J = 1, NTT)
      WRITE (6,502)
      DO 67 I2 = 1, NPTS
      67 WRITE (6,21) TXBX(I2), (EMEAN(I2, J2), J2 = 1, NTT)
      DO 93 I2 = 1, NPTS

```



```
93 WRITE(6,21) TXBX(I2), FUTX(I2,J2), J2=1,NTT
```

```
IF(NCALL.LE.6) GO TO 23
```

```
IF(NCALL.LE.12) NTT=12
```

```
IF(NCALL.LE.22) NTT=NCALL
```

```
IF(NCALL.LE.50) NTT=7,NTT)
```

```
18 WRITE(6,21) NPNTS
```

```
19 WRITE(6,502)
```

```
20 WRITE(6,21) NPNTS
```

```
21 WRITE(6,503)
```

```
22 WRITE(6,21) TXBX(I2), (EMEAN(I2,J2), J2=7,NTT)
```

```
23 CONTINUE(//,10X,15,2X,'TOTAL POINTS PLOTTED USING DRAW',/,
```

```
31 11CX15,2X,'TOTAL CALLS MADE TO DRAW')
```

```
32 14 FORMAT(//,10X,'DATA PLOT NUMBER',2X,15,'OX, 'TEST RUN NT =',
```

```
33 12X,15,'/10X, TIME =',2X,E11.4,2X,'SECONDS')
```

```
34 21 FORMAT(11X,F6.2,6X,(1PE10.3,4X))
```

```
35 22 FORMAT(11X,F6.2,8X,6(4X,'NT = ',13,1X),/,
```

```
36 122X,6(1X,E10.3,3X),/,22X,6(5X,SEC,6X))
```

```
45 DO 68 I2=1, NPNTS
```

```
46 WRITE(6,21) TXBX(I2), (FUTX(I2,J2), J2=7,NTT)
```

```
47 IF(NCALL.LE.18) GO TO 23
```

```
48 IF(NCALL.LE.18) NTT=18
```

```
49 IF(NCALL.LE.18) NTT=NCALL
```

```
50 IF(NCALL.LE.18) {NT(I), I=13,NTT}, (TIM(J), J=13,NTT)
```

```
51 WRITE(6,21) NPNTS
```

```
52 WRITE(6,502)
```

```
53 WRITE(6,21) NPNTS
```

```
54 WRITE(6,503)
```

```
55 WRITE(6,21) TXBX(I2), (EMEAN(I2,J2), J2=13,NTT)
```

```
56 IF(NCALL.LE.18) GO TO 23
```

```
57 NXIT=NXT(I2,22) NT(19),NT(20),NXT,NXT,TIM(19),TIM(20)
```

```
58 WRITE(6,501)
```

```
59 DO 20 I2=1, NPNTS
```

```
60 WRITE(6,21) TXBX(I2), FUT(I,19), FUT(I,20)
```

```
61 NTT=NCALL
```

```
62 DO 70 I2=1, NPNTS
```

```
63 WRITE(6,21) TXBX(I2), (FUTX(I2,J2), J2=19,NTT)
```

```
64 WRITE(6,502)
```

```
65 WRITE(6,21) NPNTS
```

```
66 WRITE(6,21) TXBX(I2), (FUTX(I2,J2), J2=19,NTT)
```

```
67 CONTINUE(//,10X,15,2X,'TOTAL POINTS PLOTTED USING DRAW',/,
```

```
68 11CX15,2X,'TOTAL CALLS MADE TO DRAW')
```

```
69 14 FORMAT(//,10X,'DATA PLOT NUMBER',2X,15,'OX, 'TEST RUN NT =',
```

```
70 12X,15,'/10X, TIME =',2X,E11.4,2X,'SECONDS')
```

```
71 21 FORMAT(11X,F6.2,6X,(1PE10.3,4X))
```

```
72 22 FORMAT(11X,F6.2,8X,6(4X,'NT = ',13,1X),/,
```

```
73 122X,6(1X,E10.3,3X),/,22X,6(5X,SEC,6X))
```



SUBROUTINE FOUT3  
REVISED 2 JUNE 1971







```

4 CONTINUE
      TXBX(3)=S3
      C DETECTOR RESPONSE FOLLOWING DATA MUST BE STATED IN THE BASIC PROBLEM DEFINITION
      C FOR THIS ROUTINE TO RUN SUCCESSFULLY
      S3      C LOCATION OF THE MID POINT OF THE DETECTOR
      S4      C THE HALF WIDTH OF THE DETECTOR
      SET LOGICAL VARIABLE L(16)=1 AND CALL FCUT1 TO RETURN THE
      VALUES F(U,T,X) TO THIS ROUTINE FOR EACH POSITION POINT
      L(16)=1
      NCALL3=IKNT(1)
      DO 41 IXPNT=1,5
      CALL FOUT1(TXBX(IXPNT))
      DO 42 NIVAL=1,NCALL3
      BX3=0.
      NUMX=NVIR=1
      DO 43 NIVAL=1,NUMX
      BX3=BX3+FUTX(NIVAL,NTVAL)*VSIG(NIVAL)
      43 CONTINUE
      FXT(IXPNT,NTVAL)=BX3
      42 CONTINUE
      USE ARRAY DUMMY TO STORE THE SUM OVER ALL TIME STEPS AT THE
      USE ARRAY TXBX TO STORE DATA FOR INDIVIDUAL TIME STEPS FROM MOD-5
      VALUE OF NCALL3 IS THE MAX TIME STEP POINT VALUE FROM MOD-5
      CX=0.Q
      DO 47 NTVAL=1,NCALL3
      CX1=FXT(1,NTVAL)+FXT(5,NTVAL)
      CX2=FXT(2,NTVAL)+FXT(4,NTVAL)
      CX3=(CX1+3.0*CX2+5.0*FXT(3,NTVAL))/9.0
      DUMMY(NIVAL)=CX3
      47 CONTINUE
      C REVISED 18 MAY 1971
      C TIME INTEGRATION SEQUENCE FOR THE DETECTOR RESPONSE FUNCTION
      DX=0.Q
      DO 497 NIVAL=2,NCALL3

```



```

NTMIN=NTVAL-1
DX1=(DUMMY(NTVAL)+DUMY(NTMIN))*0.50
DX2=TIM(NTVAL)-TIM(NTMIN)
DX3=DX1*DX2
DX=DX+DX3
DXBX(NTVAL)=DX

CONTINUE(1)=0.0
TXBX(1)=0.0
NOW START THE DATA PLOT TO SHOW THE DETECTOR RESPONSE
AS COMPUTED ABOUT THE MIDPOINT OF THE DETECTOR
CAL COMP PLOTTER ROUTINE ROUTINE
FOU 3-- DETECTOR RESPONSE FUNCTIONS -- OUTPUT ROUTINES
ALL OUTPUT WILL BE GIVEN BY THE THREE FORMS
PRINT ALL VALUES : PLOT ON PRINTER : PLOT USING CAL COMP
STANDARD NPS OUTPUT METHODS
NTCALL=IKNT(1)
OUTPUT ON PRINTER
PRINT OUT ALL VALUES OF THE TWO FUNCTIONS
F(X,T) AND F(X) VERSUS TIME PRIOR TO DOING THE
FINAL PLOT USING THE NPS PLOTTING ROUTINE // DRAW //
USE THE PRINTER AND THE ROUTINE -- PLOT-P TO OBTAIN A
PRELIMINARY DATA PLOT
WRITE (6,154) ((TXBX(IA),TIM(IA)),IA=1,NTCALL)
MODC=0
13 MAY 1971 ---- SIMPLE FIX-UP // REMOVE FUNNY RESULTS ON FIRST RUN
TEST SERIES TO PLOT DATA DESPITE FUNNY RESULTS ON FIRST RUN
IKFT=0
DO 1000 IK6=1,NTCALL
PDUM1(IK6)=DUMMY(IK6)
IF (PDUM1(IK6)*LT.0.0) IKFT=IKFT+1
PDUM2(IK6)=TXBX(IK6)
IF (PDUM2(IK6)*LT.0.0) IKFT=IKFT+1
1000 CONTINUE
IF (IKFT*LE.2) GO TO 1011
DO 1010 IK6=1,NTCALL
PDUM1(IK6)=(-1.0)*PDUM1(IK6)
1010 CONTINUE
1011 CONTINUE
IF (IKFT*LE.2) GO TO 1013
DO 1012 IK6=1,NTCALL
PDUM2(IK6)=(-1.0)*PDUM2(IK6)
1012 CONTINUE
1013 CONTINUE
C PROVIDE A LOG10 CONVERSION ON THE TIME SCALE TO FIT THE PLOTTER
SC=2.0
DO 1515 IK6=2,NTCALL

```



```

1015 ATIMX(IK6)=SC*(ALOG10(TIM(IK6))+10.0)
CONTINUE
ATIMX(1)=0.0
CALL PLOT(PDUM2,ATIMX,NTCALL,MODC)
WRITE(6,155) ((DUMMY(IA),TIM(IA)),IA=1,NTCALL)
155 PLOT(PDUM1,ATIMX,NTCALL,MODC)
C DEFINE SCALE VARIABLES FOR THE DRAW--CAL CUMP PLOTTER
OVER RIDE CALL DRAW SEQUENCE UNLESS NOUT = 3
IF (NOUT.NE.3) GO TO 1003
MODX=0
MODY=0
IYRT=9
IXUP=0
ITYTYPE=1
EXSCL=1.0/9.0
C PREPARE DATA TO PLOT REFLECTED AND REVERSED AXIS ON THE PLOTTER
DO 151 KPNT=1,NTCALL
151 TXBX(KPNT)*LE*0.001 TXBX(KPNT)=0.0
TXBX(KPNT)=(-1.0)*TXBX(KPNT)
CONTINUE
IWIDE=9
IHIGH=15
IGRID=1
C PLOT F(X,T) ON THE GRAPH FIRST
MODC1=1
CALL DRAW(NTCALL,TBX,TIM,MODC1,IYRT,MODX,IWIDE,IHIGH,IGRID,LAST),
152 IYSCAL,IXUP,ISCALE,ITYPE1,LBL1(1),TITLE,EXSCL,
C PLOT F(X) VS TIME
PREPARE DATA BY REFLECTING ABOUT Y-AXIS
DO 152 KPNT2=1,NTCALL
152 DUMMY(KPNT2)*LE*0.001 DUMMY(KPNT2)=0.0
DUMMY(KPNT2)=(-1.0)*DUMMY(KPNT2)
CONTINUE
153 ITYPE2=2
MODC2=3
CALL DRAW(NTCALL,DUMMY,TIM,MODC2,ITYPE2,LBL2(2),TITLE,EXSCL,
153 IYSCAL,IXUP,IYRT,MODX,MODY,IWIDE,IHIGH,IGRID,LAST),
154 154 FORMAT('/*,10X,'PRELIMINARY DATA PLOTS',/10X,'TIME INTERGRATED',/
11X,'FUNCTION-F(X) VS TIME',/20X,PF(X),/T35,TIME-SECONDS,/,/
22C(/18X,1PE10,3))
155 155 FORMAT('/*,10X,'PLUT OF DETECTOR RESPONSE FUNCTION',/10X,
120X,PF(X,T) VS TIME',/20X,PF(X,T),/T35,TIME-SECONDS,/
22C(/18X,1PE10,3),RETURN
END

```



SUBROUTINE FOUT4 - PREFORMS A DATA CHECK OF EACH INPUT MODE VALUES OF THE NEUTRON DENSITY VECTOR FOR EACH BUCKLING MODE TO TEST ALL VALUES TO DETERMINE THOSE DATA POINTS ABOVE AN ARBITRARY MINIMUM VALUE FOR THE PARAMETER BX

THIS VERSION USES BX = 0.001

AS A MINIMUM POINT VALUE IS NORMALIZED TO A VALUE OF 1.00

IN PROGRAM MOD-5

FOUT4 INITIALLY SCANS THE INPUT DATA TO OBTAIN A MAXIMUM OF THE POINT VALUES OF THE ARRAY--- FUT (IMODE,NU,NT) FOR ALL POINTS ABOVE THE LIMIT OF THE PARAMETER BX

INITIAL DRAFT OF THIS ROUTINE (10 APR 71) WILL PROVIDE A PLOT OF THE INPUT SPECTRUM OF FUT (I,J,K) VS NU--LETHARGY STATE FOR EACH MODE OF INPUT DATA

INITIAL DRAFT IS BASED ON THE MODEL UTILIZING ONLY A MAXIMUM OF SIX INPUT MODES OF DATA FROM MOD-5

ROUTINE WILL PLOT THE FUNCTION--SC\*ANALOG10 (FUT (IMODE,NU,NT)) VS. LETHARGY VALUE -- U USING THE CAL COMP PLCTTER

COMMON/CCM1/ XBX(100),U(150),E(150),TIM(20),DUMMY(150),TXBX(150)  
COMMON/CCM4/ SO1,SA1,S1,S2,S3,S4,SZER0,BUCKLE,EMIN,EMAX,DIFLEN,X  
COMMON/CCM7/ SO1,SO2,SO3,SO4,AK11,AK12,AK13,AK14,AN3A,AN3B,  
1AN3C,AN4,AN5,ANB4,AN2X,AK15,AK5X  
COMMON/CCM2/ NT(20),NPROB(54),NANS(6),L(20),LTAB(20)  
COMMON/CCM5/ MODES,IS,IData,N,NVIR,NI,NF,NCALL,NOUT,IBX,  
1INSTAT,MODE  
COMMON/CCM3/ FXT(100,20),FUTX(150,20),BUCK(6),FUT(6,71,20)  
COMMON/CCM6/ NTOP(6,20),NM(6,20),IKNT(6)

REAL\*8 RMOD(6), MODE-1, MODE-2, MODE-3, MODE-4, MODE-5,  
1 MODE-6,  
1 REAL LABEL  
REAL\*8 TITLE(1,2), 'MILLS', 'BOX--M', 'FOUT4', 'TEST', ONE', /  
1, , , FUT(U,T, ), VS, U, LETHARGY, /

C START OF THE MULTIPLE PROCESSOR MODE TO CLEAN UP INPUT DATA  
C AND PROVIDE GRAPHICAL OUTPUT OF NEUTRON SPECTRUM  
C

SC = 3.0  
BX = 1.0E-03  
NPTS = NVIR - 3  
ITYPE = 0  
EXSCAL = 1.0  
NMAX = NVIR - 1



```

YSCALE = U(NMAX)/14.0
IXUP = 0
IYRT = 0
MODX = 0
MODY = 0
IWIDE = 9
IHIGH = 15
IGRID = 1
DO 300 J=1, MODES
  TITLE(6) = 'MODS(J)'
  DO 92 NXX = 1, 15
    DUMNY(NXX) = 0.0
  FX = 0.0
  NCALL1 = IKNT(J)
  DO 555 K=2, NCALL1
    NNU = NM(J,K)
    NNU DEFINES THE MAXIMUM STEP VALUES INPUT FROM MOD 5
    FX = 0.0
    IF(FUT(J,NU,K)*GT.FX) FX=FUT(J,NU,K)
    THIS SCAN WILL DETERMINE THE MAX VALUE IN THE INPUT DATA
    CX= BX*FX
    THIS WILL DETERMINE THE MINIMUM VALUES TO CONSIDER
    WRITE(6,108) CX,BX,FX,K
    108 FORMAT(1X,1X,CX=1PE10.3,2X,1PE10.3,2X,1PE10.3,
    12X,FX=1PE10.3,2X,1PE10.3,2X,1PE10.3,
    C
    MIL1 = NTOP(J,K)
    MIL2 = NN(J,K)
    MIL3 = MIL1 + MIL2
    DO 51 NU=1, NNU
      AX = FUT(J,NU,K)
      SET MINIMUM VALUE FOR LOG CONVERSION
      IF(AX.LT.BX) AX = BX
      MIL4 = MIL1 + NU
      FUTX(MIL4,K) = SC*ALOGIC(AX) + 9.0
    C
    51 CONTINUE
    C THIS STEP WILL AUTOMATICALLY ZERO OUT ANY POINT AT THE
    C LOW END OF THE SPECTRUM(HIGH-ENERGY; LOW - LETHARGY) TO
    C THE CORRECT PROFILE OF THE SPECTRUM
    C IF (MIL1.LE.2) CUTO 53
    DO 52 KN = 1, MIL1
      FUTX(KN,K) = 0.0
    C
    52 CONTINUE
    C
    53 CONTINUE
  C

```



SUBROUTINE FOUT5  
REVISED 15 JUNE 1971  
SUBROUTINE FOUT5 PROVIDES A FOURIER POSITION FUNCTION  
WEIGHTED SPECTRAL PLOT OF THE INPUT DATA FROM MOD-5 AND  
PLOTS F(I MODE, U, T)\*PHI (I MODE, X1) VS U (LET HAGY) FOR EACH  
MODE OF INPUT DATA AND A SUMMED SPECTRUM OF ALL MODES OF INPUT  
L ( 4 ) = 1 DEFINES THE EXCEPTIONAL CASE WHERE THE USER WOULD  
DESIRE TO PLOT THE SPECTRAL RESPONSE FUNCTION



```

FOURIER EXPANSION COEFFICIENT FOR EACH MODE AND POSITION OTHER
TN4N AT THE DEFAULT CASE OF X = 0.0, CONSIDER IS
TCTAL POINTS ( POSITION ) THAT THE PROGRAM WILL FOR THE NUMBER OF
INPUTS USING THE CONTROL VARIABLE NANS ( 6 ) FOR THE ACTUAL POINTS
POINTS AND THE INPUT ARRAY XBX( 1 ) FOR THE ACTUAL POINTS
PRE CALCULATE THE FOURIER SPACE DEPENDENT VALUES FOR EACH MODE
FOR EACH POINT X1 = XBX( 100 ), U( 150 ), E( 150 ), DUMMY( 150 ), TXBX( 1
COMMON/COM1/ XBX, SA, S1, S2, S3, S4, SZERO, BUCKLE, EMIN, DIFLEN,
COMMON/COM4/ SO1, SO2, SO3, SO4, AK1, AK12, AK13, AK14, AN3A, AN3B,
COMMON/COM7/ AN3C, AN4, ANA4, ANB4, AN2X, AK15, AK5X
COMMON/COM2/ NT( 20 ), NPROB( 54 ), NANS( 6 ), L( 20 ), LTAB( 20 )
COMMON/COM5/ MODES, IS, IDATA, N, NVIR, NI, NF, NCALL, NOUT, IBX,
COMMON/COM3/ FXT( 150, 20 ), FUTX( 150, 20 ), BUCK( 6 ), FUT( 6, 71, 20 )
COMMON/CCM6/ NTOP( 6, 20 ), NM( 6, 20 ), KNT( 6 )
REAL*8 RTITLE( 12 ), MILLS, BOX, M, FOUT5, FOUT6, FOUT7, FOUT8, FOUT9, FOUT10
1*TOTAL*8 ANAME( 12 ), ALL, MODES, M, /, /, /, /, /, /, /, /, /, /
1*REAL*8 BMODE( 5 ), /, 1-2, 1-2-3, 1-2-3-4, 1-2-3-4-5, /, 123456, /
REAL*8 APNT( 9 ), /, POINT-1, POINT-2, POINT-3, POINT-4, POINT-5, POINT-6, POINT-7, POINT-8, POINT-9, /
1*POINT-5, /, MODE-1, MODE-2, MODE-3, MODE-4, MODE-5, /
1*REAL*8 AMODE( 6 ), /, MODE-1, MODE-2, MODE-3, MODE-4, MODE-5, /
1*MODE-6, /
1 REAL LABEL DIMENSION PHIMX( 6 )
START EXECUTION DEFINE PARAMETERS--SCALE FACTORS, ETC FOR EXECUTION
IPNTR=NANS( 2 )
IF ( ( IPNTR.LE.0 ) .OR. ( IPNTR.GE.9 ) ) IPNTR=9
RTITLE( 12 )=APNT( IPNTR )
ANAME( 12 )=APNT( IPNTR )
NRUNS = NANS( 6 )
THIS DEFINES THE TOTAL NUMBER OF DATA POINTS THAT WILL
HAVE STCENTRAL RESPONSE DATA POINTS PLOTTED
IF ( ( NRUNS.LE.6 ) .OR. ( NRUNS.GE.3 ) ) GO TO 550U
IF ( L( 4 ).EQ.0 ) NRUNS = 1
DO 5001 NTIMX = 1, NRUNS
X1= XBX( NTIMX )
IF ( L( 4 ).EQ.0 ) X1=0.0
WRITE( 6, 1013 ) X1
1013 FORMAT( 6, 1, 1, X, * POSITION POINT FOR SPECTRAL PLOT = *, 5X, F8.2 )
DO 1011 JXX=1, MODES
PHIMX( JXX )=PHI1( JXX, IS, X1 )
WRITE( 6, 1012 ) JXX, PHI1( JXX, IS, X1 )
1012 FORMAT( 6, 1, 1, X, * MODE = *, 2X, 15, * PHIX = *, E15.8 )
1011 CONTINUE

```



```

SC = 3.0
TYPE = 0
EXSCAL = 1.0
NPNTS = NVIR-1
TESTS SERVES CORRECTION FOR MIN SCALING VALUE TO PLOT
THIS SHOULD BE ADJUSTED TO BECOME PART OF THE INPUT
NAME DATA-- DATA -- TO ALLOW THE USER TO SELECT
AN ESTIMATE OF THE BEST VALUE --- THIS WILL REQUIRE ADDITION
TO LINE OF THE ESTABLISHMENT OF A DEFUALT OPTION IN INCN1
PLUS THE ESTABLISHMENT OF A DEFUALT OPTION IN INCN1
REVISED 05 JUNE 1971
BX=1.0E-05
YSCALE = U(NPNTS)/14.0
IXUP = 0
IXRT = 0
MCDX = 0
MCDY = 0
IWIDE = 9
IHIGH = 15
IGRID = 1
DO 92 NX=1,15
DO 92 NT=1,20
FUTX(NXX,NTT) = 0.0
92 CONTINUE IX=1, MODES
A NAME(6) = AMODE(IX)
NCALL1 = IKNT(IX)
DO 55 K=2, NCALL1
DO 49 IX=1,15
DUMMY(IX) = 0.0
49 FX=0.0
MIL1 = NTOP(IX,K)
MIL2 = NM(IX,K)
MIL3 = MIL1+MIL3
DO 51 NU = 1, MIL2
MIL4 = MIL1+NU
AX = PHIMX(IX)*FUT(IX,NU,K)
FUTX(MIL4,K) = AX + FUTX(MIL4,K)
IF ((AX*LE-BX) AX = BX
C CX*ALOGIC(AX) + 15.0
IF (CX.GT.9.0) CX = 9.0
C IF (CX.LT.BX) CX=0.0
DUMMY(MIL4) = CX
51 CONTINUE
LABEL = LTAB(K)

```



```

IF (( K.EQ.2) MODC = 1 LAST = C
IF (( K.EQ.3) AND. (LAST.EQ.C)) MODC = 2
CALL DRAW(NPNTS,DUMMY,U,MODC,ITYPE,LABEL,ANAME,EXSCAL,YSCALE,
1 IXUP,IYRT,MODY,IWIDE,IHIGH,IGRID,LAST)
IF ((LAST.EQ.1) GO TO 105
IF ((LAST.EQ.3) GO TO 105
55 CONTINUE
105 IF ((LAST.EQ.1) WRITE (6,106)
IF ((LAST.EQ.3) WRITE (6,107)
300 CONTINUE
C LAST = C
C SECOND PART OF THE DRAW ROUTINE TO PLOT THE SUM OF THE TOTAL
C HARMONIC EXPANSION OF THE NEUTRON DENSITY FUNCTIONS
NKN = IKNT(1)
NKN = MCDE$-1
IF (NX.EQ.0) MX=1
IF (NX.EQ.6) = BMODE(MX)
DO 310 NXT = 2, NKN
DO 309 NNTX = 1, NPNTS
AX = FUTX(NNTX,NXT)
IF (AX.EQ.BX) AX = BX
DUMMY(NNTX) = SC*ALOG10(AX) + 9.0
CONTINUE
309 LABEL = LTAB(NXT)
IF (NXT.EQ.2) MODC = 1
IF ((NXT.EQ.2) AND. (LAST.EQ.C)) MODC = 2
IF (NXT.EQ.NKN) MODC = 3
LAST = 0
NFLAG=NFLAG+1
IF (NFLAG.GE.4) GO TO 550
CALL DRAW(NPNTS,DUMMY,U,MODC,ITYPE,LABEL,RTITLE,EXSCAL,YSCALE,
1 IXUP,IYRT,MODY,IWIDE,IHIGH,IGRID,LAST)
IF ((LAST.EQ.1) GO TO 320
IF ((LAST.EQ.3) GO TO 320
310 CONTINUE
IF ((LAST.EQ.1) WRITE (6,106)
IF ((LAST.EQ.3) WRITE (6,107)
106 FORMAT ('//',1X,'LAST = 1',3X,'RECHECK DATA')
107 FORMAT ('//',1X,'LAST = 3',3X,'RECHECK DATA')
5501 CONTINUE
5506 CONTINUE
RETURN
END

```



SUBROUTINE READ11(LASTN)  
 REVISED 5 MAY 1971  
 THIS SUBROUTINE HAS NOT BEEN COMPLETELY TESTED AND DEBUGGED  
 READ11 INPUTS THE SAME TEST DATA AS ROUTINE READ10  
 BY PROCESSING NINE (9) TRACK MAGNETIC TAPE  
 INITIALLY AT A DENSITY OF 800 BPI  
 IDENTICAL INPUT/OUTPUT MESSAGES AND DATA ARE FURNISHED  
 AS BY READ10

```

REAL MTITLE (54)
DIMENSION NNRUN(20),NNT(20)

C COMMON/COM1/ XBX(100),U(15C),E(15C),TIM(20),DUMMY(150),TXBX(100)
C COMMON/COM4/ S01,SA1,S1,S2,S3,S4,SD4,SD3,SD2,SD1,SO4,SO3,SO2,SO1
C COMMON/CCM7/ AN3C,AN4C,AN5C,AN6C,AN7C,AN8C,AN9C,AN10C,AN11C,AN12C,AN13C,AN14C,AN15C,AN16C,AN17C,AN18C,AN19C,AN20C
C COMMON/CCM2/ N1,N2,N3,N4,N5,N6,N7,N8,N9,N10,N11,N12,N13,N14,N15,N16,N17,N18,N19,N20
C COMMON/COM5/ MODES,IS,IDA,TAB(20),LTAB(20),NOUT,IBX,INSTATION,MODE
C COMMON/COM3/ FXT(100,20),FUTX(15C,20),BUCK(6),FUT(6,71,20)
C COMMON/COM6/ NTCP(6,20),NM(6,20),IKNT(6)

C SET DEFAULT OPTION FOR LASTN = 0
LASTN=0
DO 8 ISET=1,MODES
  READ(4,20,END=8,ERR=9) (MTITLE(IA),IA=1,54)
  READ(4,21,END=8,ERR=9) N,NVIR,NI,NF,BUCK(ISET)
  IF (ISET*GT•1) GO TO 1
  READ(4,22) (U(IA),IA=1,NVIR)
  READ(4,22) (E(IA),IA=1,NVIR)
  READ(4,22) (E(IA),IA=1,NVIR)
  1 NCOUNT=0
  2 NCALL1=NCALL1+1
  IF (NCALL1*GE•21) GO TO 18
  READ(4,23,END=8,ERR=9) IRUN,TIM(NCALL1),NNT(NCALL1),
  1 INTOP(ISET,NCALL1),NM(ISET,NCALL1)
  NNRUN(NCALL1)=IRUN
  MILI=NM(ISET,NCALL1)+1
  READ(4,24,END=8,ERR=9) (FUT(ISET,IB,NCALL1),IB=1,MILI)
  IF (IRUN*EQ•1) GO TO 2
  END OF TAPE READ IN LOOP FOR NORMAL TIME STEP INPUT DATA

C DEFAULT OPTION IF MORE THAN 20 DATA SETS IS PRODUCED PER
C TIME STEP MOD 5
C WRITE(6,27)
C NKNT=0
  
```



```

19 READ (4,23,ERR=9) IXRUN,XTIM,NXNT,NXNM
NKNT = NXNM + 1
MIL2 = (4,24,END=8,ERR=9) (DUMMY(IA),IA=1,MIL2)
READ (4,28) IXRUN,XTIM,NXNT,NXTOP,NXNM
WRITE (NKNT,GT=20) GO TO 9
IF (NKNT .GT. 20) GO TO 9
PROTECTION STEP TO GET OUT OF INFINITE TAPE READ LOOP
IF (IXRUN .EQ. 1) GO TO 19
WRITE (6,29) NKNT
END OF TAPE READ ROUTINE TO INFORM USER OF DATA SETS INPUT
DO (1SET,GT=1) GO TO 4
DO (IX=1,NCAL,1)
NT(IX) = NNT(IX)
IF (1SET .GT. 1) GO TO 4
3 WRITE (6,25) (MTITLE(IA),IA=1,54)
IF (1SET .GT. 1) GO TO 6
NPX= (NVIR/2) + 1
DO 5 IK1= 1, NPX
IK1P = NPX+IK1
5 WRITE (6,30) IK1, U(IK1), E(IK1), E(IK1P), U(IK1P)
6 WRITE (6,31) NCALLI
DO 7 IX= 1, NCALLI
7 WRITE (6,29) IX, TIM(IX), NNT(IX), NTOP(ISET,IX), NM(ISET,IX)
8 CONTINUE (6,32)
8 WRITE (6,34)
GO TO 50
9 LASTN=1E
50 CONTINUE
20 FORMAT (18A4)
21 FORMAT (415,E15.8)
22 FORMAT (8E10.3)
23 FORMAT (15,E15.8,315)
24 FORMAT (3(10X,18A4,/,)
25 FORMAT (/10X,1N=,9X,15,/10X,'NVIR =',6X,15,'/10X,'NI =',
18X,15,'/10X,'NF =',8X,15,'/10X,'BUCKLE =',3X,F10.5),
26 FORMAT (//10X,'TOTAL INPUT TIME DATA FRCM MOD-5',//,
110X,'EXCEEDS 120 DATA SETS')
27 FORMAT (//10X,15,E15.3,315)
28 FORMAT (//10X,2(13,5X,2(1PE10.3,5X)) )
30 FORMAT (10X,2(13,5X,2(1PE10.3,5X)) )
31 FORMAT (10X,2(13,5X,2(1PE10.3,5X)) )
32 FORMAT (//13X,'SET',3X,'TIME',3X,'INT',3X,'NTOP',2X,'NM')
110X,3(8X,'***',3X),/,'10X,'END OF INPUT DATA SET',/,,
34 FORMAT (/,10X,'END OF TAPE READ CYCLE')

```



RETURN  
END



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## 13. ABSTRACT

The slowing down of fast neutrons was analyzed by a multi-group method of discrete time and energy states coupled with a spatial harmonic expansion method to determine the neutron density in a homogeneous, isotropically scattering slab. Five neutron source geometries were studied for both a fissioning and a non-fissioning system.

Numerical results were obtained for the neutron flux, mean neutron energy and the neutron spectra for the one dimensional system using a harmonic mode expansion of up to six terms to determine the time-energy-space dependence.



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